



Gravitational Wave Lensing: Current Searches and Future Prospects

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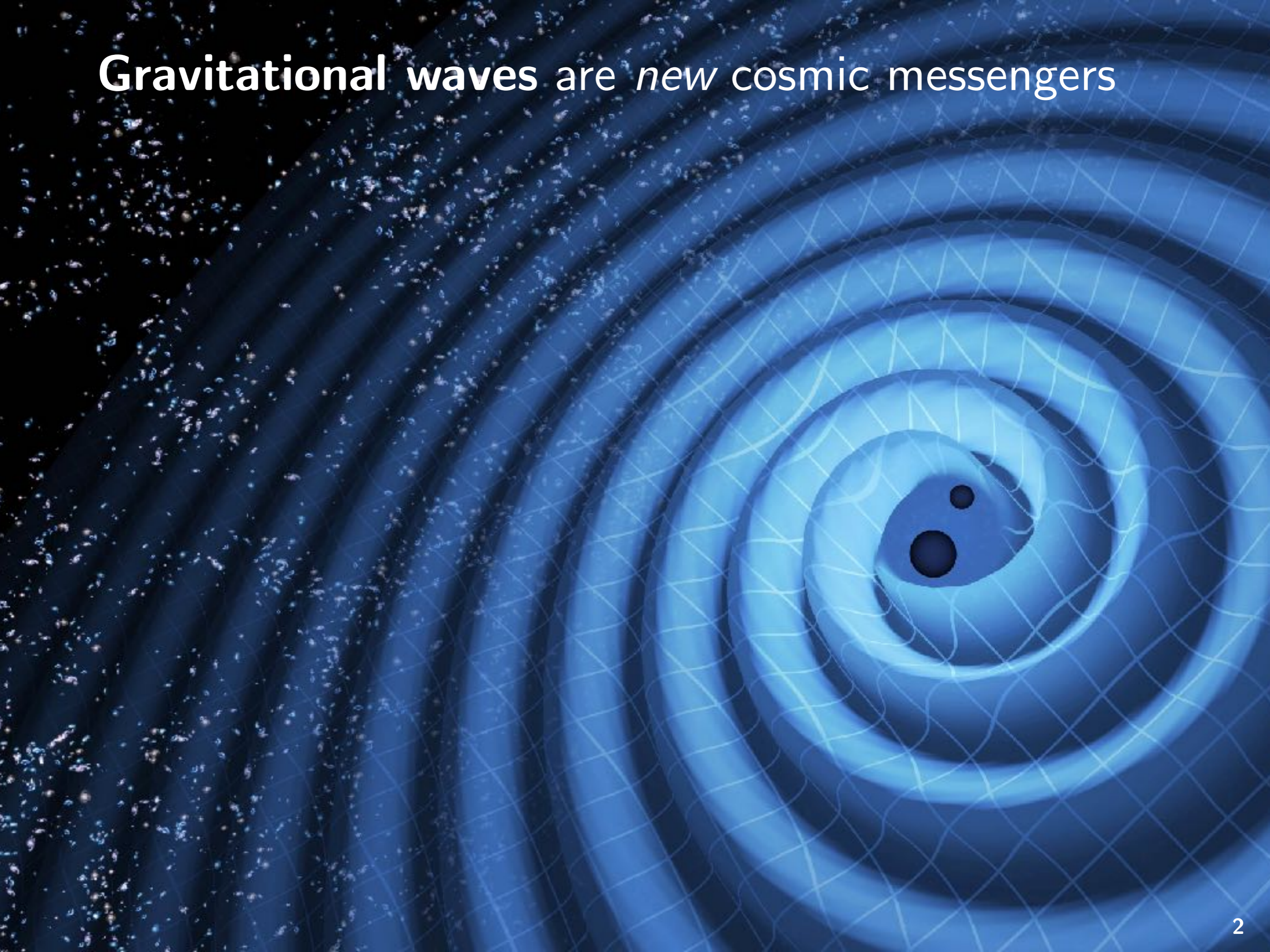
VILLUM FONDEN



KØBENHAVNS
UNIVERSITET

[Gustav Klimt]

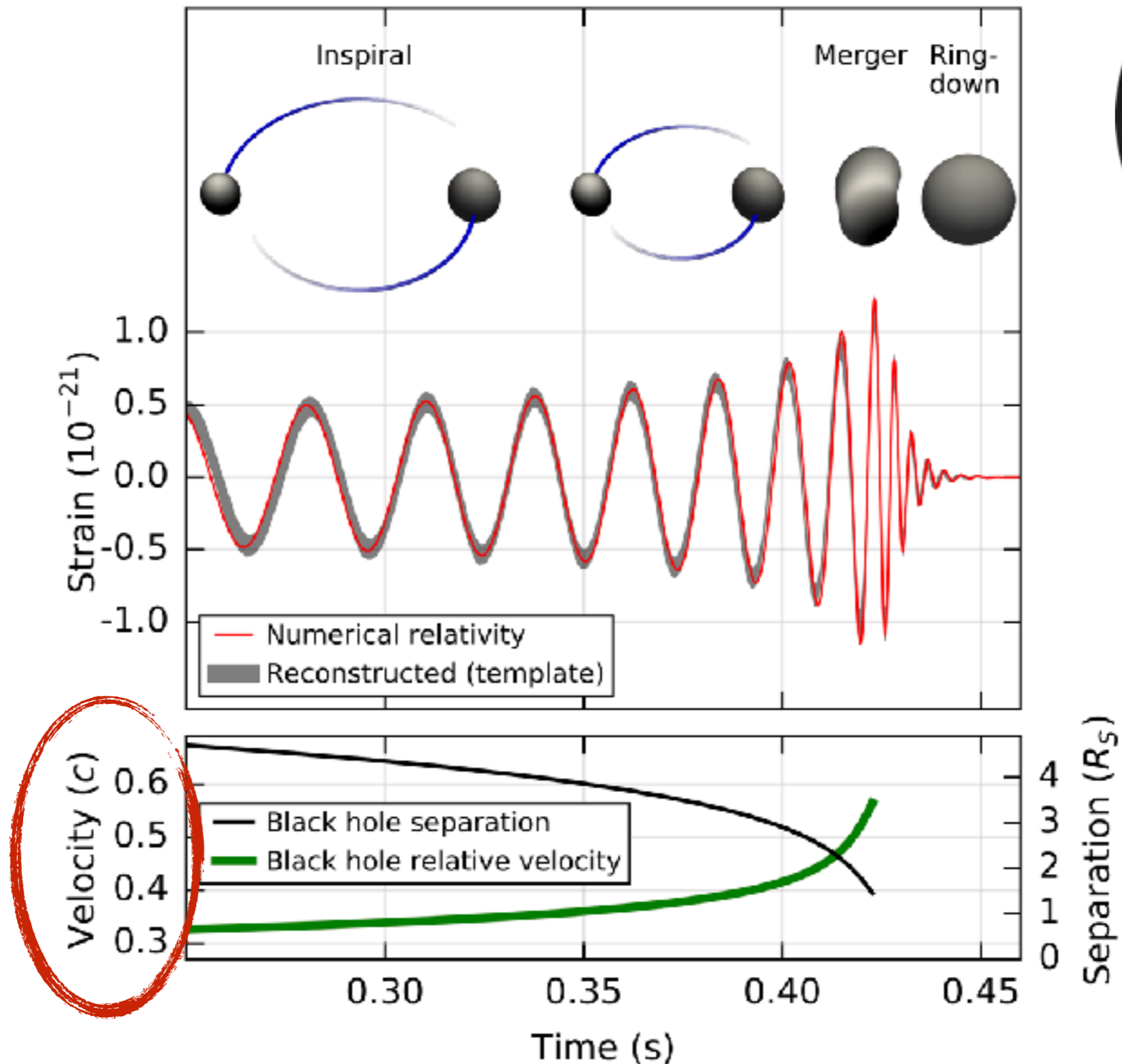
Gravitational waves are *new* cosmic messengers



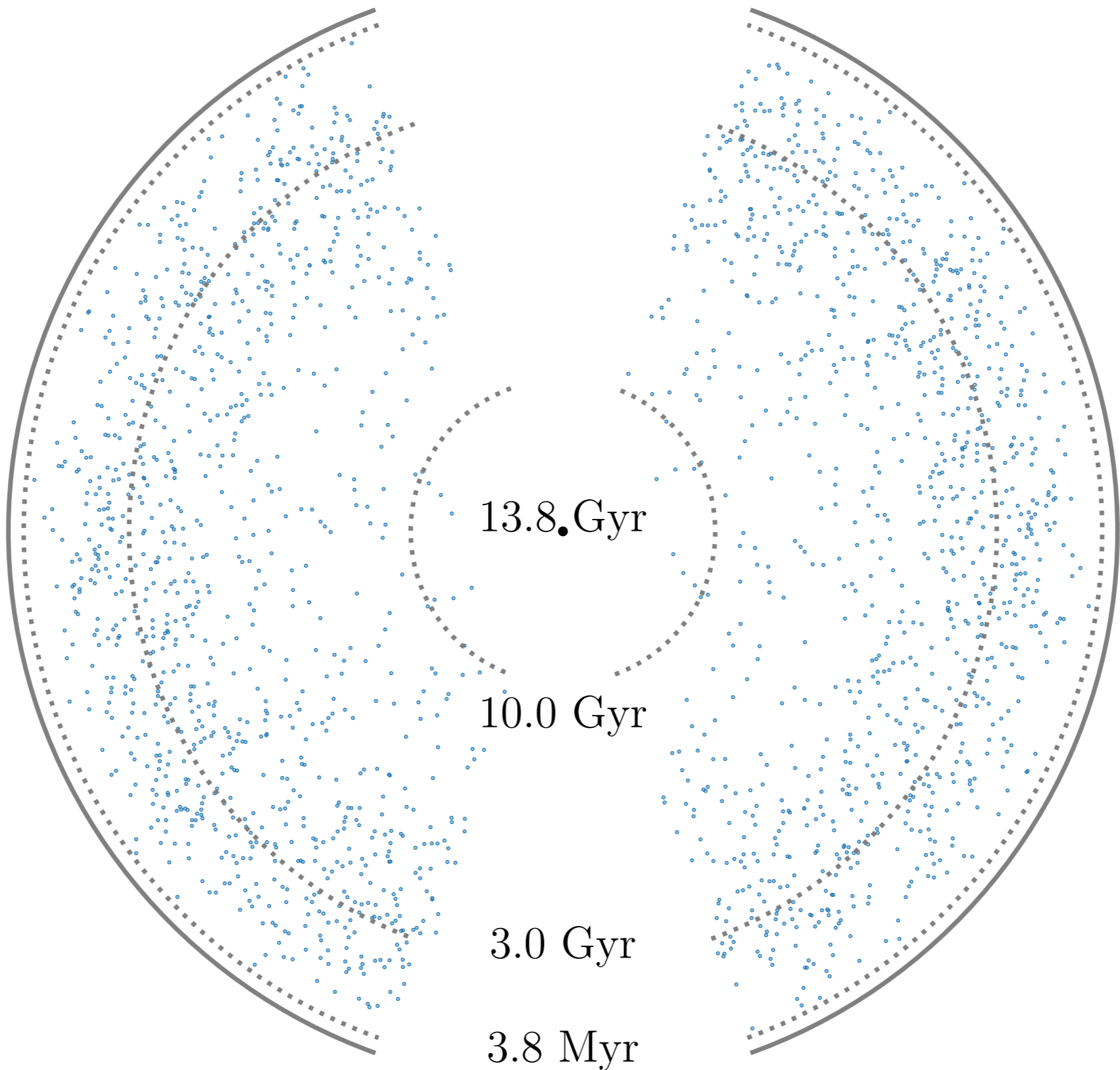
Gravitational waves from stellar-mass **binary black holes**



Strong-field gravity



[First detection, GW150914]

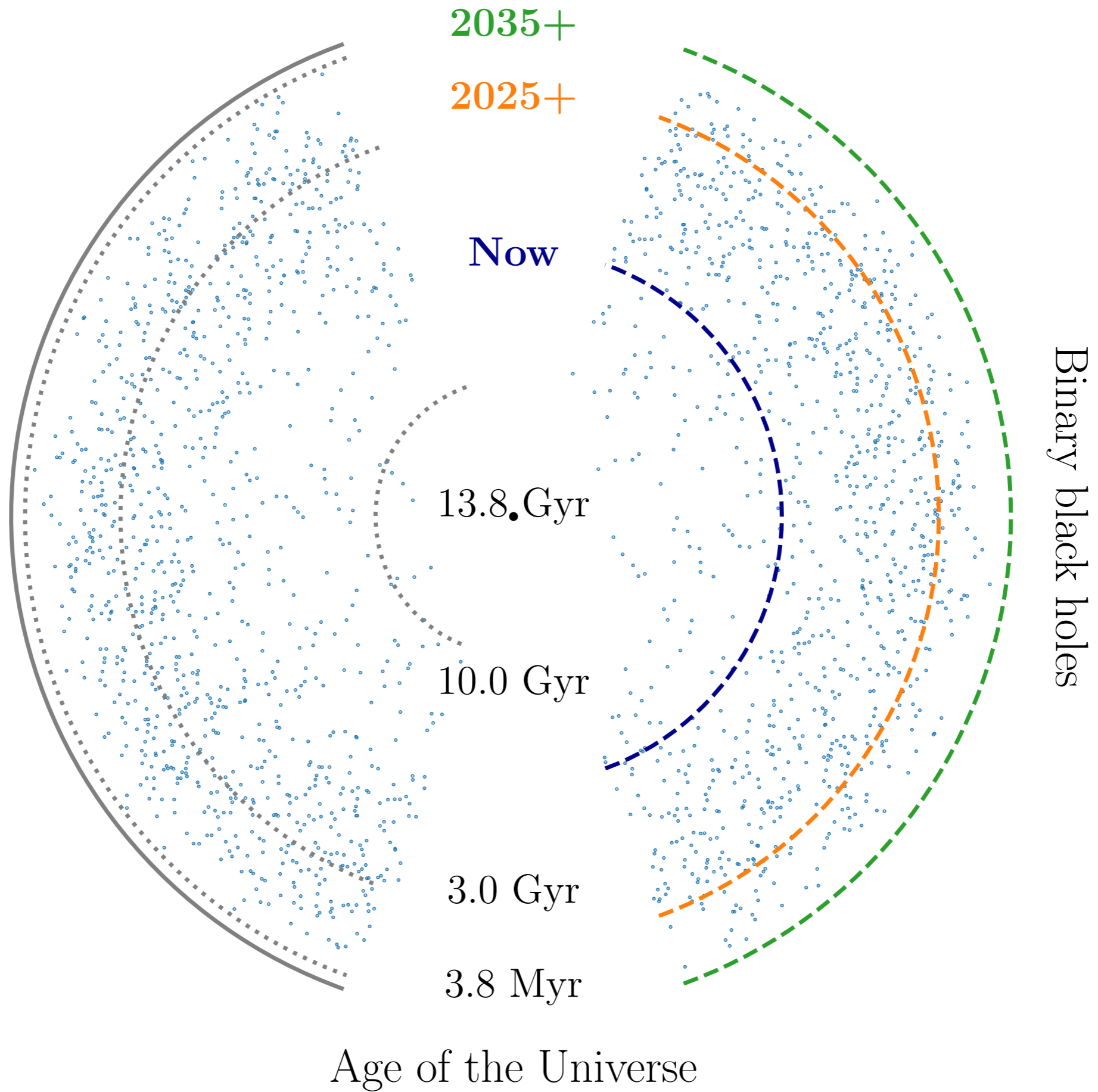


Binary black holes

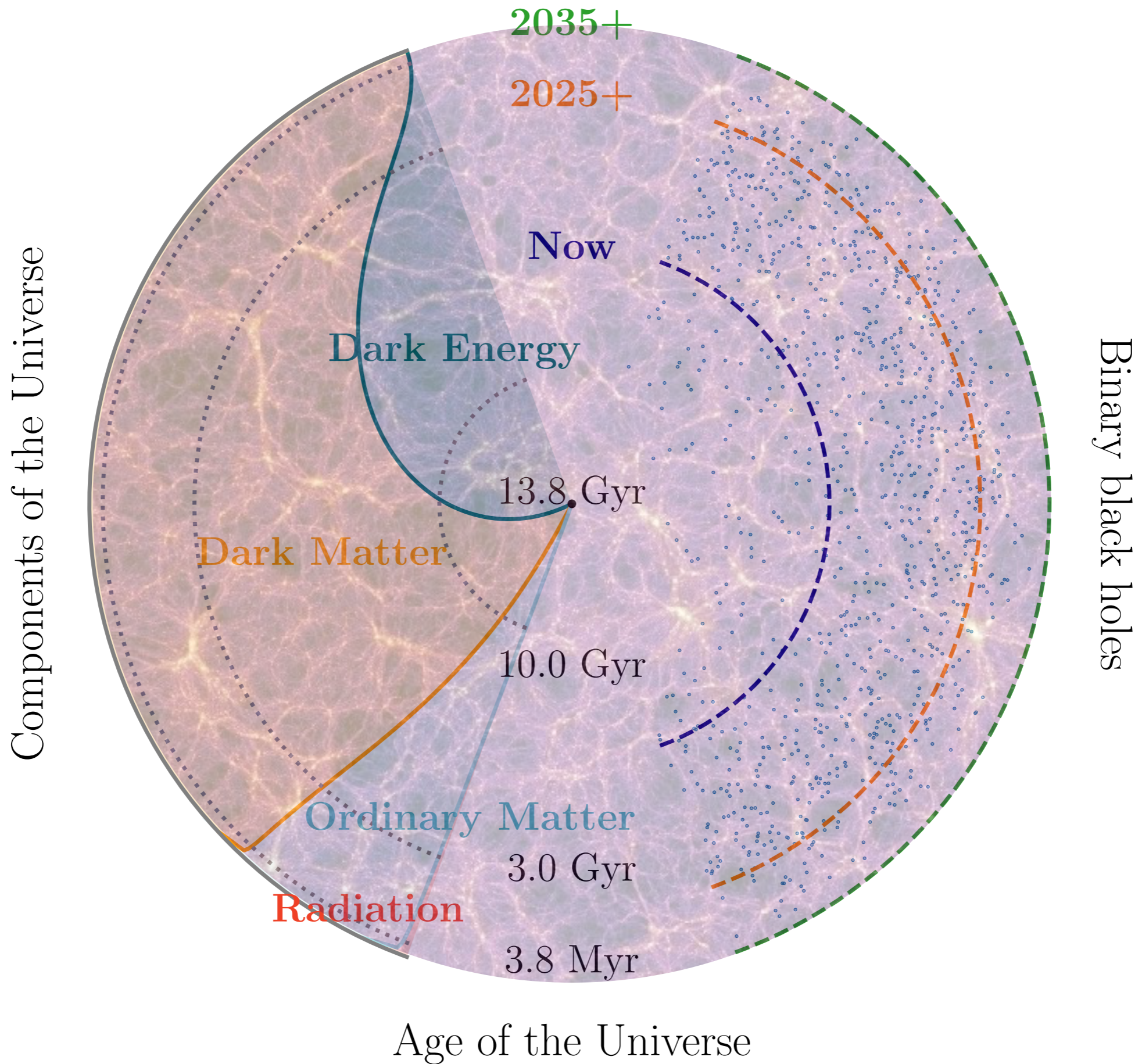
Age of the Universe

*stellar mass
binary black holes

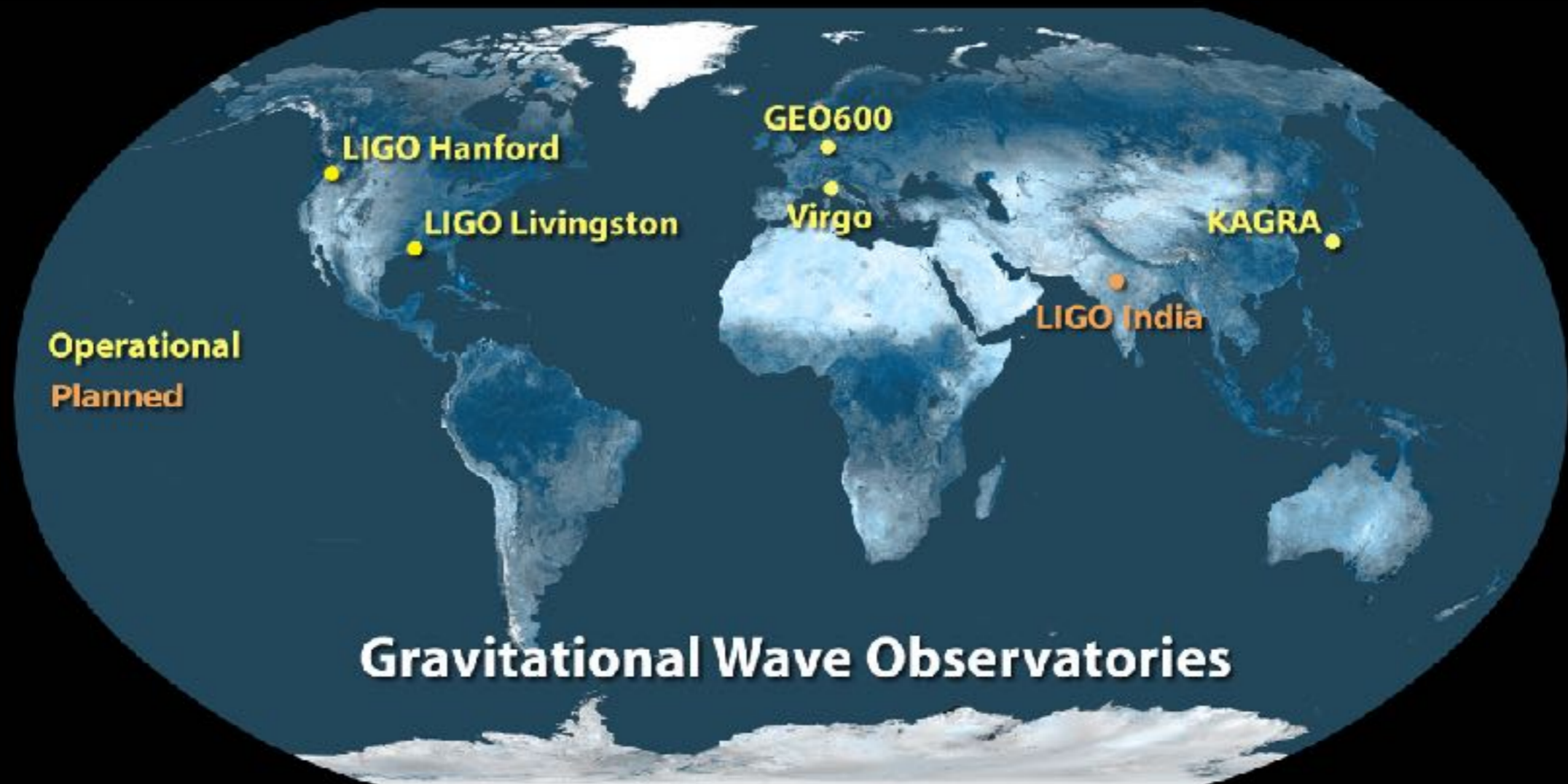
Gravitational Wave horizons



Gravitational Wave horizons



The era of gravitational wave astronomy is **here!**



[Hanford, US]



[Livingston, US]



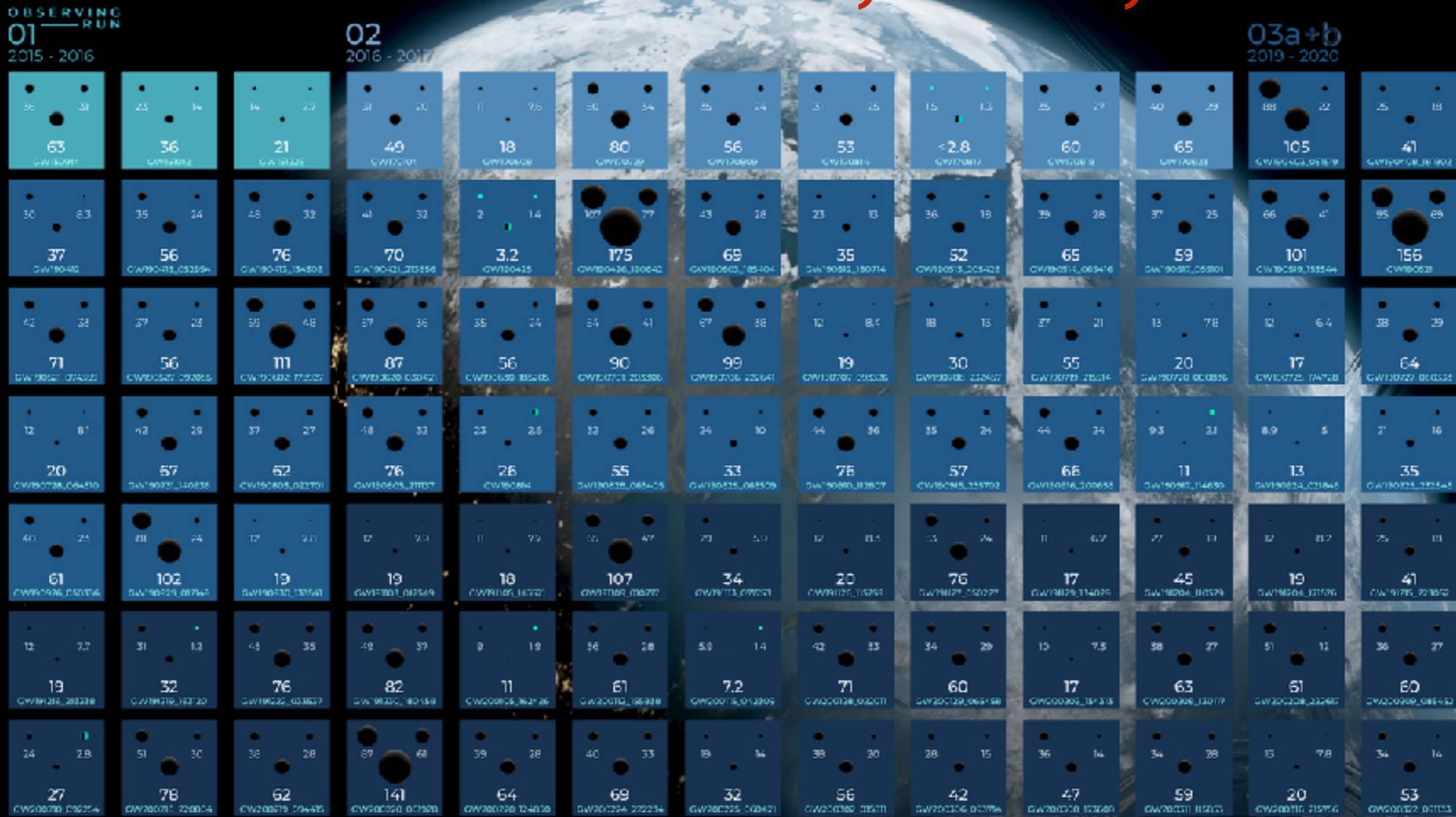
[Virgo, Italy]



[KAGRA, Japan]

The era of gravitational wave astronomy is here!

~100 events: BBH, BNS, NSBH

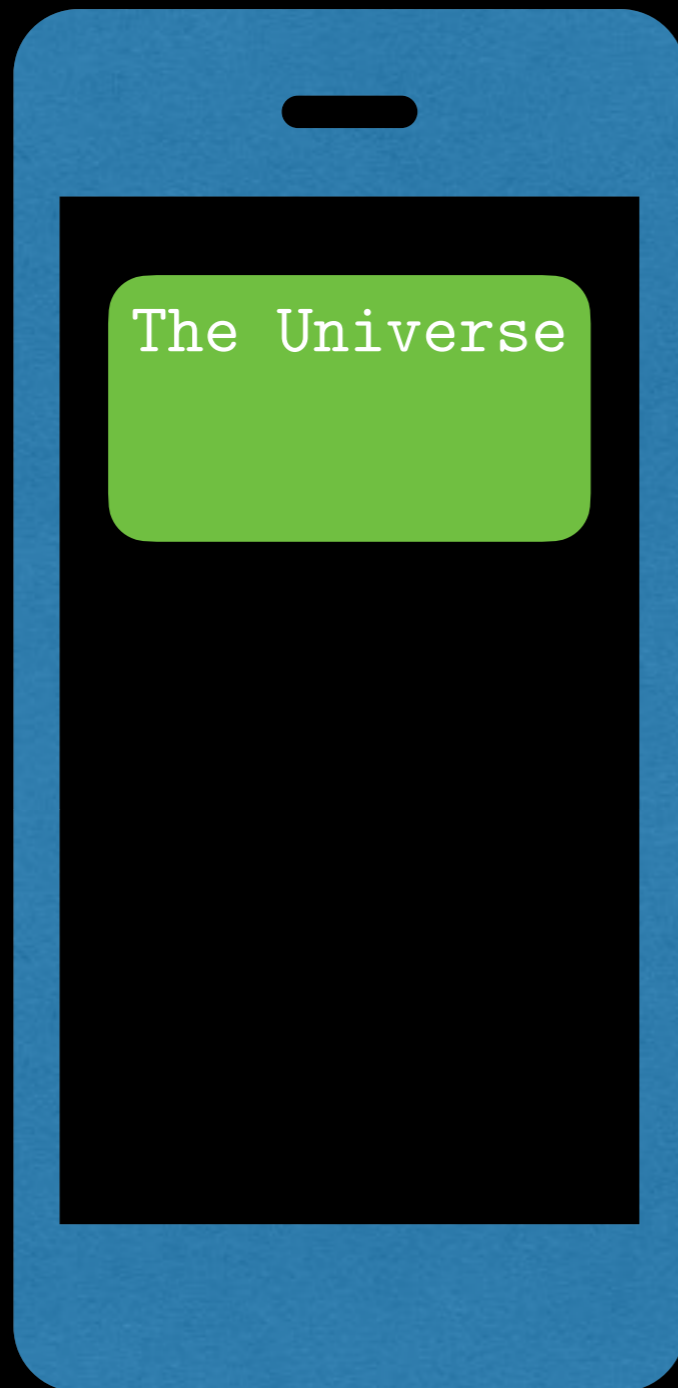


GRAVITATIONAL WAVE
MERGER
 DETECTIONS
 SINCE 2015



O4 is happening!

<https://gracedb.ligo.org/superevents/public/O4/#>



GraceDB Public Alerts ▾ Latest Search Documentation Login

Please log in to view full database contents.

LIGO/Virgo/KAGRA Public Alerts

- More details about public alerts are provided in the [LIGO/Virgo/KAGRA Alerts User Guide](#).
- Retractions are marked in **red**. Retraction means that the candidate was manually vetted and is no longer considered a candidate of interest.
- Less-significant events are marked in **grey**, and are not manually vetted. Consult the [LVK Alerts User Guide](#) for more information on significance in O4.
- Less-significant events are not shown by default. Press "Show All Public Events" to show significant and less-significant events.

O4 Significant Detection Candidates: 167 (186 Total - 19 Retracted)

O4 Low Significance Detection Candidates: 2839 (Total)

Show All Public Events

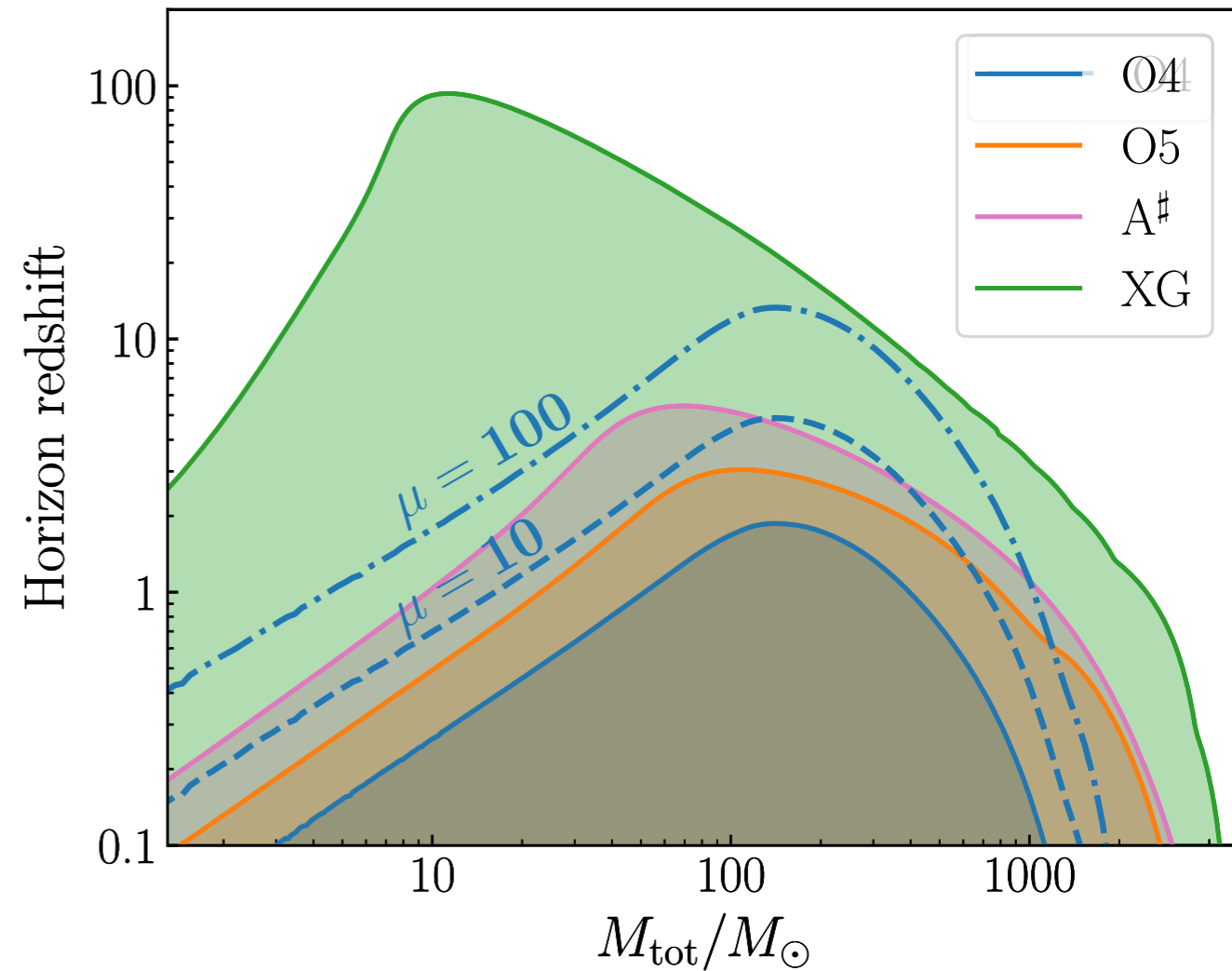
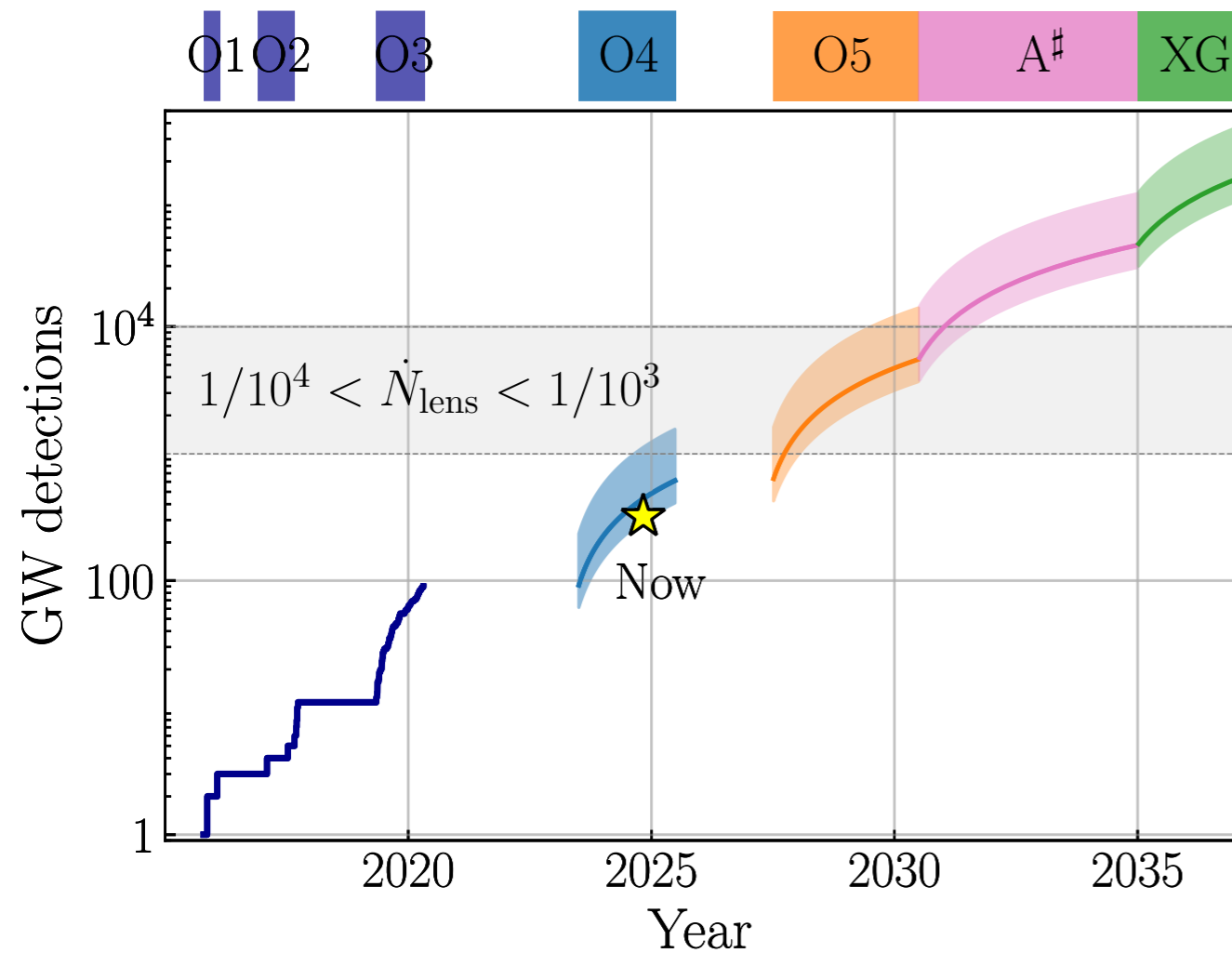
Page 1 of 1 new alerts

Event ID	Possible Source (Probability)	Significant	UTC	GCN	Location
S241201ac	BBH (97%), Terrestrial (3%)	Yes	Dec. 1, 2024 05:57:58 UTC	GCN Circular Query Notices VDE	

+160 candidates

Gravitational wave lensing:

First detection *approaching*, expanding *horizons*



The plan

*[this is an overview.
Ask the experts in the
room!]*



[ezquiaga.github.io/slides/
ezquiaga_vienna_24.pdf]

0. Motivation: gravity, astrophysics, cosmology

1. Gravitational waves are Standard Sirens

Waveforms from first principles, understood selection function

2. A crash-course on gravitational lensing

The diffraction integral, stationary phase approximation, repeated gravitational wave chirps

3. Current searches

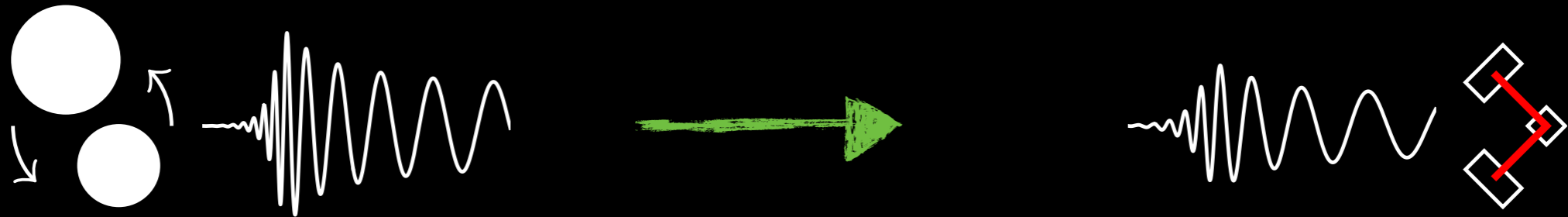
Multiple chirps, distorted waveforms, type II events, highly magnified gravitational waves

4. Future prospects

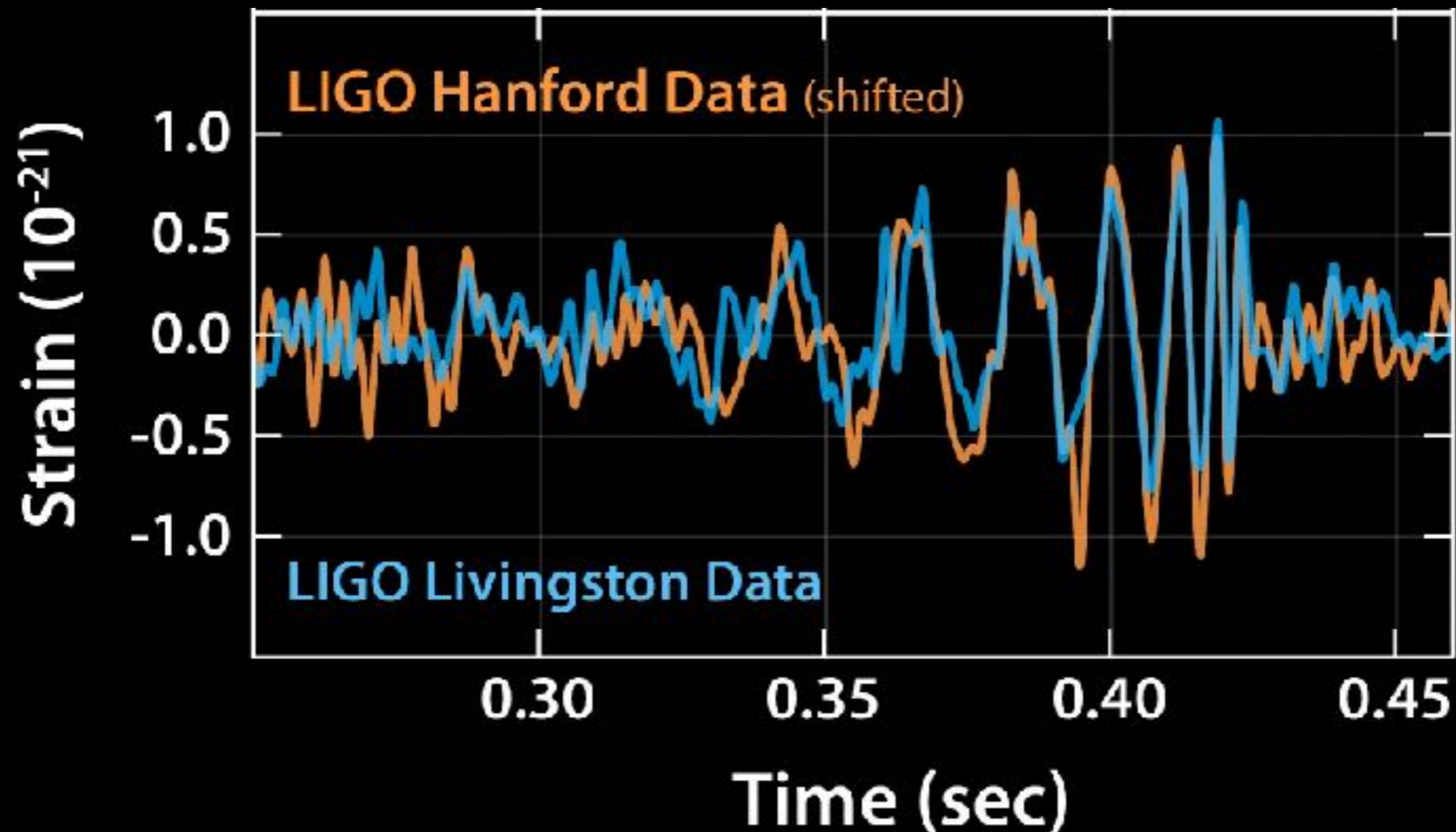
Substructure, multi-messenger & wave optics, source & lens populations, false violations of general relativity

1. Gravitational waves are
standard sirens

Gravitational waves are **standard sirens**



[general relativity predicts waveform]



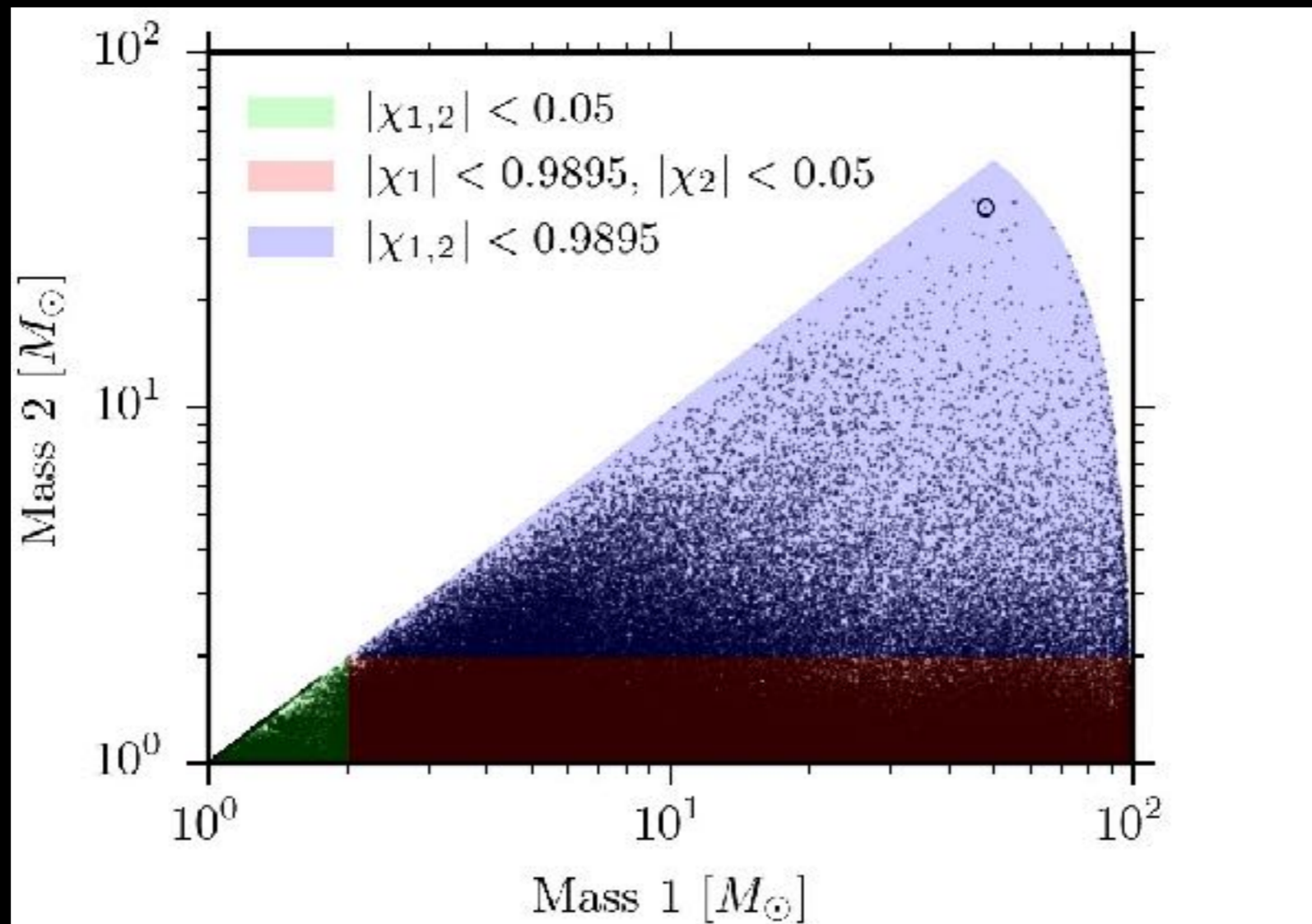
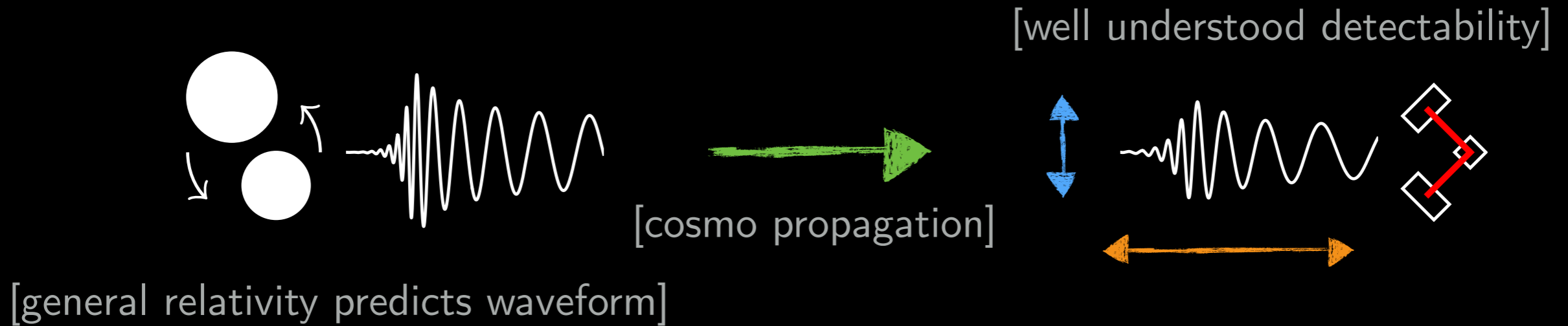
Gravitational waves are **standard sirens**



[general relativity predicts waveform]

$$h_c(t_{\text{obs}}) \sim \frac{\mathcal{M}_z^{5/3} f_{\text{obs}}^{2/3}}{d_L^{\text{gw}}}$$

Gravitational waves are **standard sirens**



Gravitational waves are **standard sirens**

[well understood detectability]



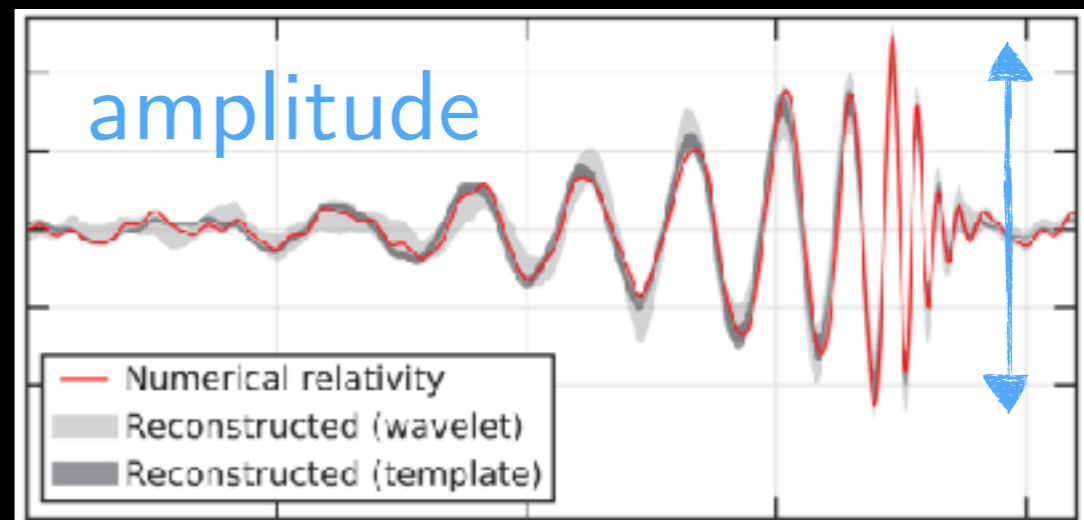
$$d_L(z)$$

[GW Hubble diagram]

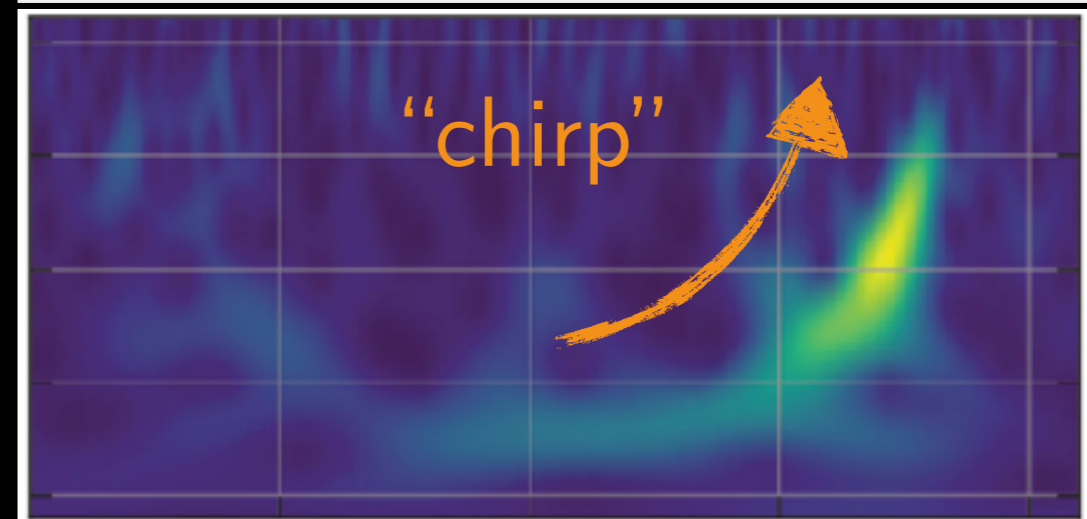
$$m_{\text{det}} = (1 + z)m$$

[Interplay with astrophysics]

strain

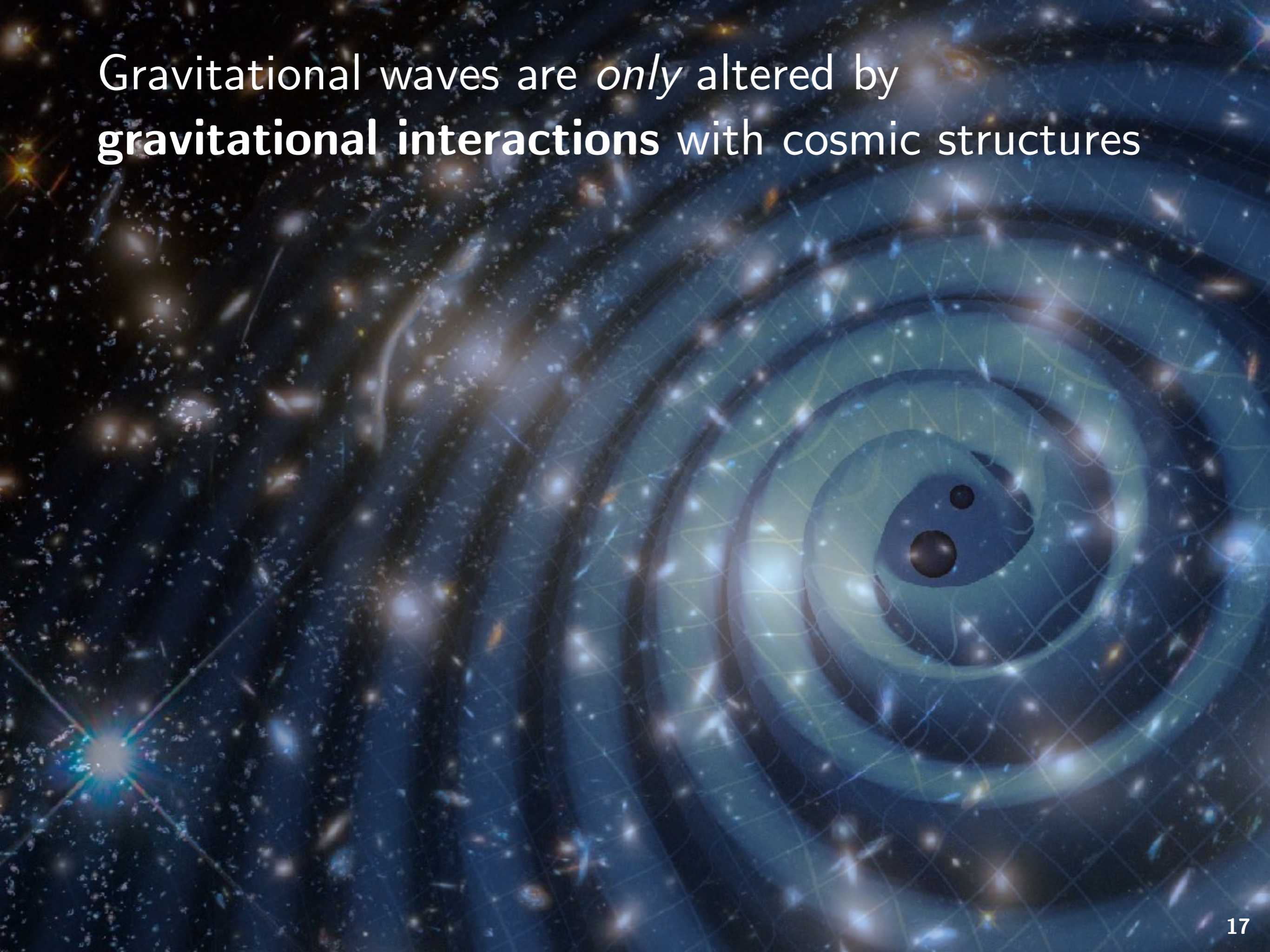


frequency



time

Gravitational waves are *only* altered by **gravitational interactions** with cosmic structures



2. A crash-course on gravitational lensing

Gravitational lensing - electromagnetic spectrum



[multiple images]



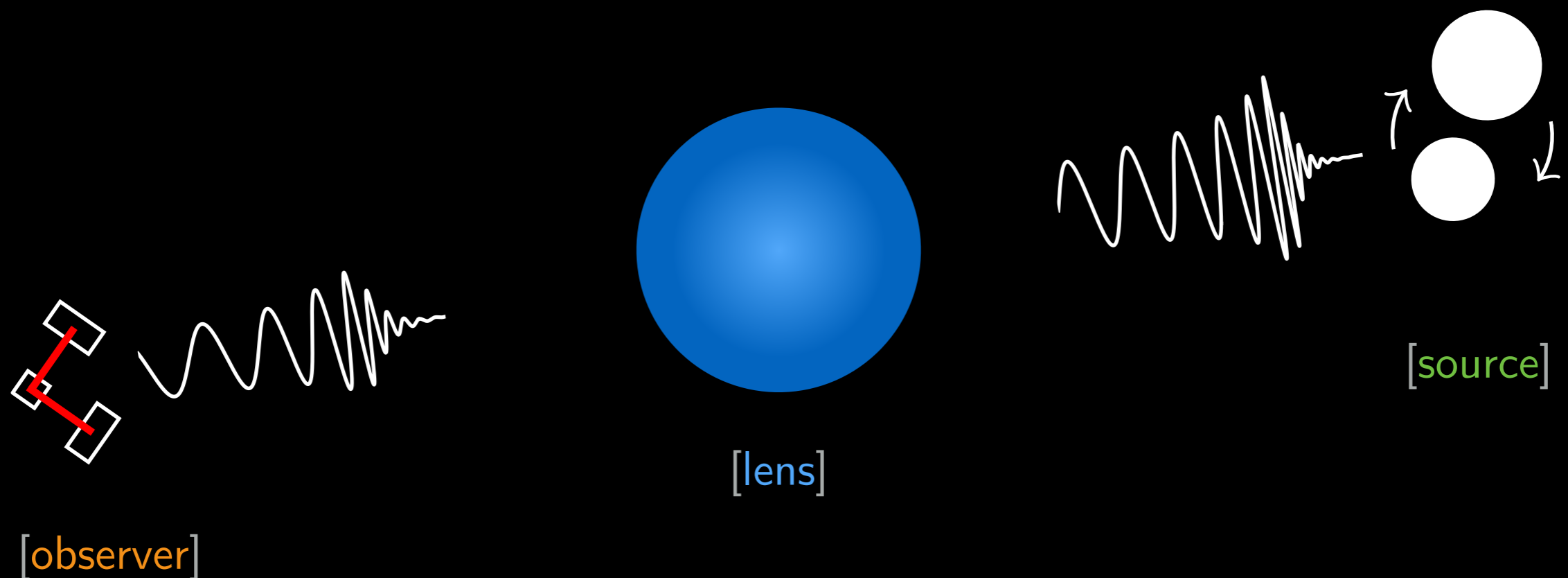
[arcs and rings]

Gravitational lensing

- Solve GW propagation on a **curved** background

$$\square \bar{h}_{\mu\nu} + 2\bar{R}_{\alpha\mu\beta\nu} \bar{h}^{\alpha\beta} = 0$$

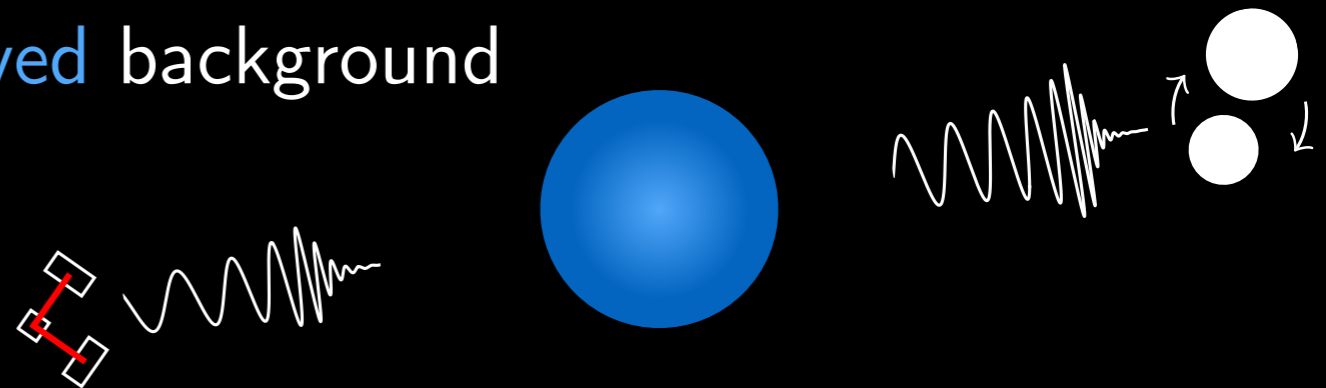
- We want to make a mapping between the **source** and the **observer** through the **lens**



Gravitational lensing

- Solve GW propagation on a **curved** background

$$\square \bar{h}_{\mu\nu} + 2\bar{R}_{\alpha\mu\beta\nu} \bar{h}^{\alpha\beta} = 0$$



- Cosmological background + **gravitational potential**

$$ds^2 = a(\eta)^2 \left(-(1 + 2\Phi)d\eta^2 + (1 - 2\Phi)d\vec{x}^2 \right)$$

- Focus on **weak-field** limit

$$\Phi \sim r_{\text{Sch}}/r \ll 1$$

- Equations simplify, same propagation for both **polarizations**

$$\nabla^2 h_A - (1 - 4\Phi)\partial_0^2 h_A = 0$$

[see Cusin's and Motohashi's talks for spin effects]

Gravitational lensing

- Within *weak-gravity*, solve in *Fourier* space:

$$(\nabla^2 + \omega^2) \tilde{h}_A = 4\Phi\omega^2 \tilde{h}_A \quad R_L \ll D_L, D_{LS}$$

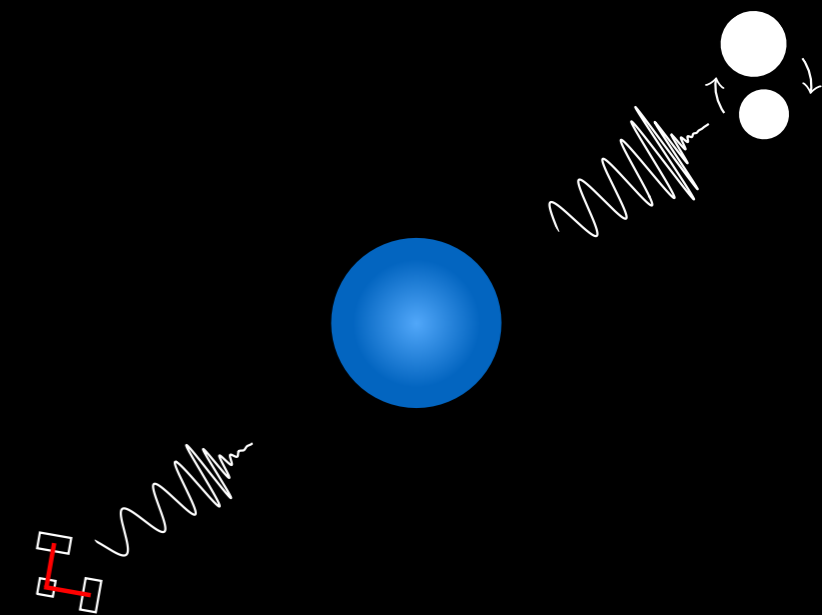
- For cosmological lenses, impose *thin lens* approximation.

- Integral solution: $h_L(\omega) = F(\omega, \theta_S) \cdot h(\omega)$

$$F(\omega, \vec{y}) = \frac{\omega}{2\pi i} \int d^2x \exp[i\omega T_d(\vec{x}, \vec{y})]$$

[Dimensionless variables] $\vec{x} \equiv \vec{\theta}/\theta_*$, $\vec{y} \equiv \vec{\theta}_S/\theta_*$, $\omega \equiv \tau_D \theta_*^2 \omega$

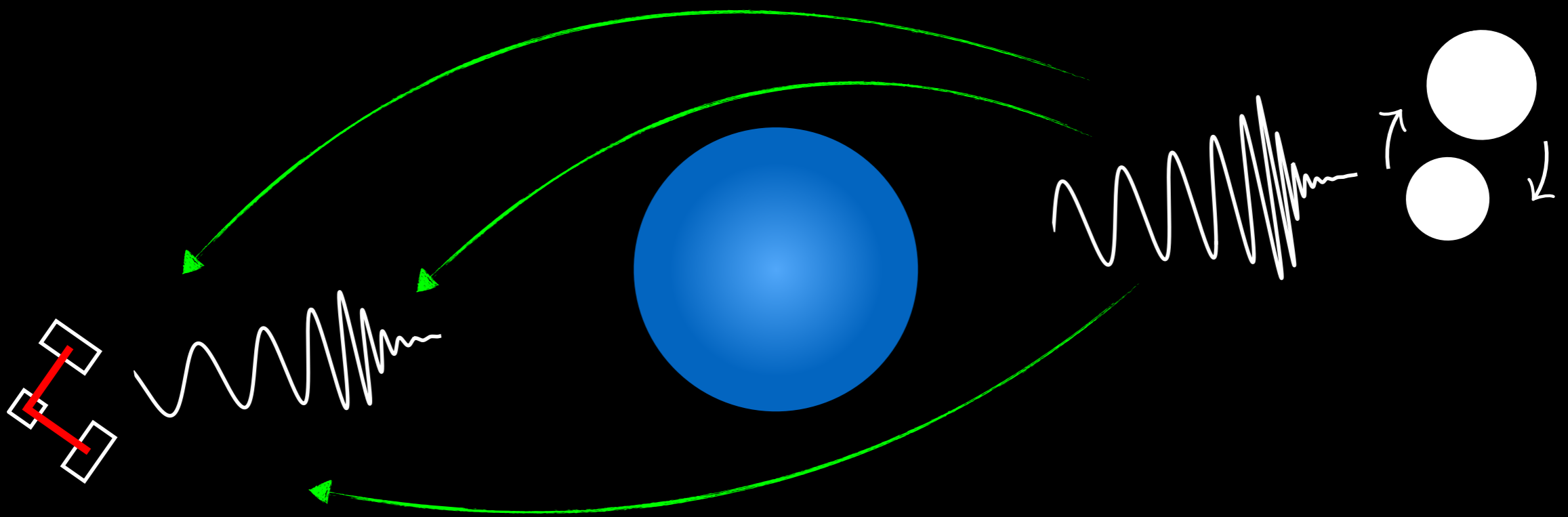
$$T_d \equiv t_d/\tau_D \theta_*^2 \quad \tau_D \equiv (1 + z_L) D_L D_S / c D_{LS}$$



Gravitational lensing

$$h_L(\omega) = F(\omega, \theta_S) \cdot h(\omega)$$

$$F(\omega, \vec{y}) = \frac{\omega}{2\pi i} \int d^2x \exp[i\omega T_d(\vec{x}, \vec{y})]$$



Stationary Phase Approximation

- Solve integral in the limit of *highly oscillatory* integrand

$$F(w, \vec{y}) = \frac{w}{2\pi i} \int d^2x \exp[iwT_d(\vec{x}, \vec{y})]$$

- Stationary points define the **images**:

$$\left. \frac{\partial t_d}{\partial \theta_a} \right|_{\vec{\theta}=\vec{\theta}_j} = 0$$

$$T_d(\vec{\theta}) \approx T_d(\vec{\theta}_j) + \frac{1}{2} \sum_{(a,b)=1}^2 \delta\theta_a \delta\theta_b \frac{\partial^2 T_d(\vec{\theta}_j)}{\partial \theta_a \partial \theta_b} + \dots$$

- Hessian matrix determines magnifications

$$\mu(\theta_j) = 1/\det(T_{ab}(\theta_j))$$

$$T_{ab} \equiv \tau_D^{-1} \partial^2 t_d / \partial \theta_a \partial \theta_b$$

Multiple chirps

$$\Delta t_d \cdot \omega \gg 1$$

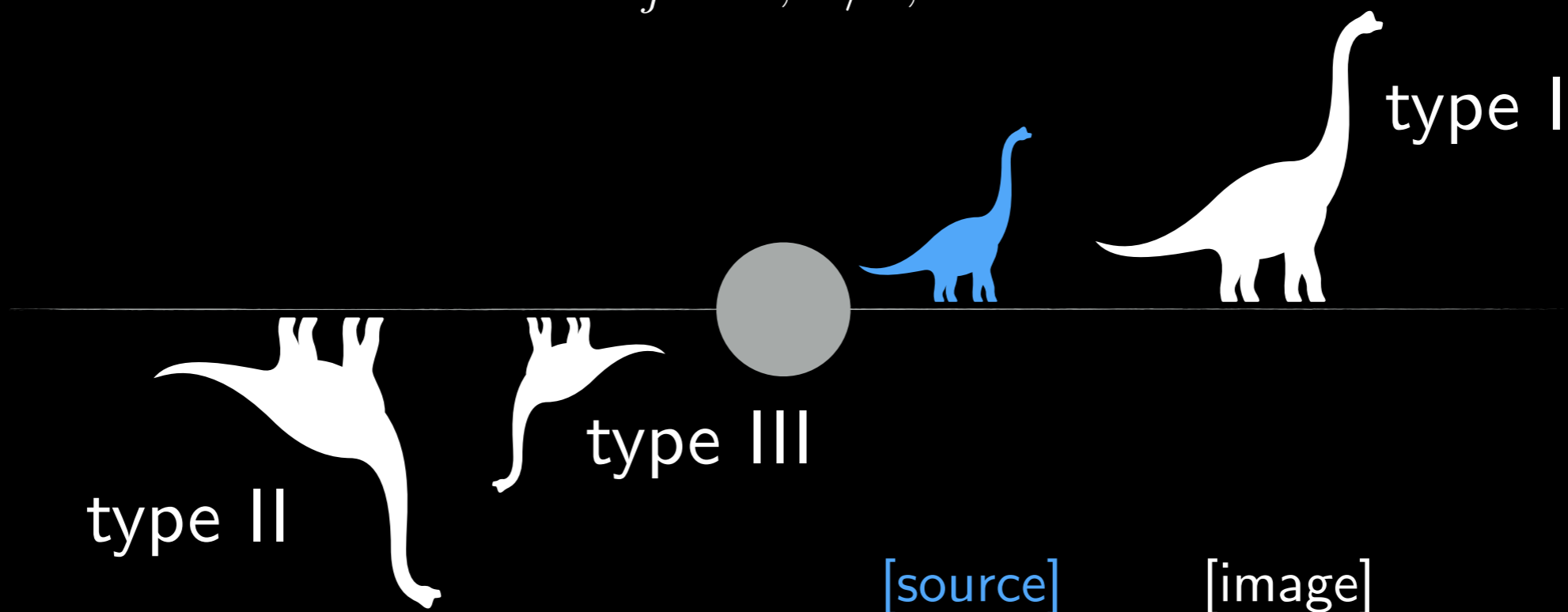
$$h_L(\omega) = F(\omega, \theta_S) \cdot h(\omega)$$

$$F \approx \sum_j |\mu_j|^{1/2} \exp(i\omega t_j - i\pi n_j)$$

Magnification
Time delay
Phase shift

- Lensed signals acquire a different phase shift

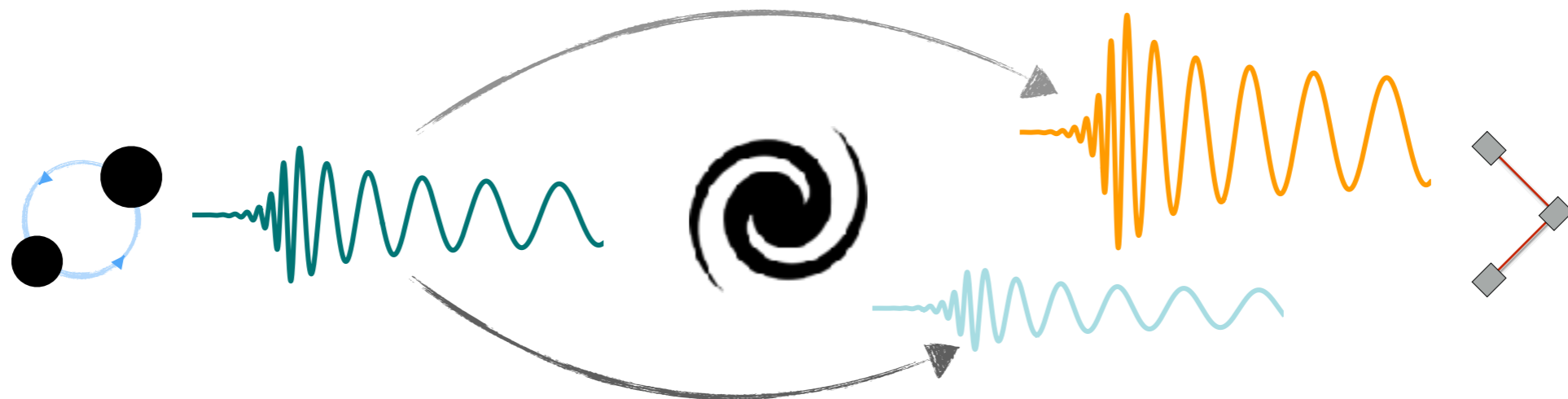
$$n_j = 0, 1/2, 1$$



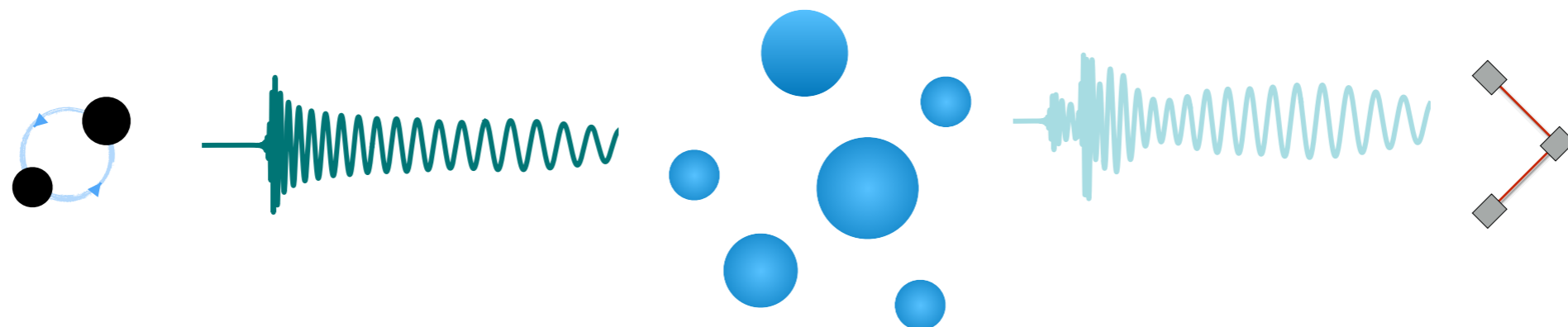
3. Current searches

Gravitational lensing - gravitational wave spectrum

Repeated chirps due to strong lensing



Waveform distortions by substructures

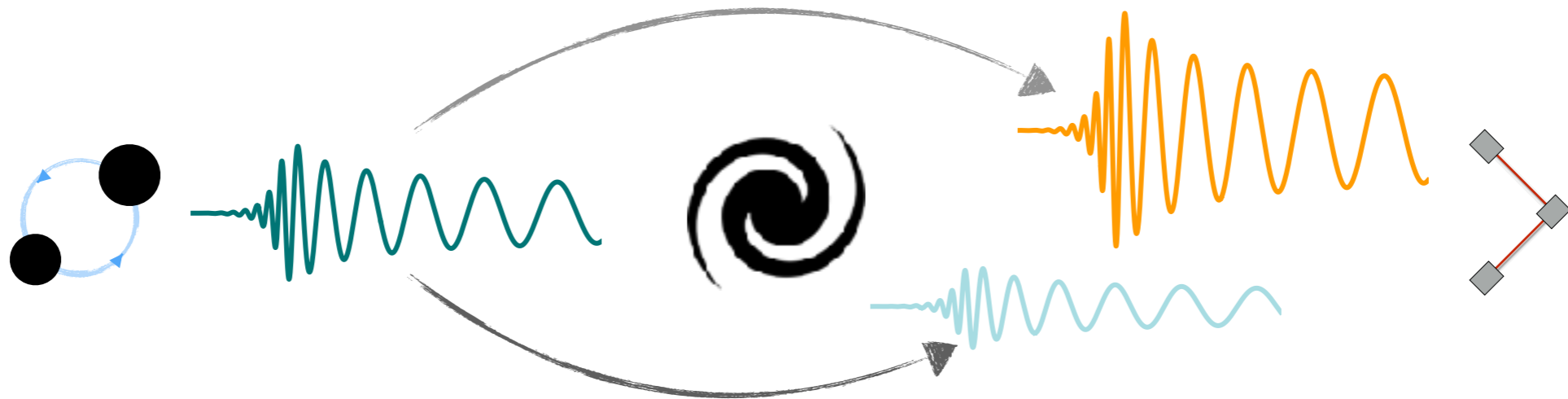


Source

Lens

Detector

Repeated chirps due to strong lensing



- The properties of the j -th chirp

$$d_L^j = d_L / \sqrt{|\mu_j|}$$

$$t_{\text{ref}}^j = t_{\text{ref}} + \Delta t_j$$

$$m_{\text{det}}^j = m_{\text{det}}$$

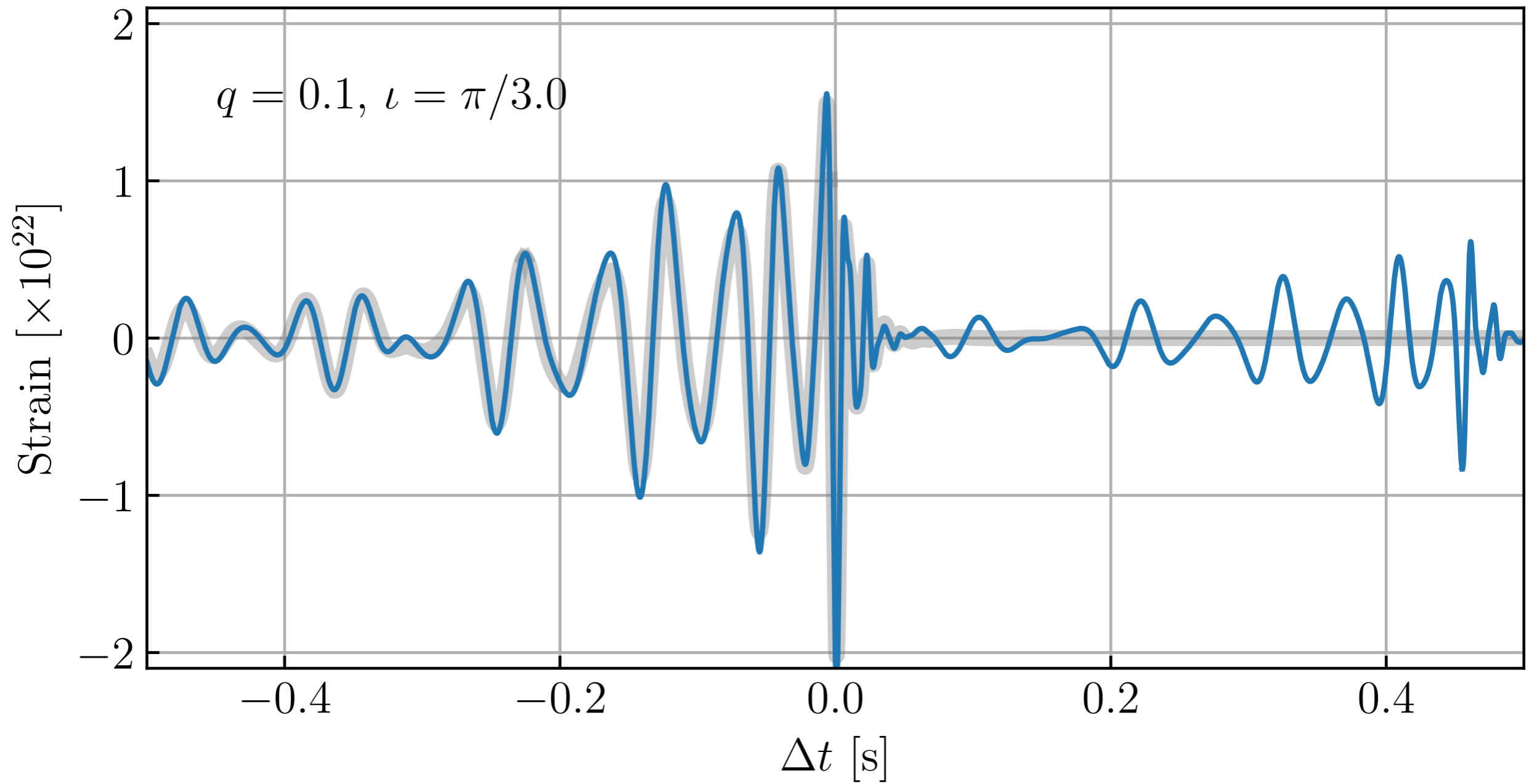
$$\phi_{\text{ref}}^j = \phi_{\text{ref}} - \pi/2$$

- If not identified as lensed, a *magnified* events appears *closer* and *more massive*

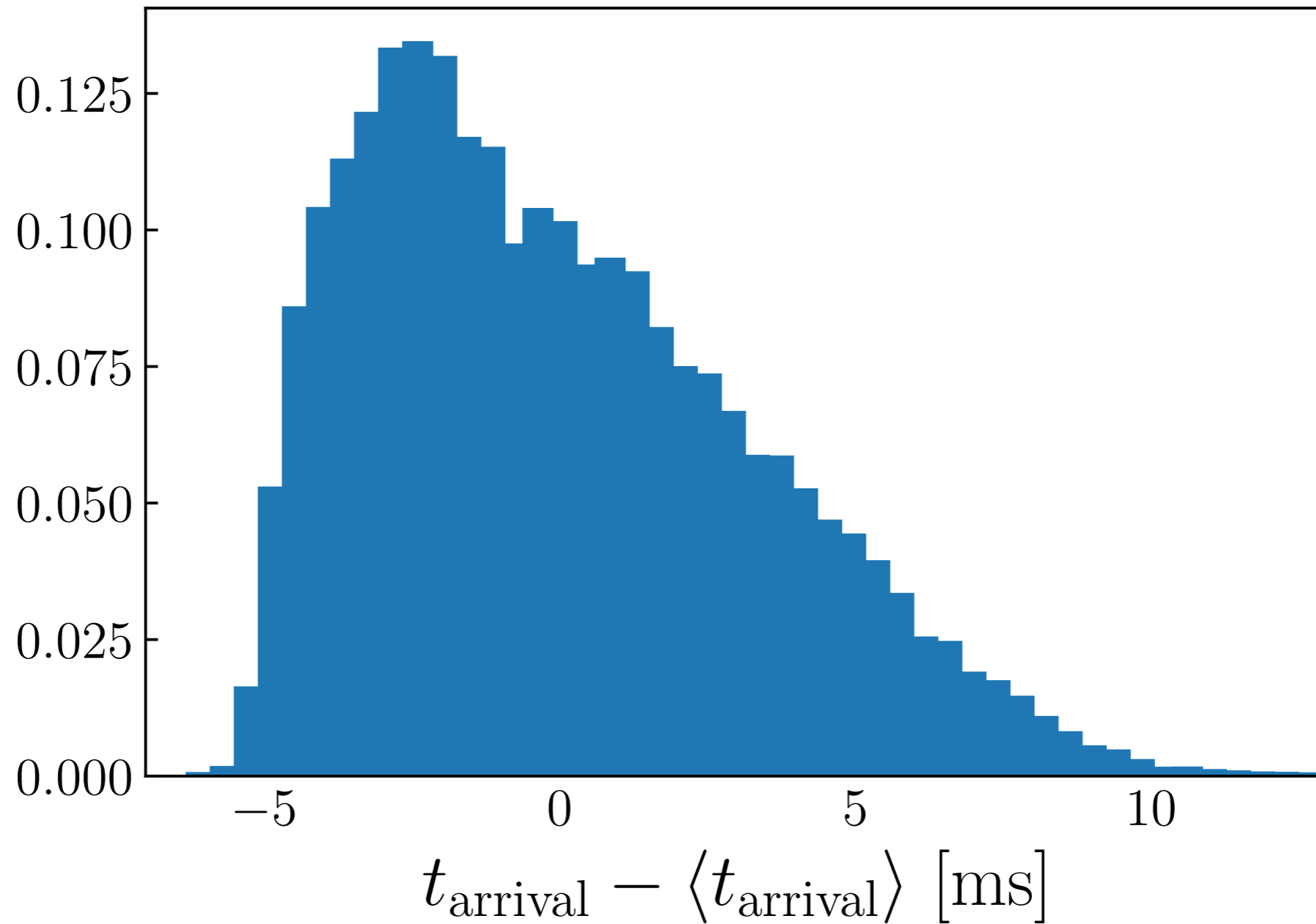
$$m_{\text{src}}^j = m_{\text{det}} / (1 + z(d_L^j))$$

[see Chen's talk]

Repeated, *coherent* signals

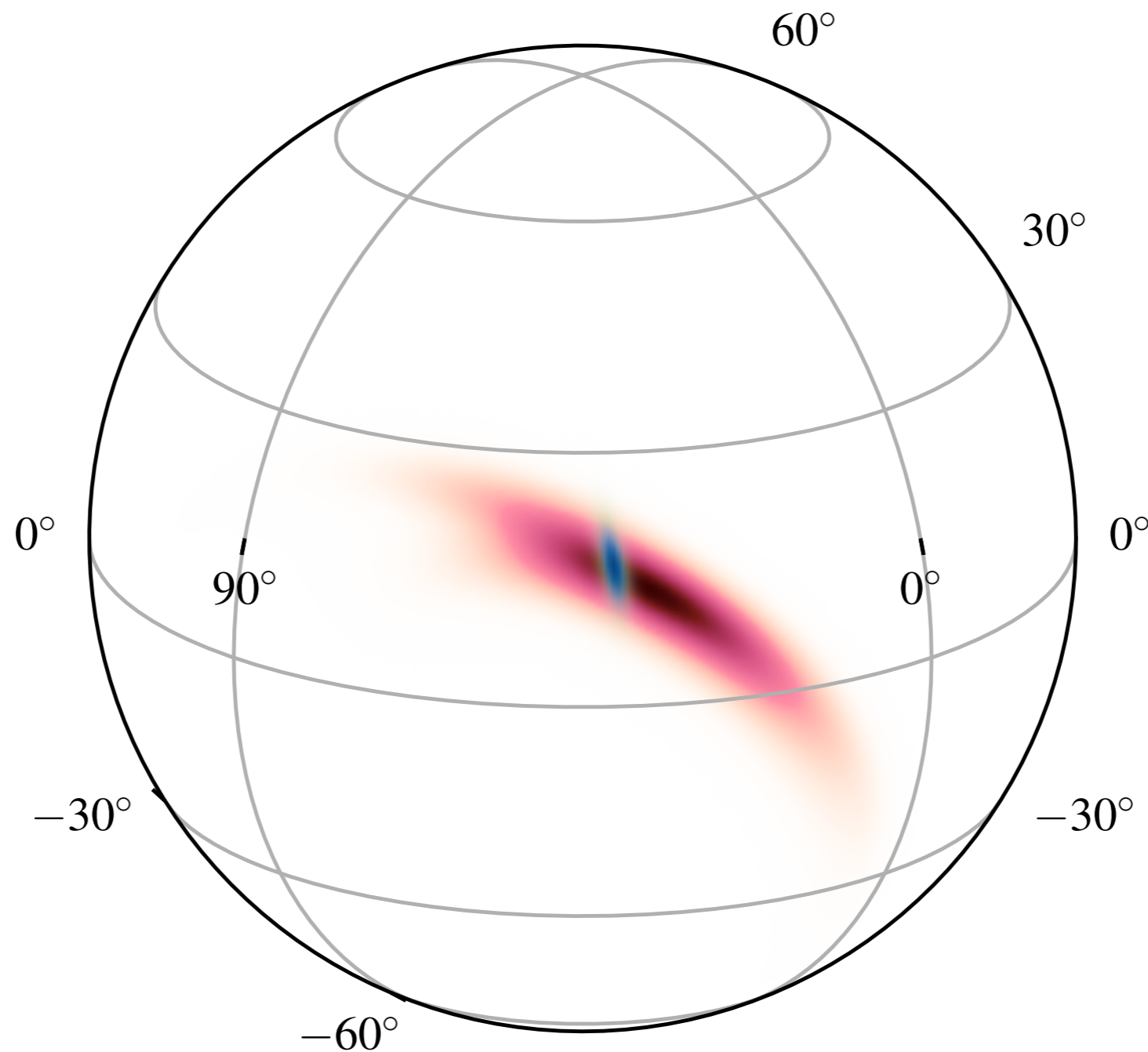


Precise timing

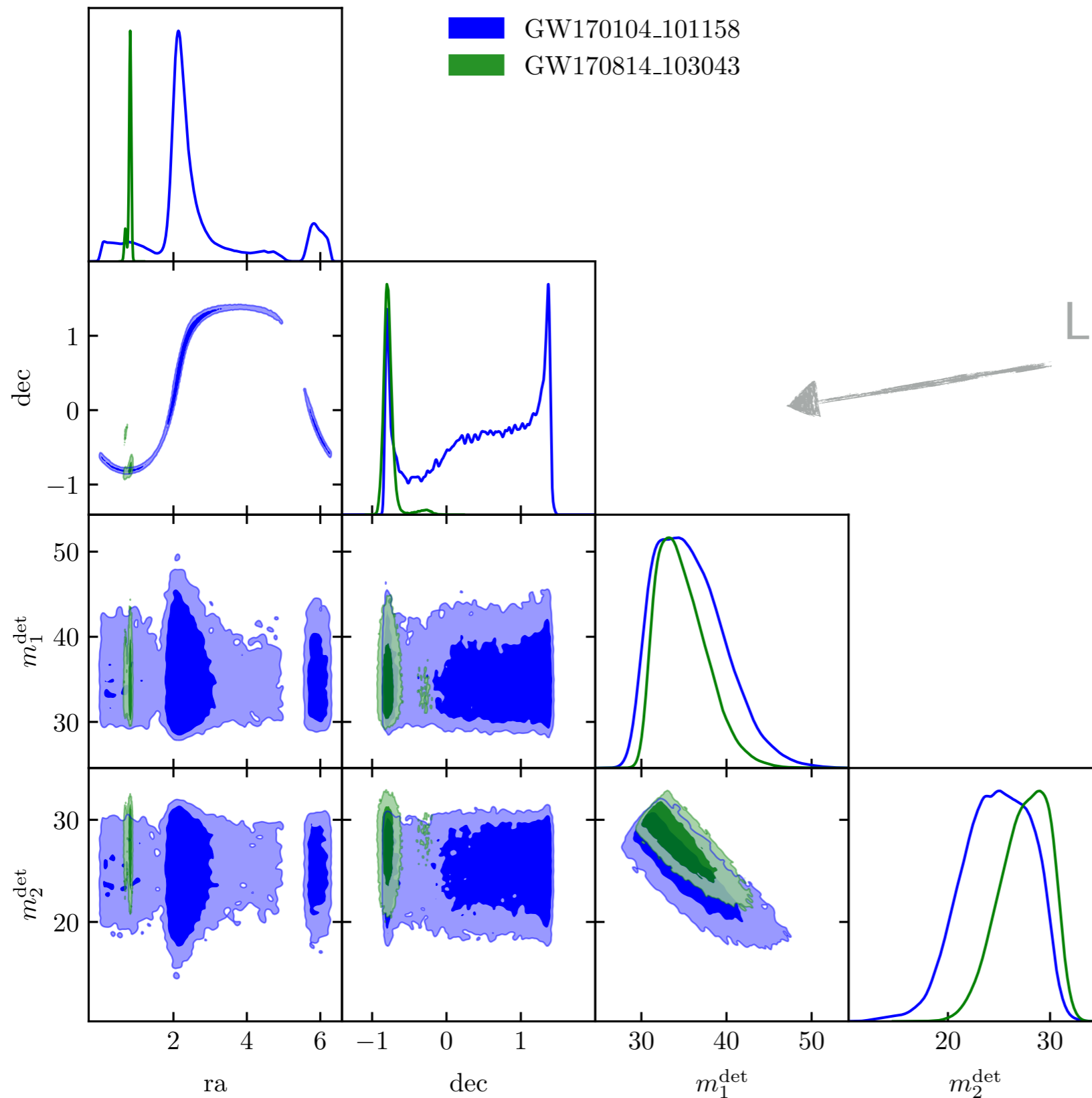


Poor sky localization

$$\theta_E \sim 1'' \sqrt{\frac{M}{10^{12} M_\odot}} \sqrt{\frac{1 \text{ Gpc}}{D}}$$

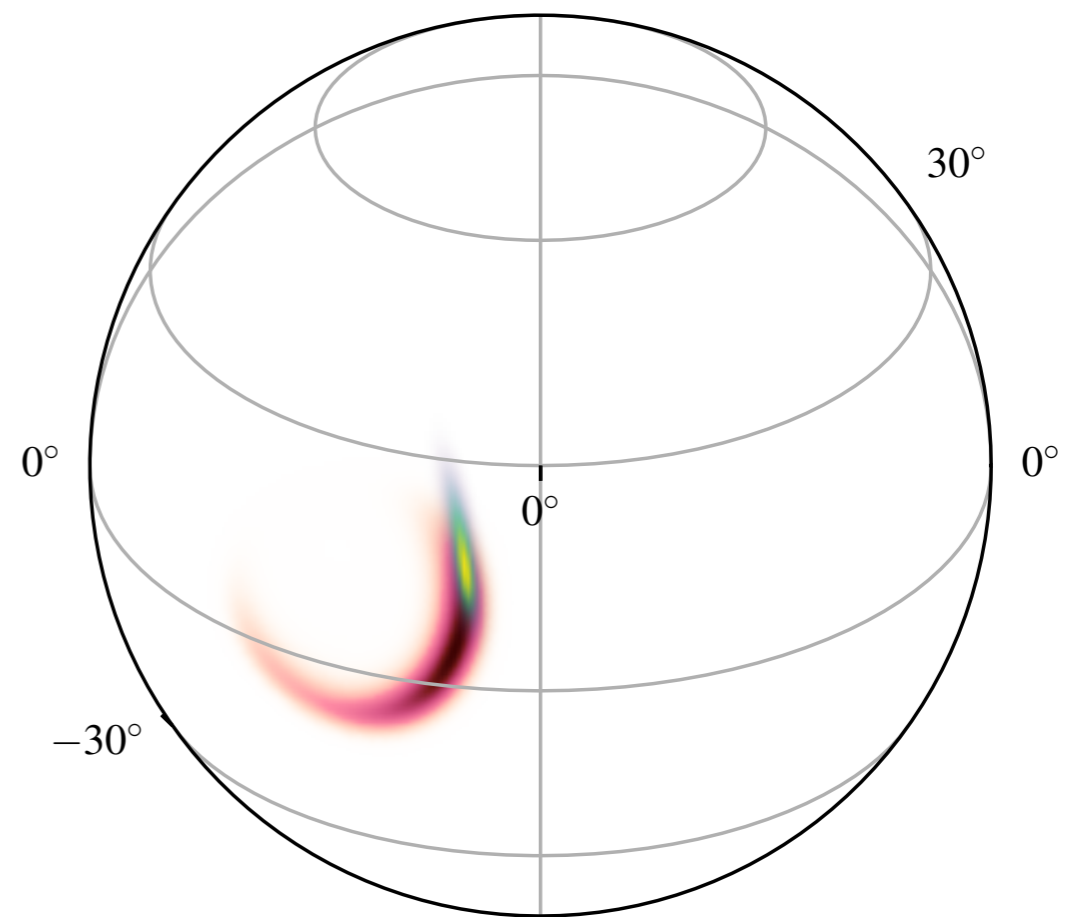
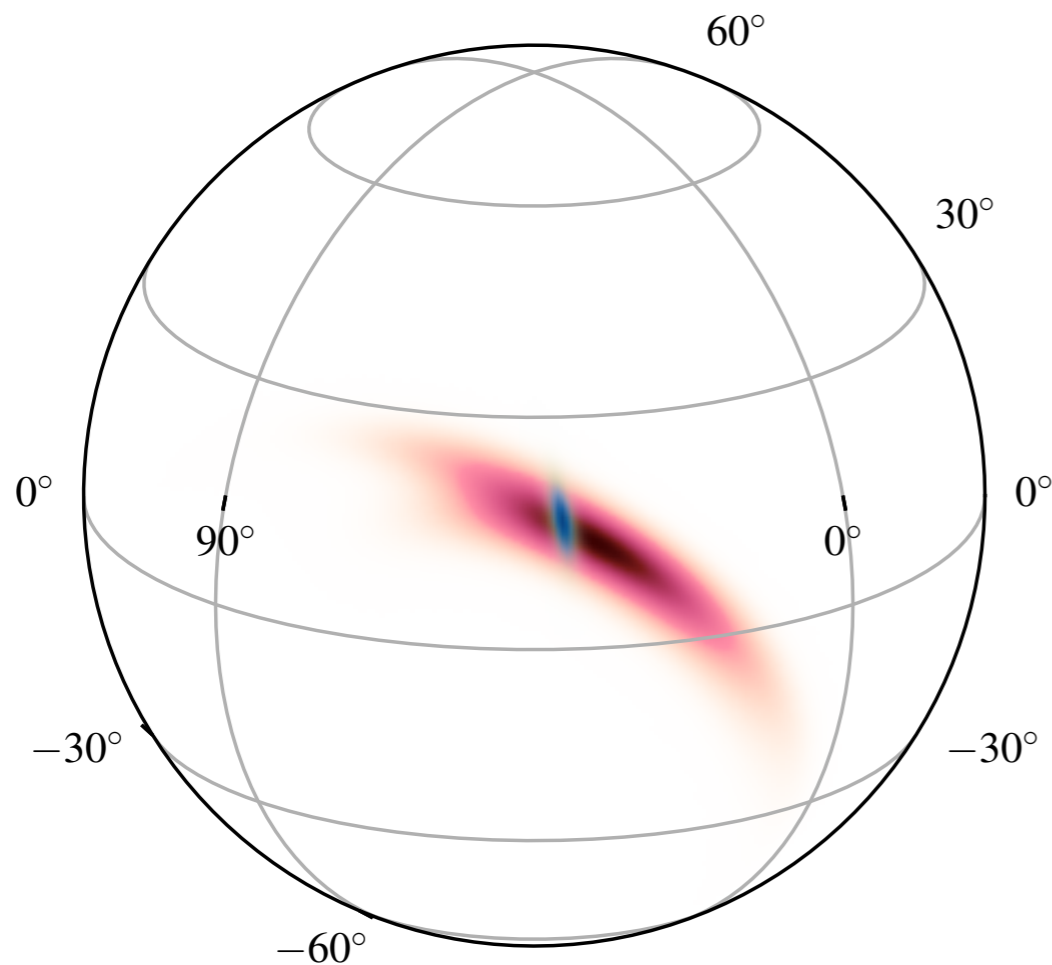
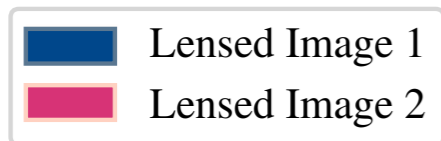


Searching for repeated chirps



Look for events with similar properties: masses, sky positions, spins...

Searching for repeated chirps: false alarms



$$N_{\text{false alarm}} \sim N^2$$

Searching for repeated chirps

- Given the large number of pairs, need quick methods to identify promising pairs

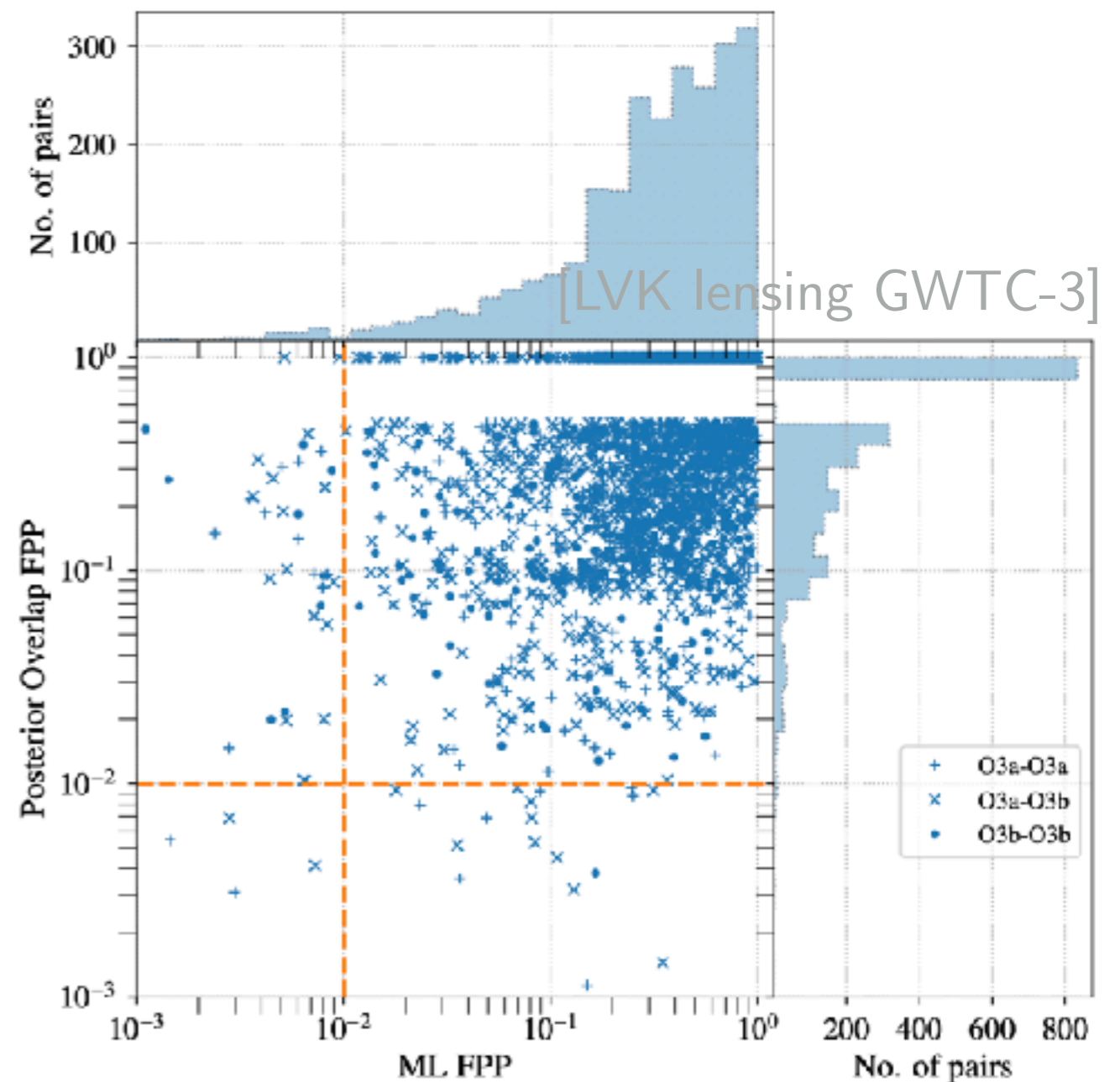
- Compute the posterior overlap

[Haris et al.; 2018]

- Use machine learning (ML) summary statistic

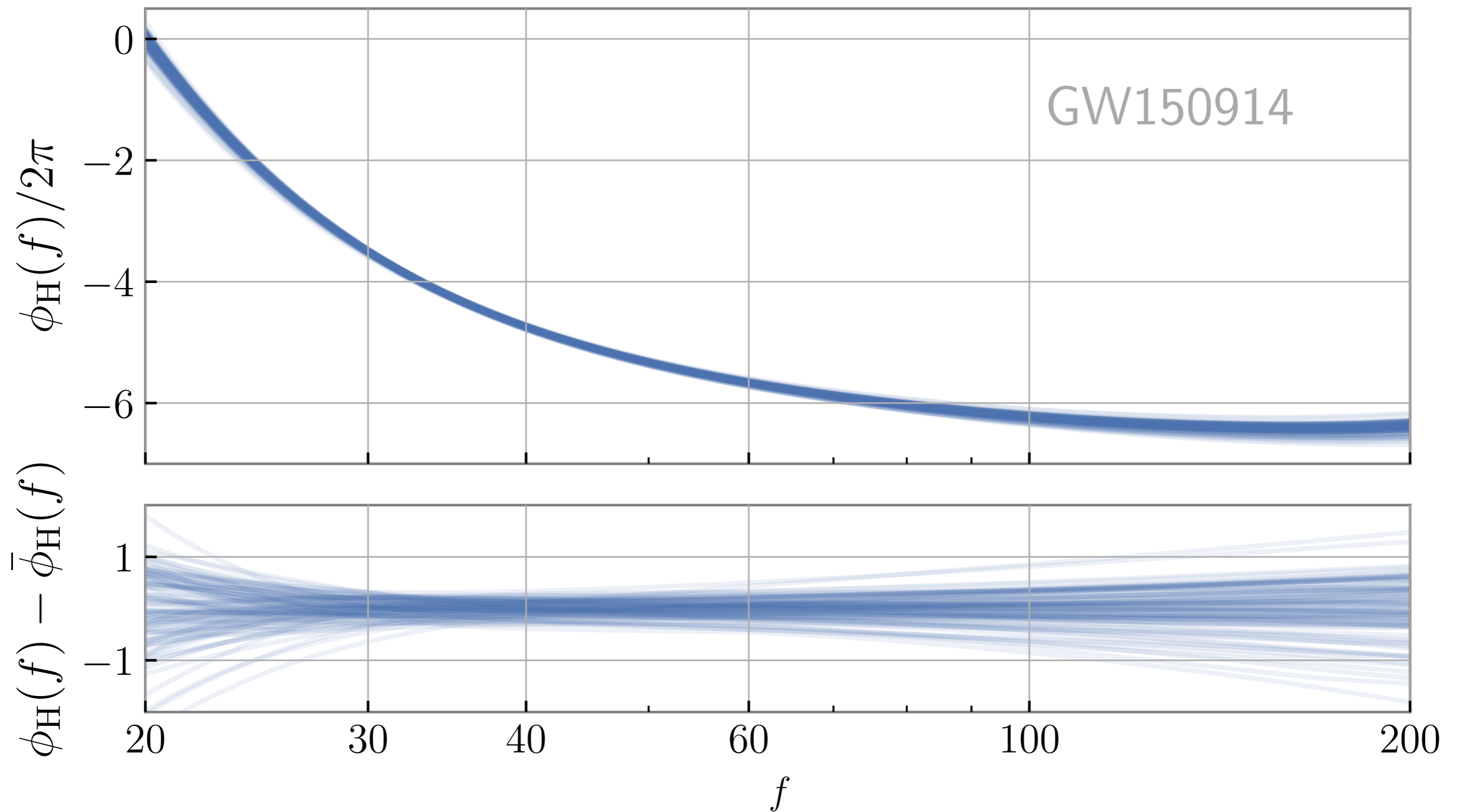
[Goyal et al.; 2021]

[see Li's talk]

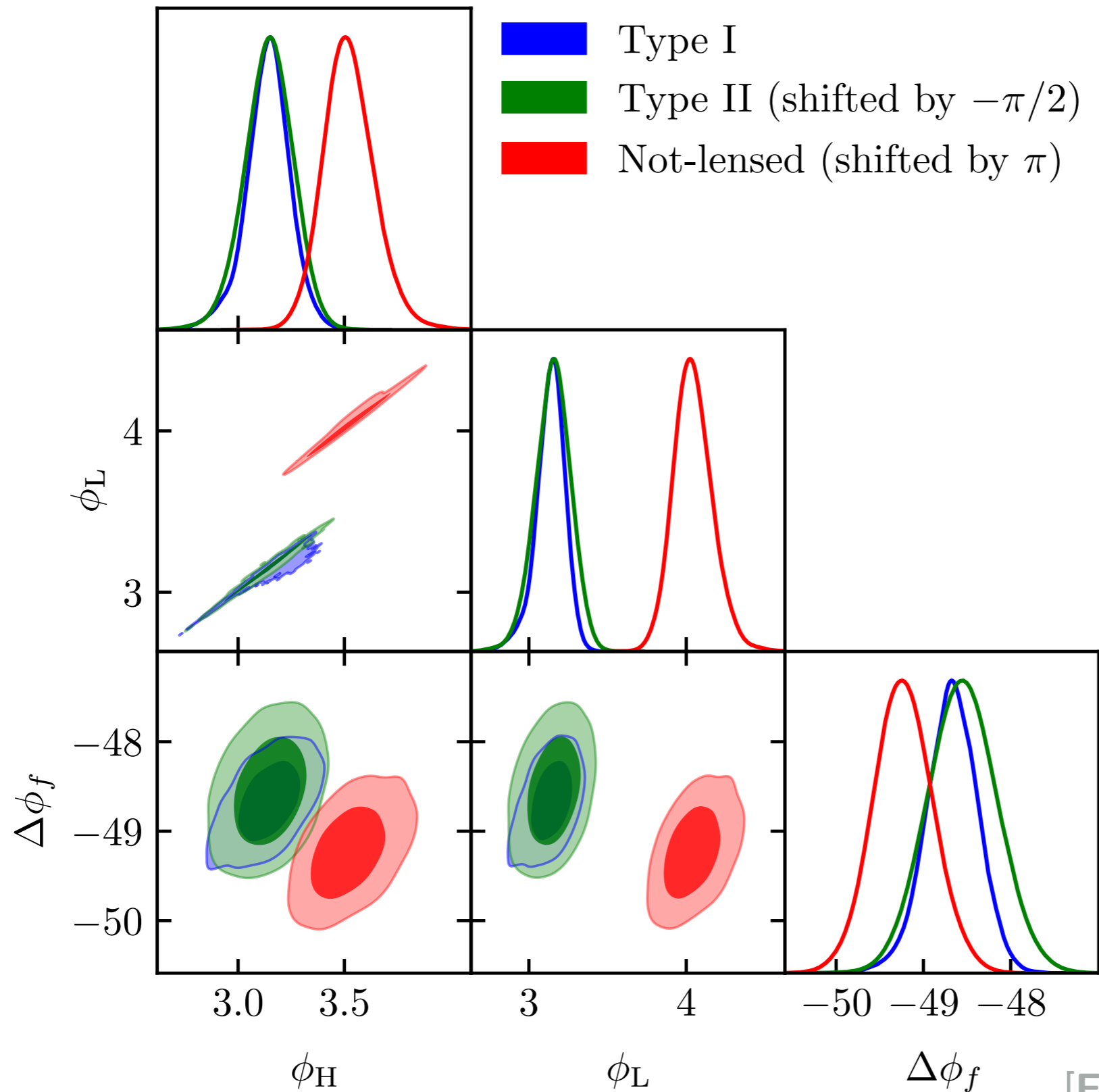


FPP = false-positive probability

Fight false alarms: **phase consistency**

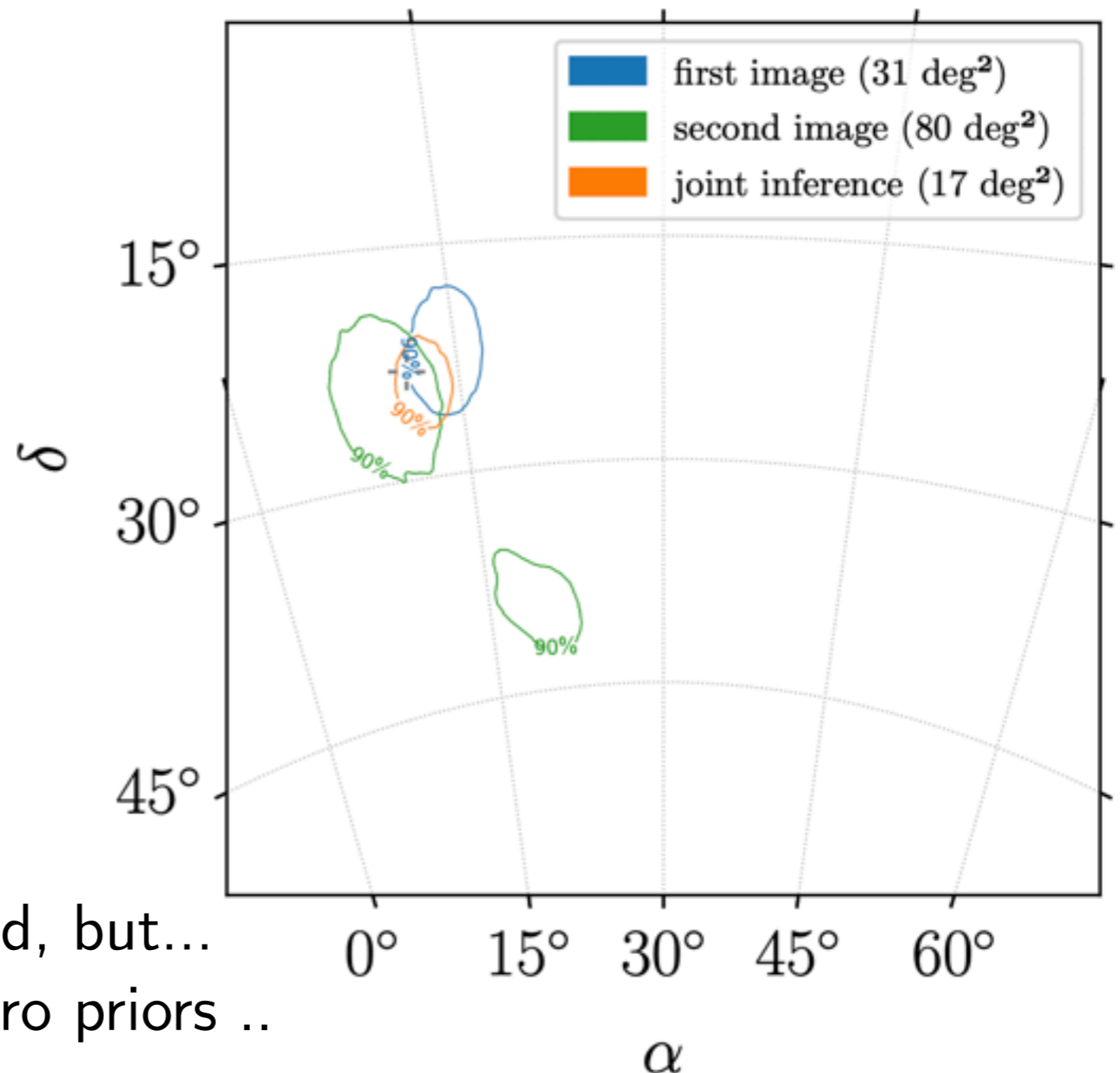


Fight false alarms: **phase consistency**



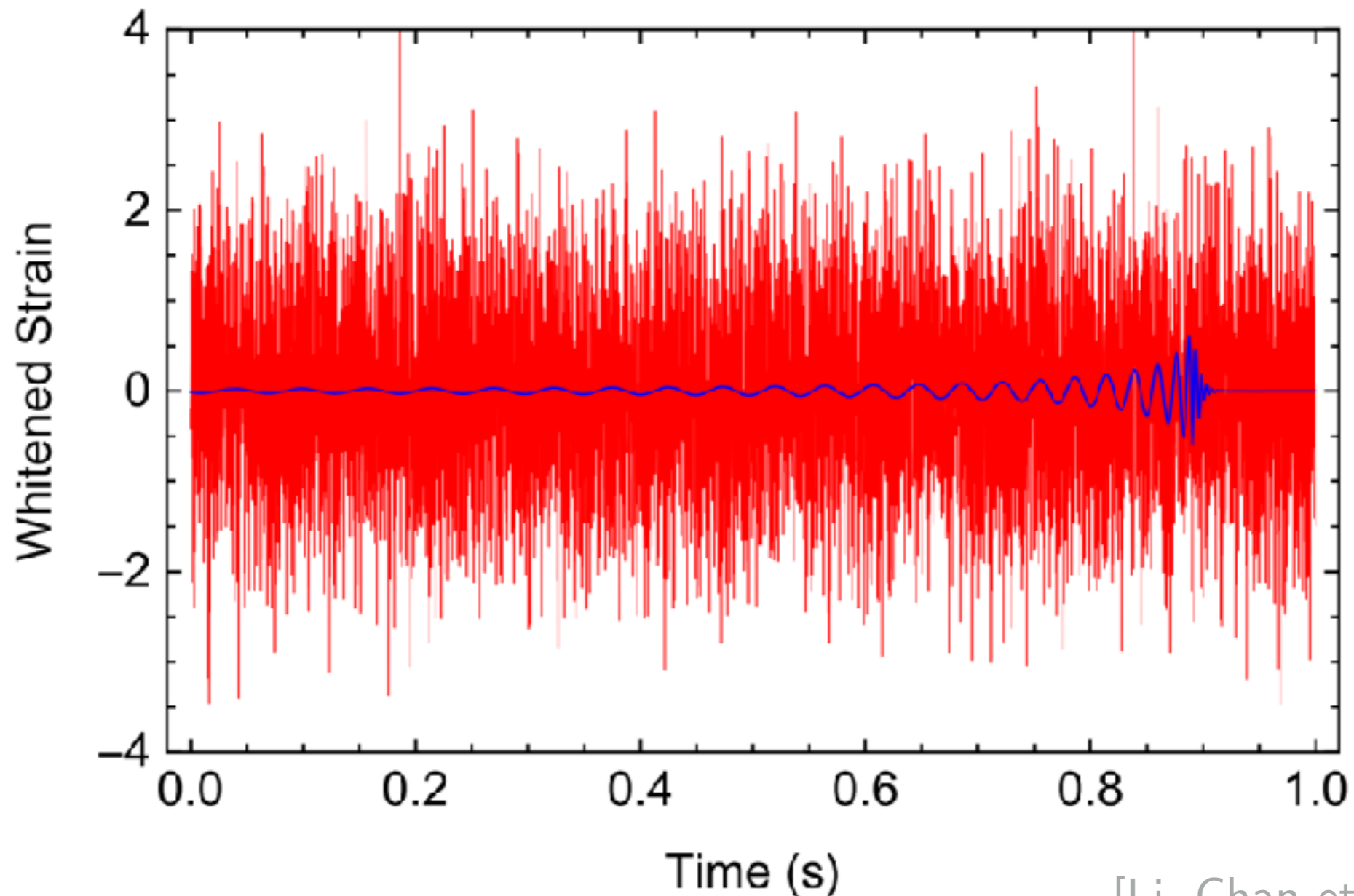
Joint parameter estimation

- Infer the parameters of the source under the *lensing hypothesis* using data from *multiple* events
- Allows for Bayesian model comparison
 - Measure consistency of events: coherence ratio (C^{L_U})
 - With source/lens populations priors, compute Bayes factor (B^{L_U})
- This is in principle all we need, but...
computationally costly ... astro priors ..



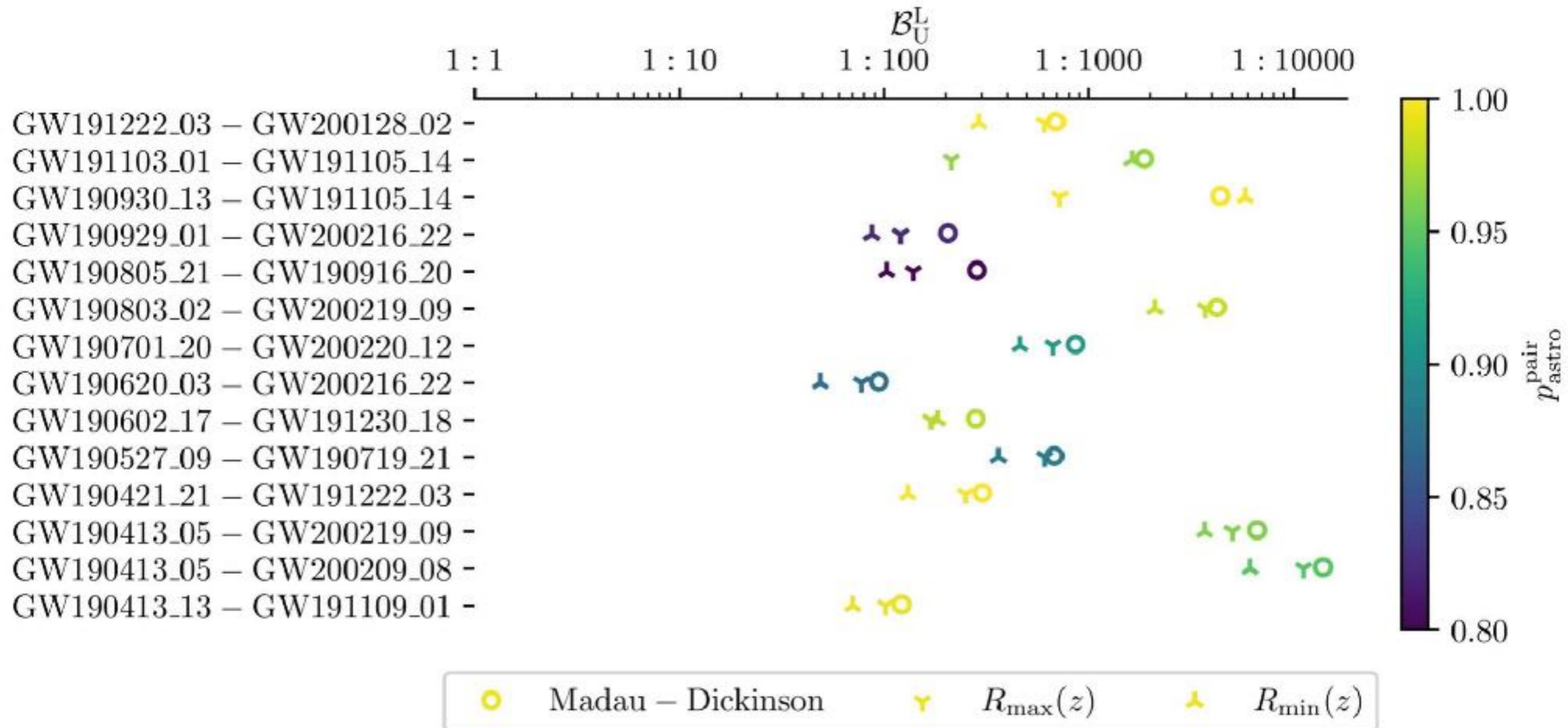
Sub-threshold searches

- Demagnified events could be under the noise (sub-threshold event)
- *Targeted* searches following super-threshold events reduce template bank and increase sensitivity



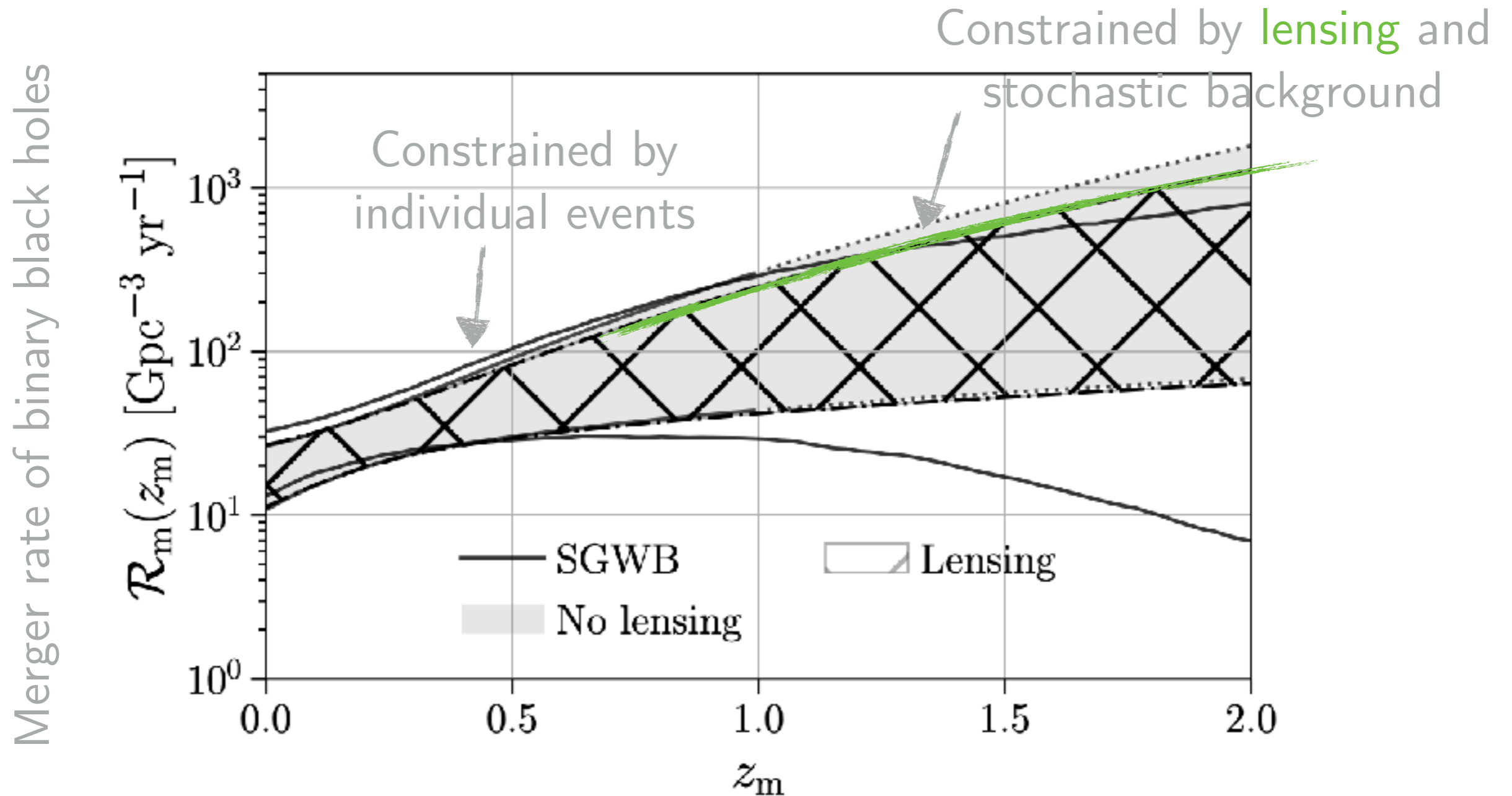
GWTC-3 results

- No evidence of repeated chirps in the data



GWTC-3 results

- *Upper bound* on binary black hole merger rate



Phase shifts & higher modes

- A gravitational wave is a superposition of *frequency modes*

$$h = \sum_{\ell, m \geq 0} \mathcal{A}_{\ell m} \cos[m(\Omega \Delta t + \varphi_c) - \chi_{\ell m}]$$

- A lensed signal of **type I** has different **amplitude** and **arrival time**

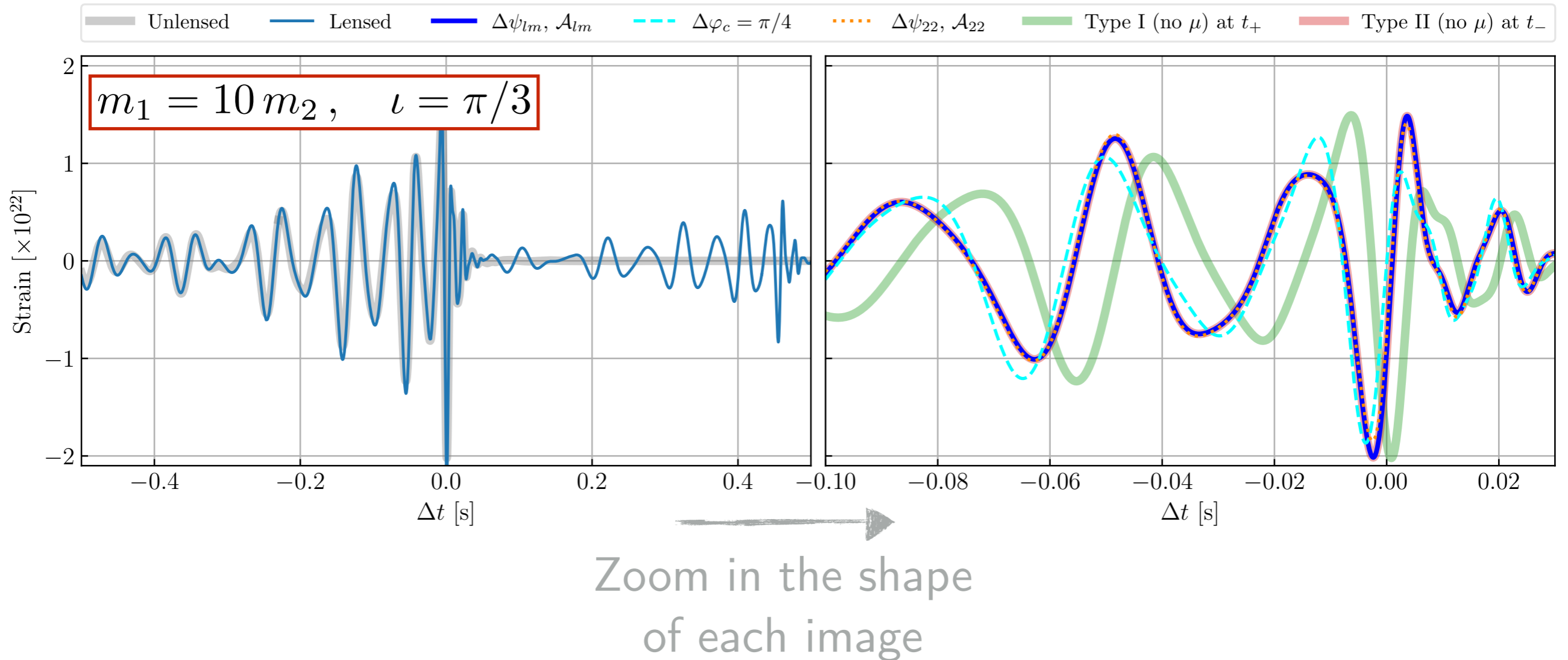
$$h_{\text{I}} = \sum_{\ell, m \geq 0} |\mu_{\text{I}}|^{1/2} \mathcal{A}_{\ell m} \cos[m(\Omega \Delta t_{\text{I}} + \varphi_c) - \chi_{\ell m}]$$

- A lensed signal of **type II** has also a **phase shift**

$$h_{\text{II}} = \sum_{\ell, m \geq 0} |\mu_{\text{II}}|^{1/2} \mathcal{A}_{\ell m} \cos \left[m (\Omega \Delta t_{\text{II}} + \varphi_c) - \chi_{\ell m} + \frac{\pi}{2} \right]$$

Waveform distortions in **type II** images

- Lensing imprints *small* but *characteristic* modifications in the signals that cannot be mapped to other astrophysical parameters



Caustics

- For point sources, there are **singular** points in the lens mapping

$$\det \left(\frac{\partial^2 T_d(\theta_j)}{\partial \theta_a \partial \theta_b} \right) \rightarrow 0 \quad \Rightarrow \quad \mu(\theta_j) \rightarrow \infty$$

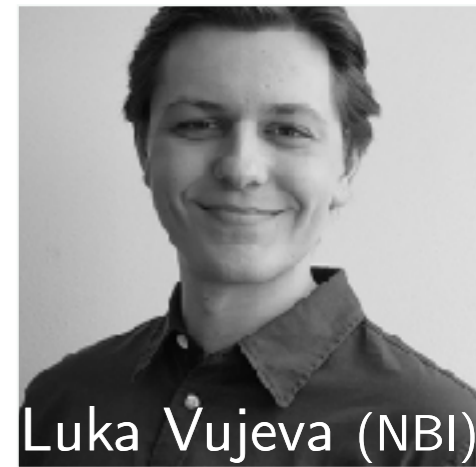
- Caustics exhibit **universal** behaviors (described by catastrophe theory)

$$\mu_{\pm} \sim 1/\sqrt{\Delta\theta_S} \sim \Delta t^{-1/3}$$

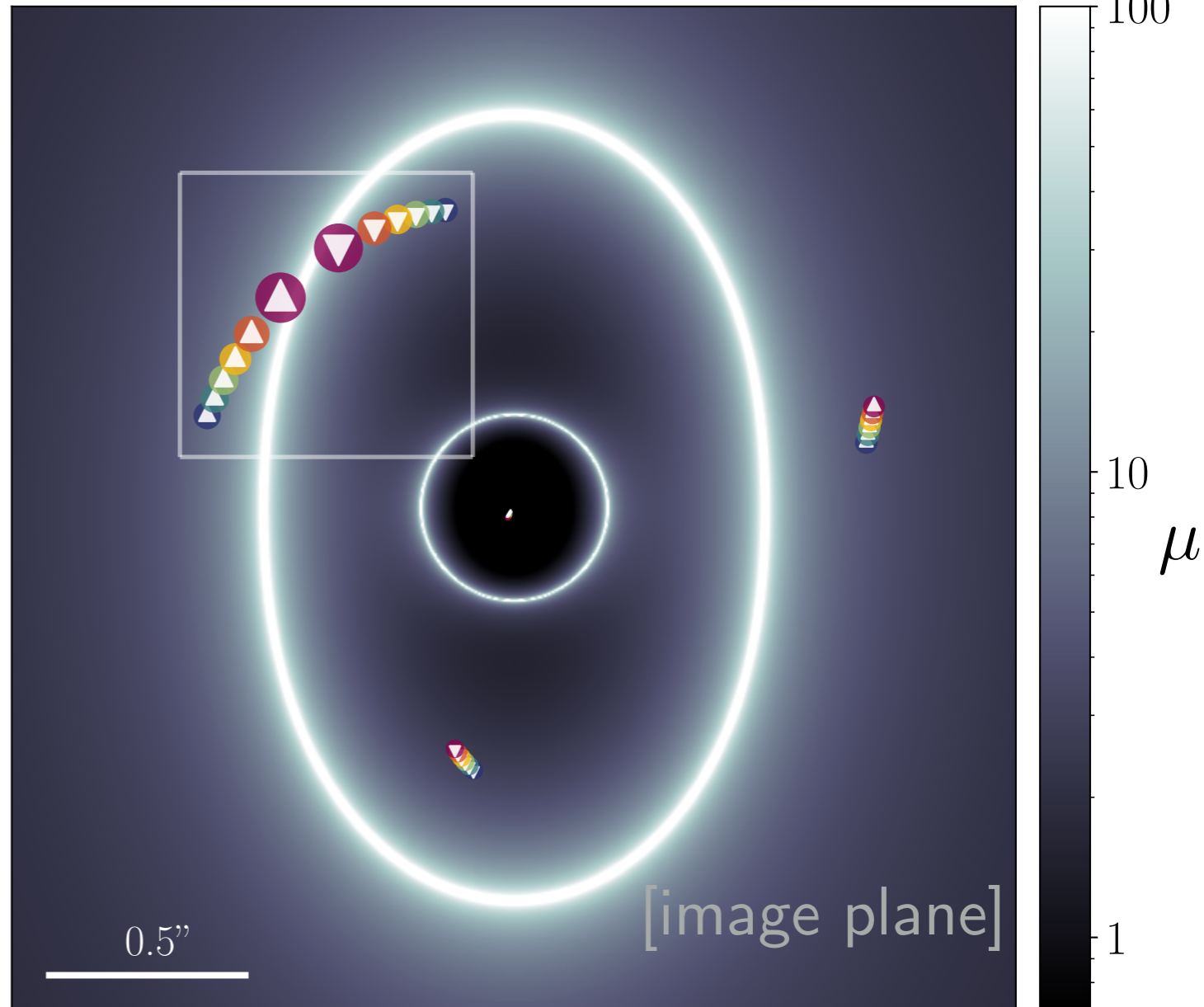
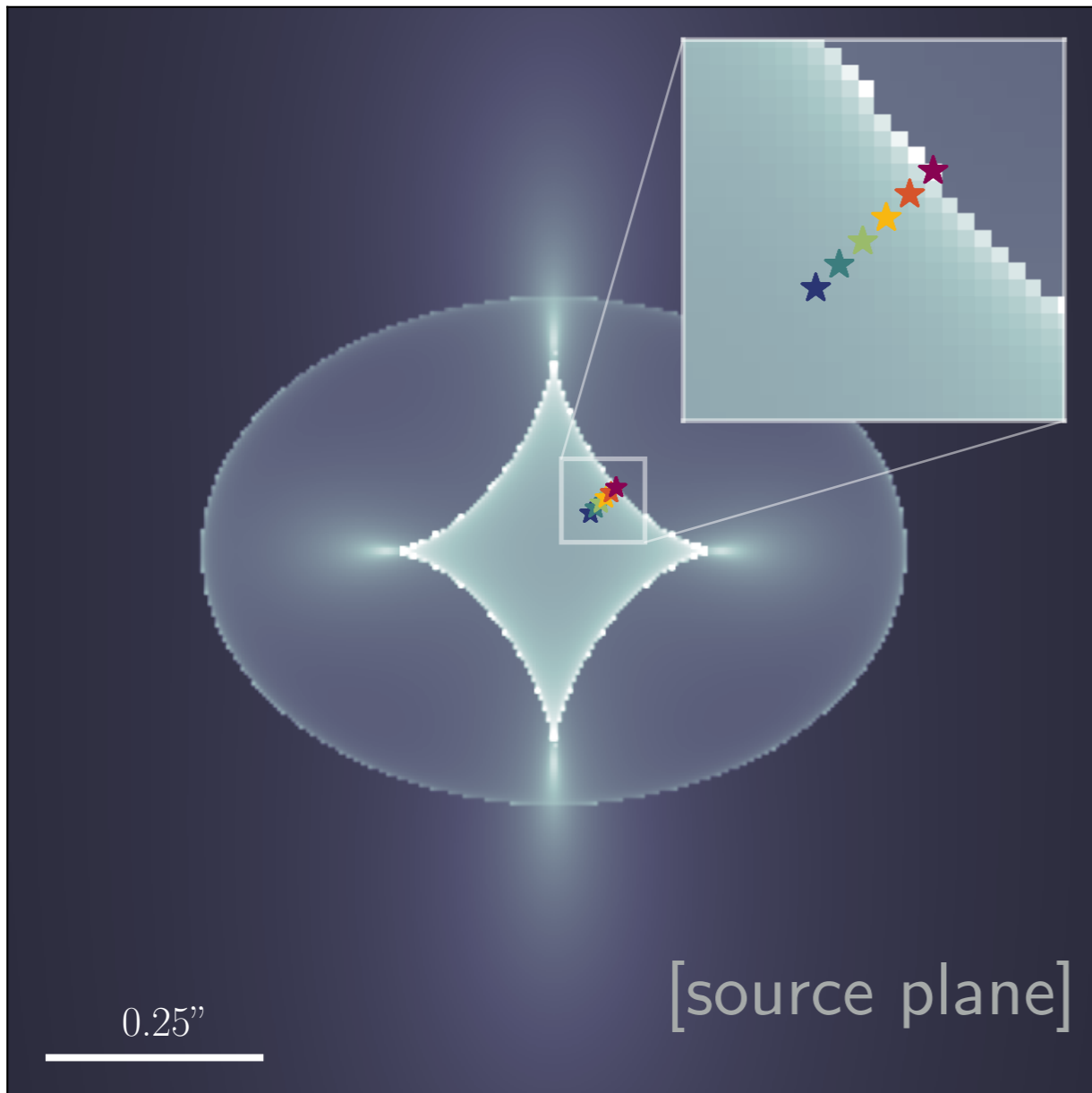
- SPA is **broken** when approaching to a caustic
- **Maximum magnification** set by diffraction



Approaching a (fold) caustic

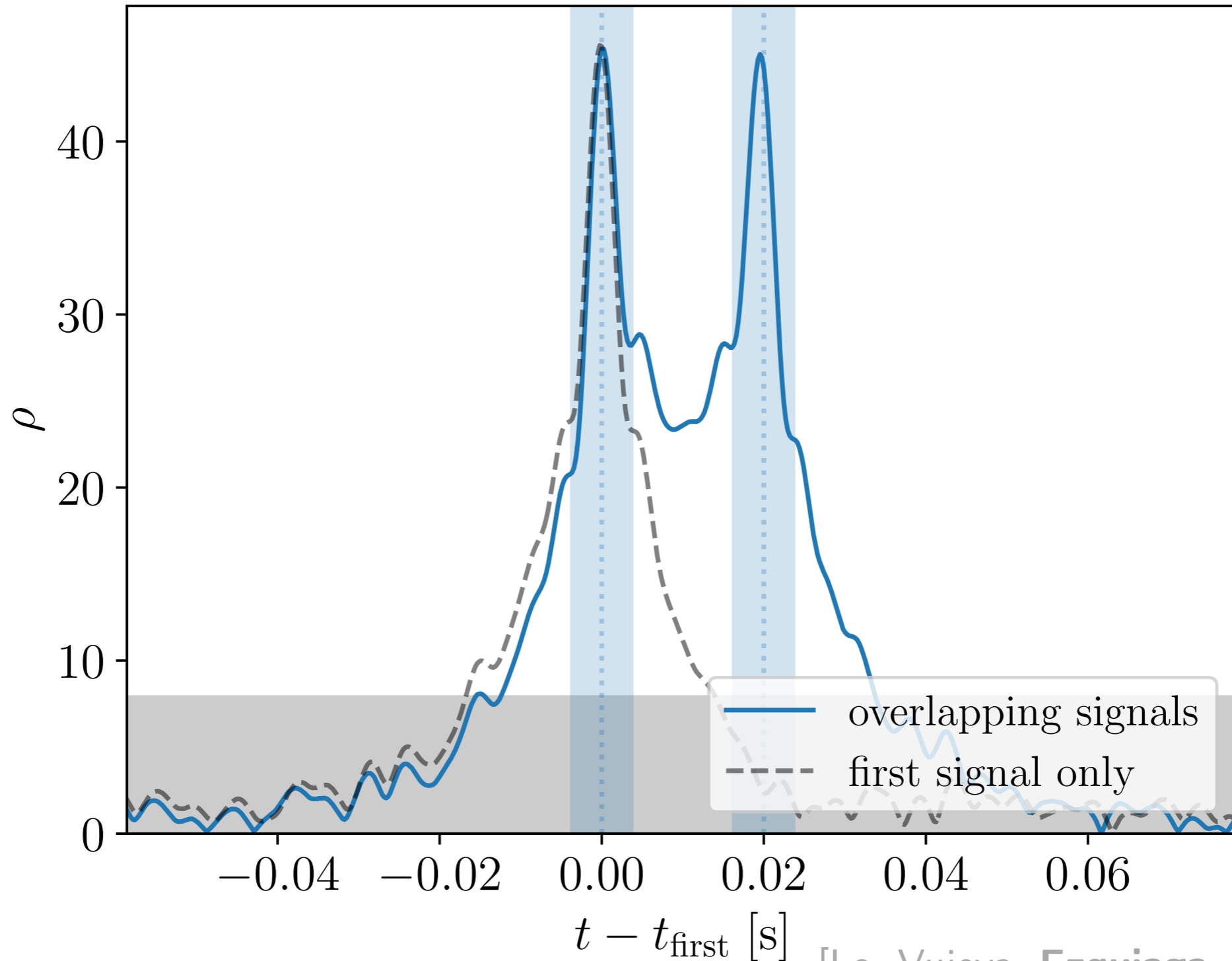
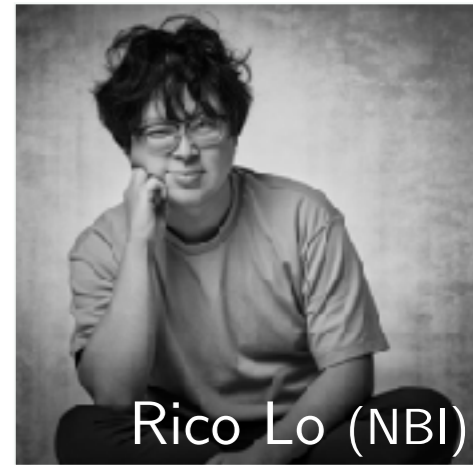


Luka Vujeva (NBI)



[galaxy lens with a cored singular isothermal ellipsoid density profile]

Highly magnified, overlapping signals



Wave optics

$$\Delta t_d \cdot \omega$$

- Time delay scales with the lens mass

$$\Delta t_d(y = 1) \simeq 4 \left(\frac{(1 + z_L) M_L}{100 M_\odot} \right) \text{ ms} \quad [\text{point mass lens}]$$

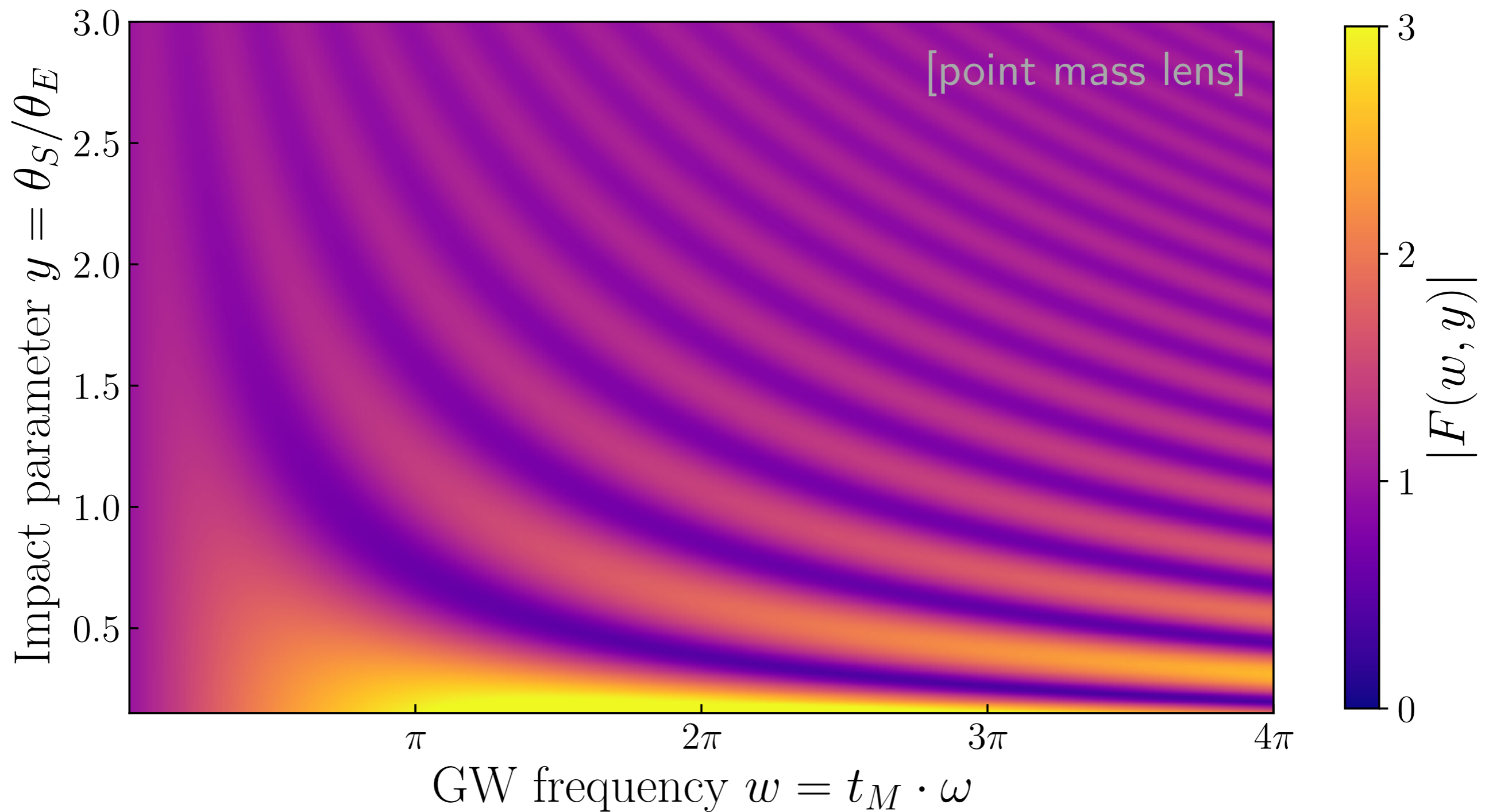
- GW frequency scales with binary mass (*has astrophysical size!*)

$$f \sim \frac{1}{2\pi} \frac{1}{2t_{\text{Sch}}} \sim 800 \text{ Hz} \left(\frac{10 M_\odot}{M} \right)$$

- Wave optics regime: $\Delta t_d \cdot \omega \sim 1$

- Low-frequency limit has small lensing $\omega \rightarrow 0 \Rightarrow F \rightarrow 1$

Wave optics: **diffraction**



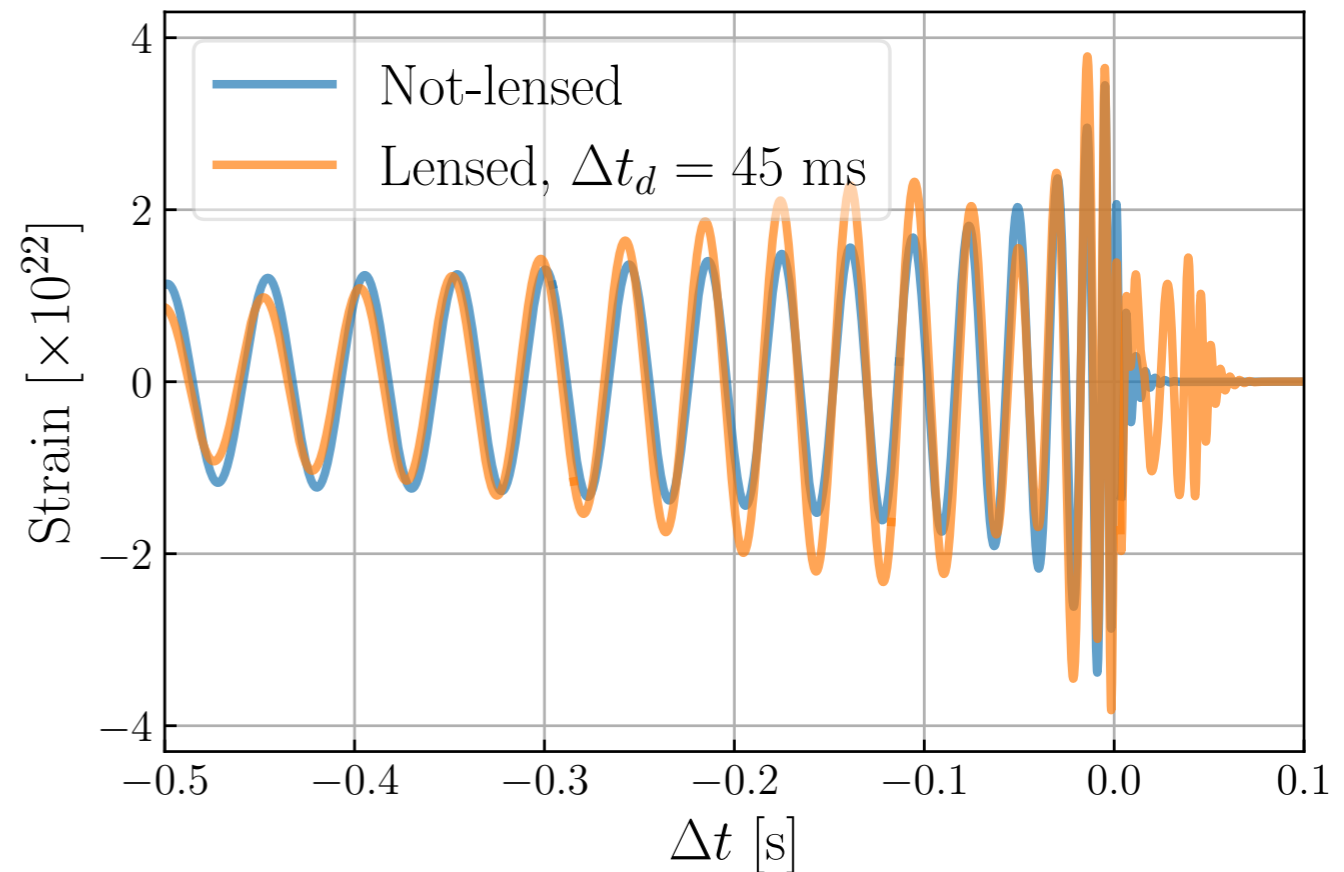
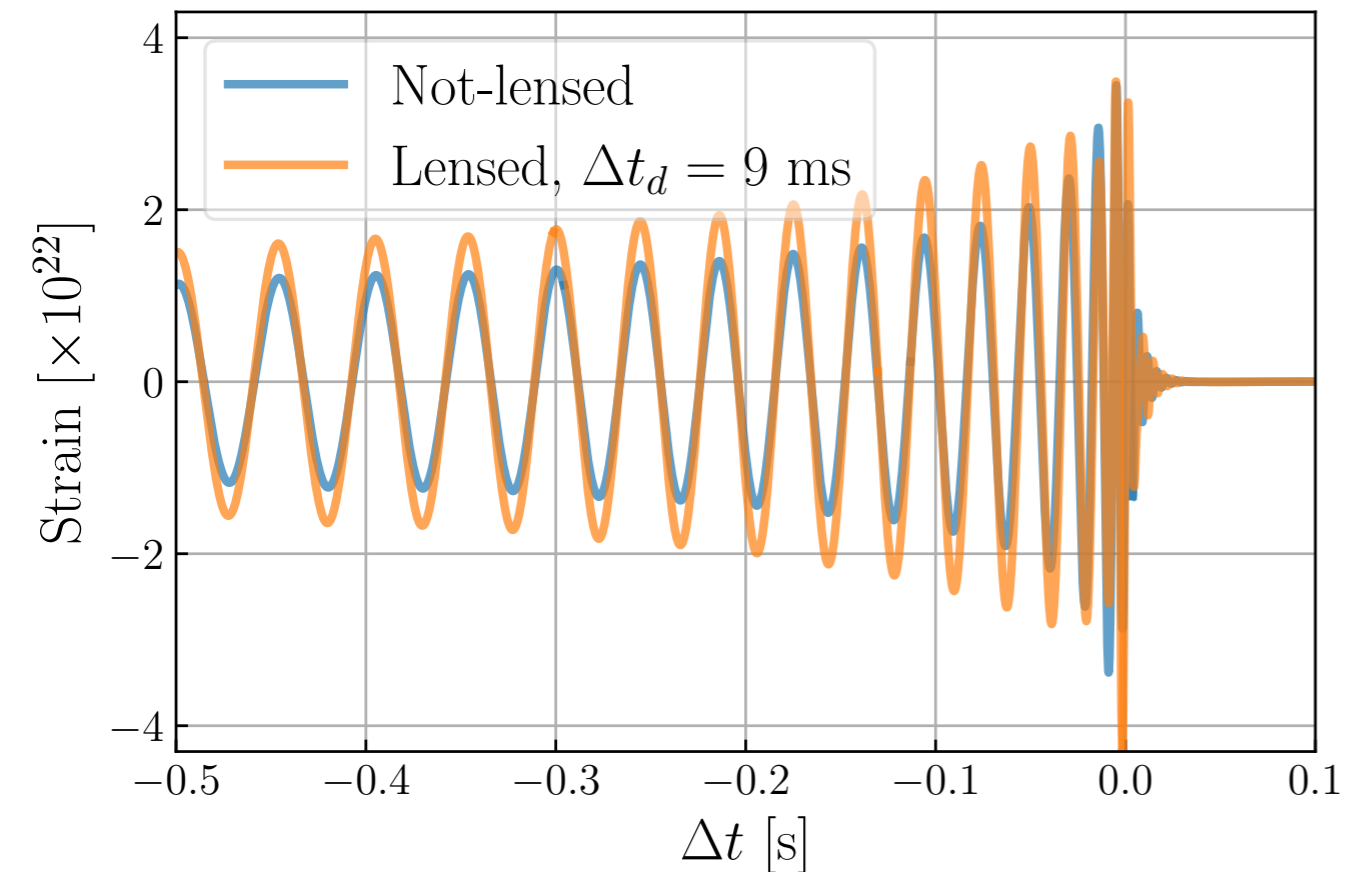
[see Ubach's talk]

E.g. compact (point) lenses

$$\Delta t_d(y = 1) \simeq 4 \left(\frac{(1 + z_L) M_L}{100 M_\odot} \right) \text{ ms}$$

Diffraction

Interference



- Most lens models require solving the diffraction integral numerically. Great recent progress [see Villarubia-Rojo's talk]

Parameter estimation of lensed signals

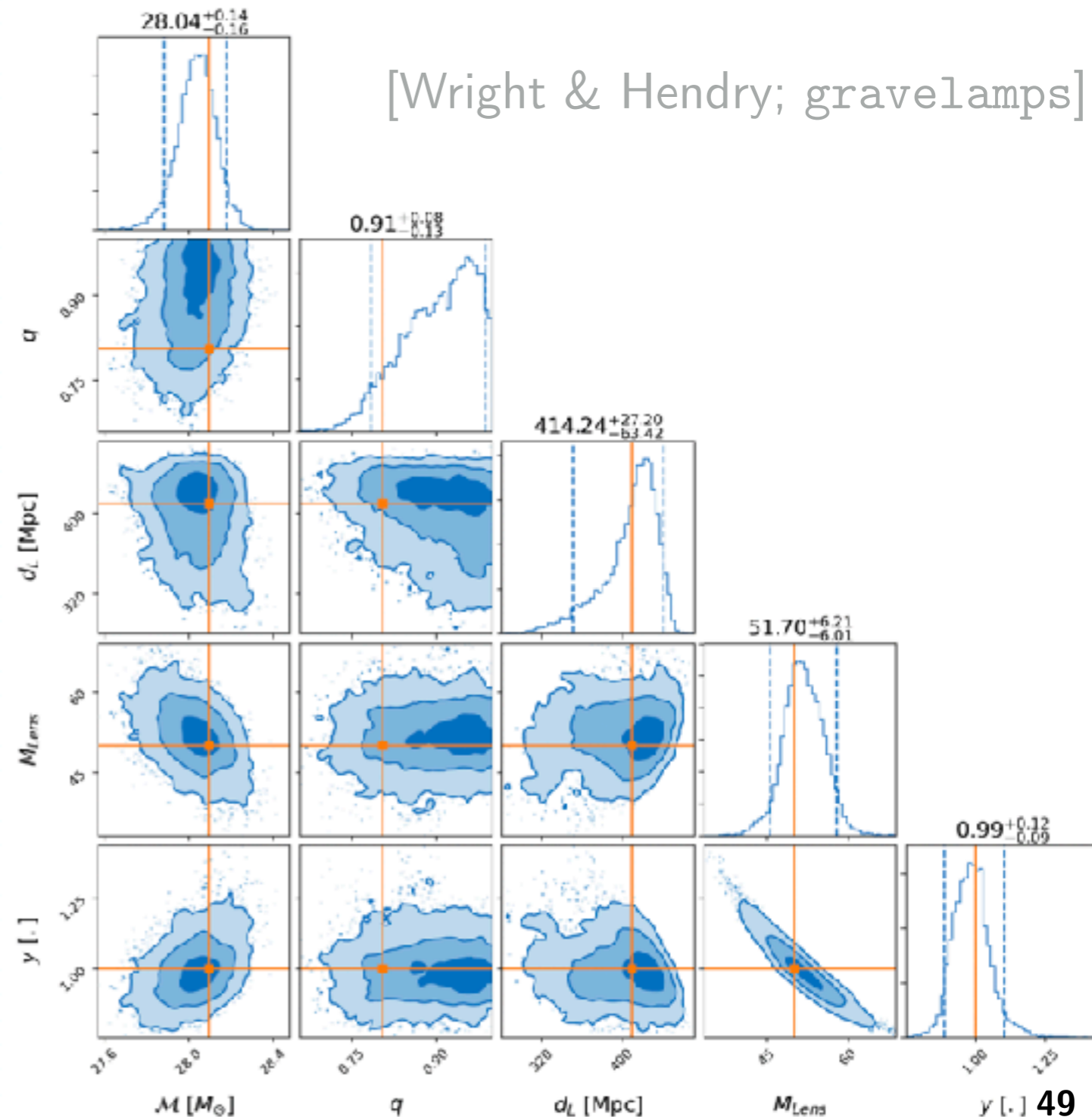
- Include lensed parameters in the inference

- Requires efficient models for the computation of the amplification function.

- Allows to make Bayesian model comparison (B^L_U)

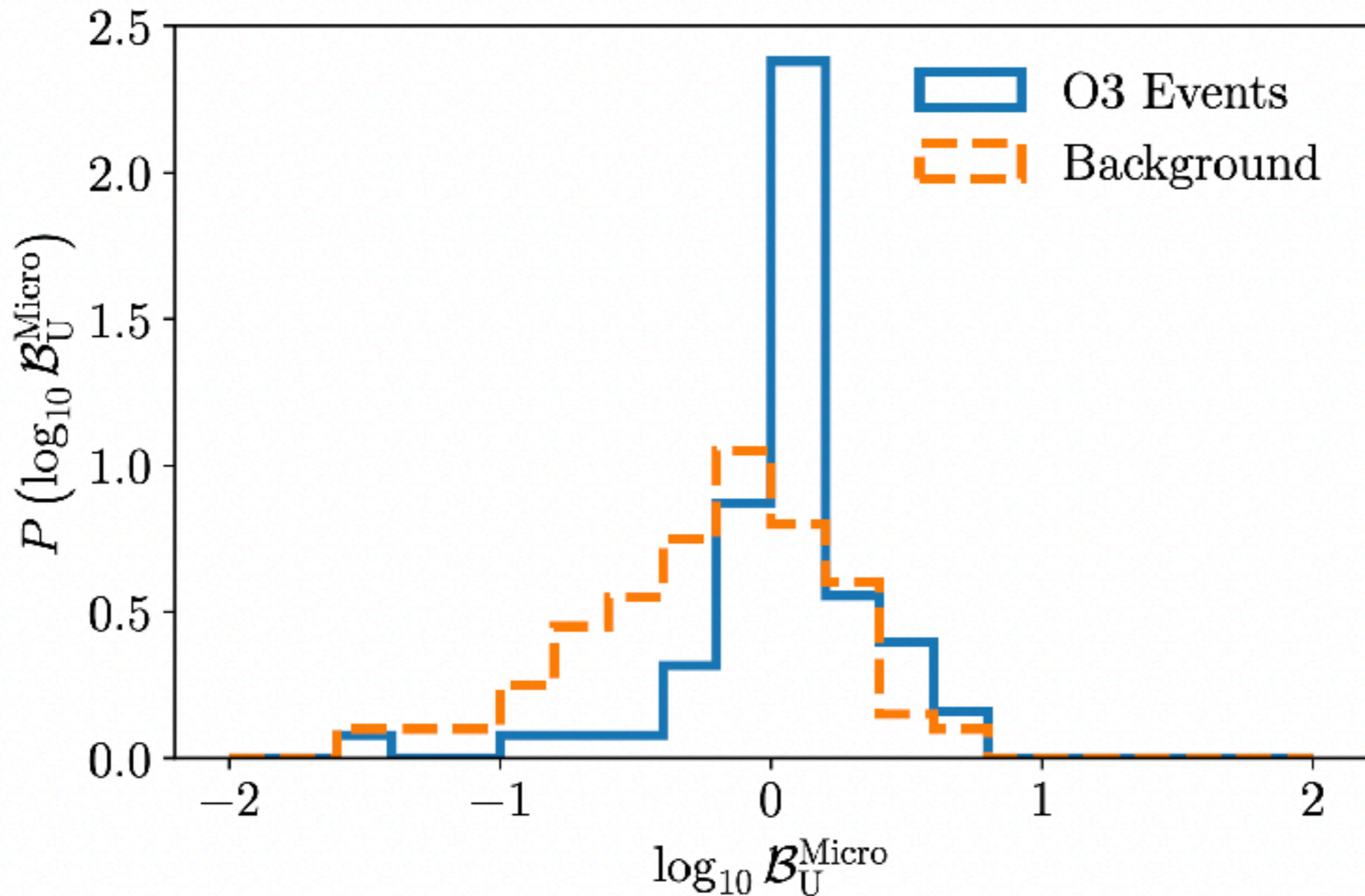
- Addressing possible *waveform systematics* and *noise artifacts* is crucial!

[Janquart et al.; MNRAS'23]



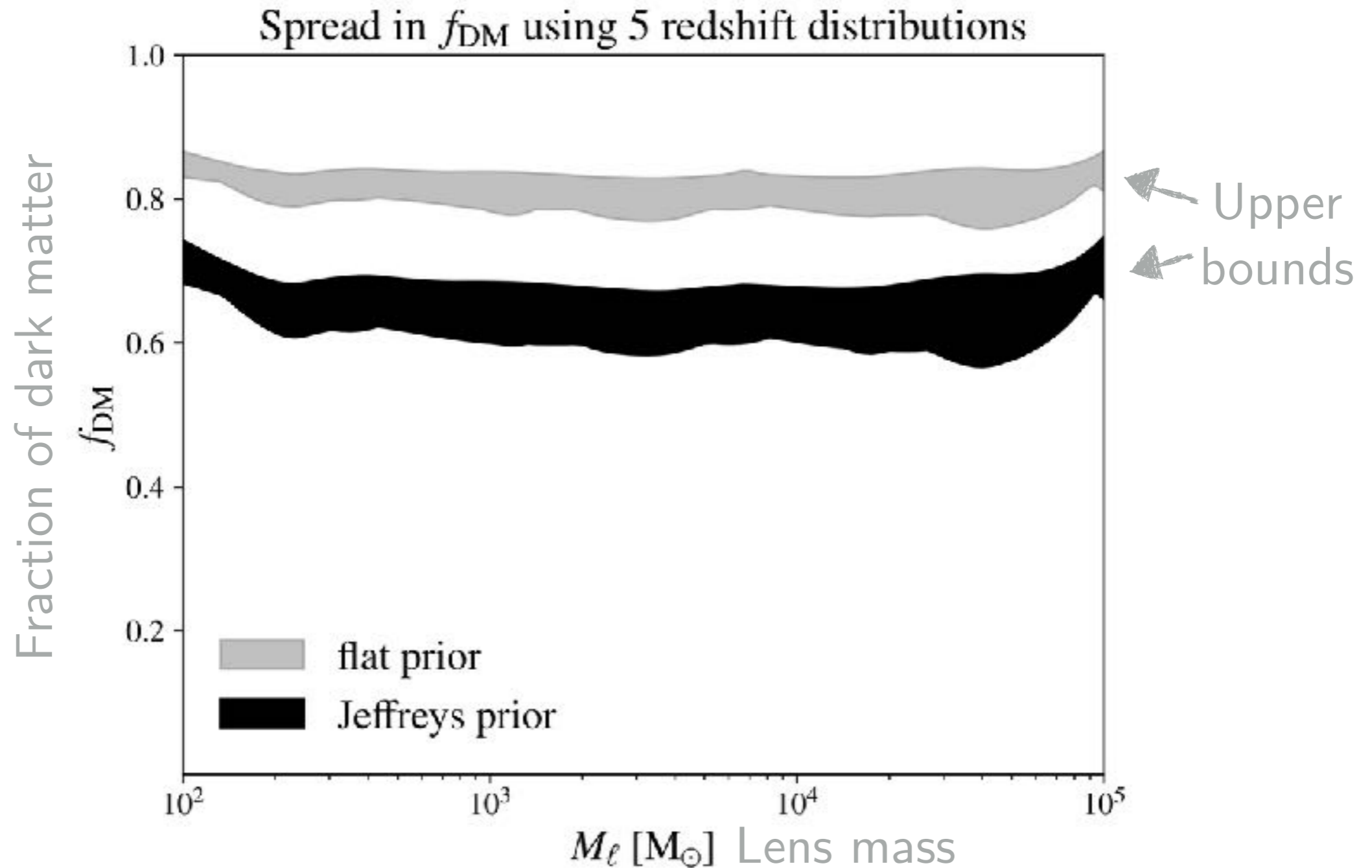
GWTC-3 results

- No evidence of distorted waveforms by lensing (“microlensing”)



GWTC-3 results

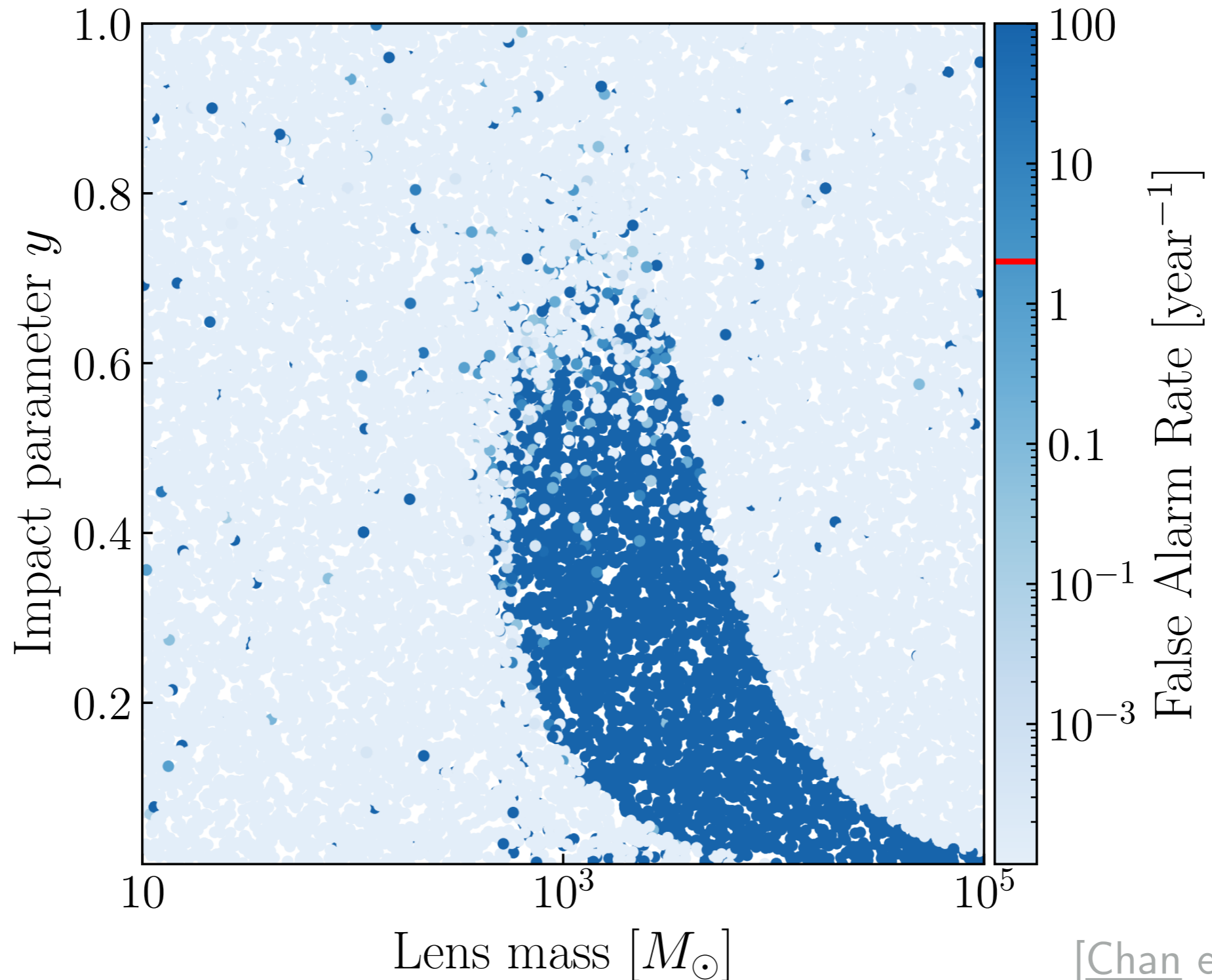
- *Upper bound* fraction of compact lenses (w.r.t. dark matter)



Searching for lensed GWs



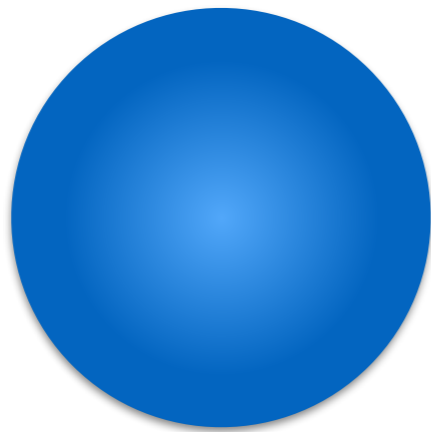
- Distorted waveforms could be missed by current searches!



4. Future prospects

Substructures

- Gravitational waves are effectively point sources. They are very sensitive to *small scales*



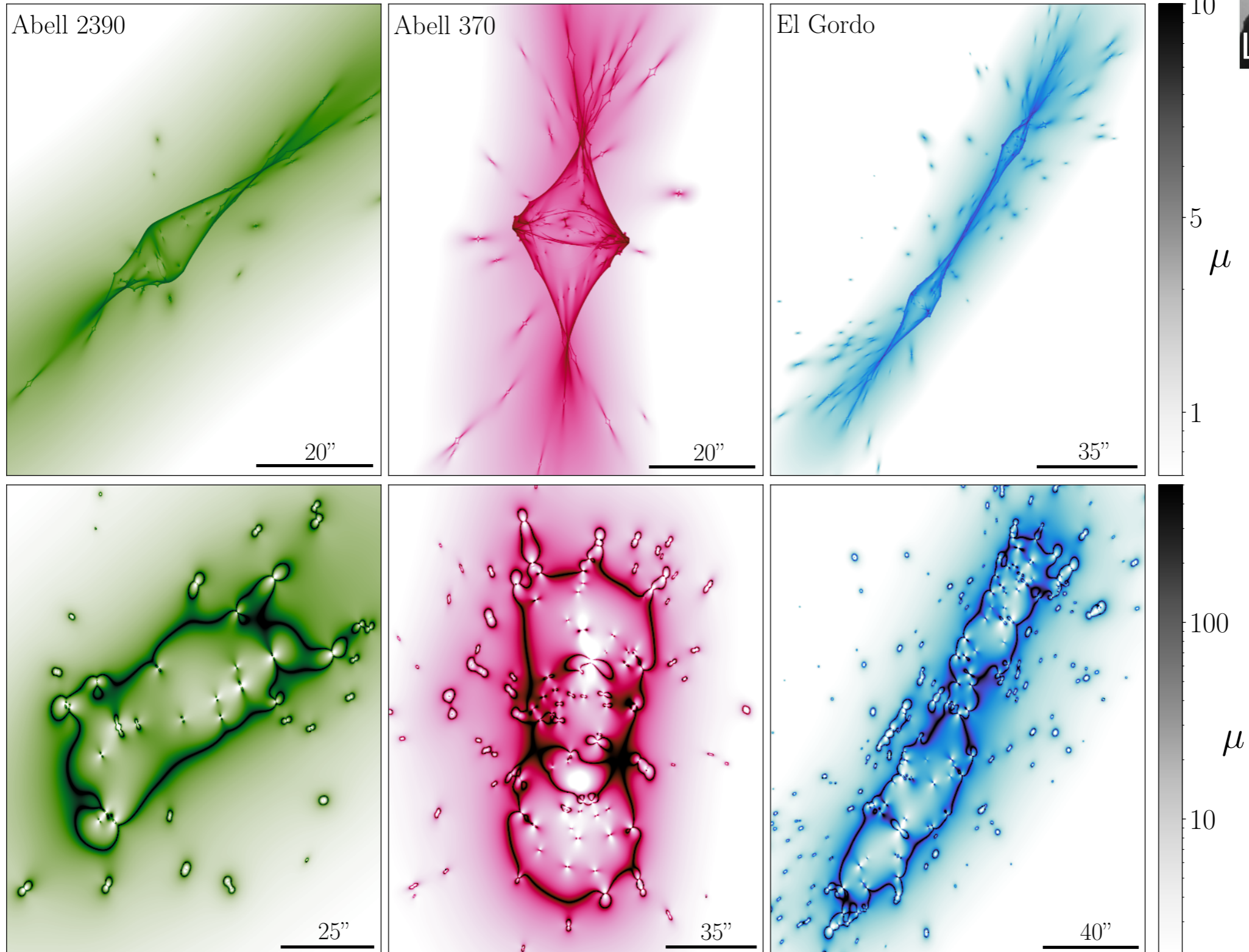
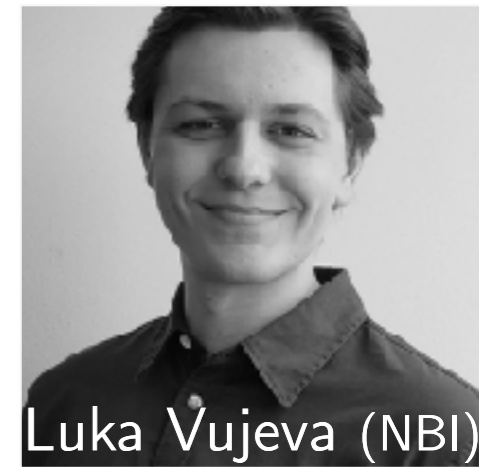
[Singular Isothermal Sphere]

v.s.



[Hubble Space Telescope]

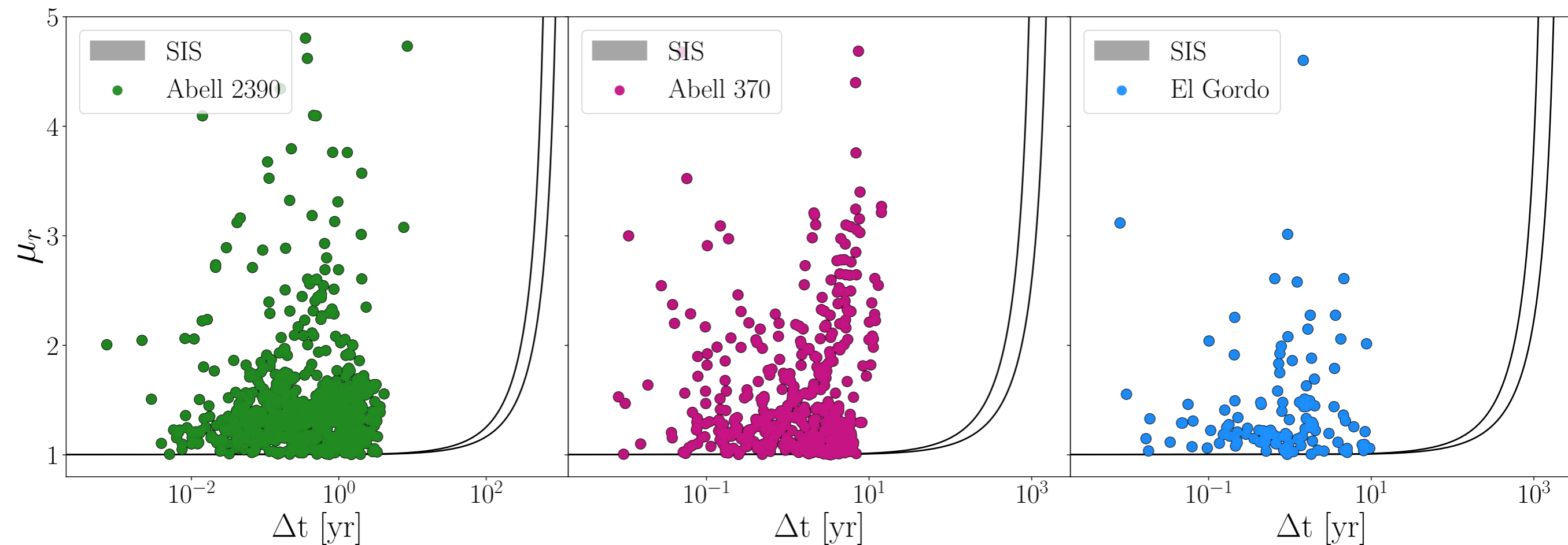
Substructures - clusters



Substructures - clusters

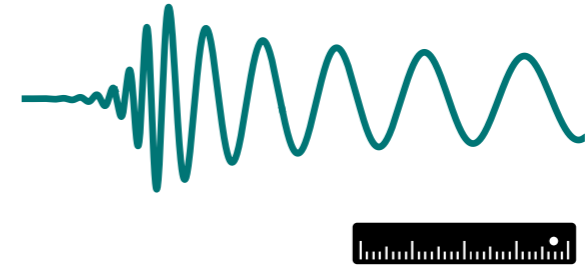
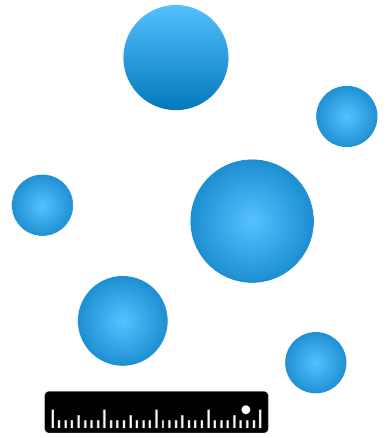


- In real clusters, relative magnifications and time delays change dramatically compared to singular isothermal sphere (SIS)

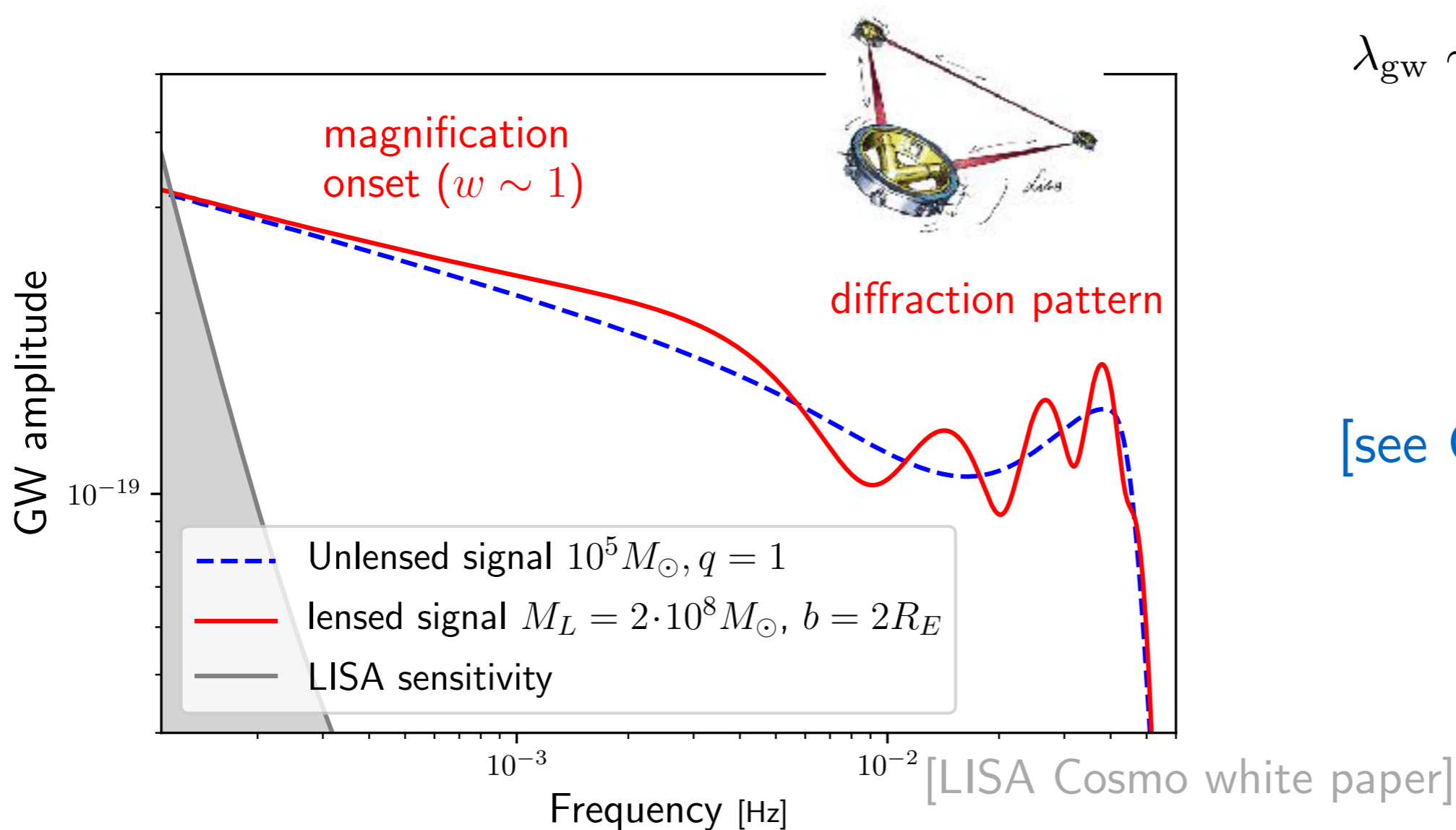


Substructures - subhalos

- Dark matter halos are made of smaller halos
- Gravitational waves could interfere!

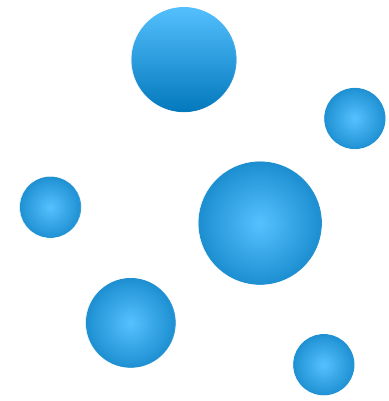


$$\lambda_{\text{gw}} \sim 10^3 \text{ km} \left(\frac{M_{\text{bbh}}}{10 M_{\odot}} \right)$$

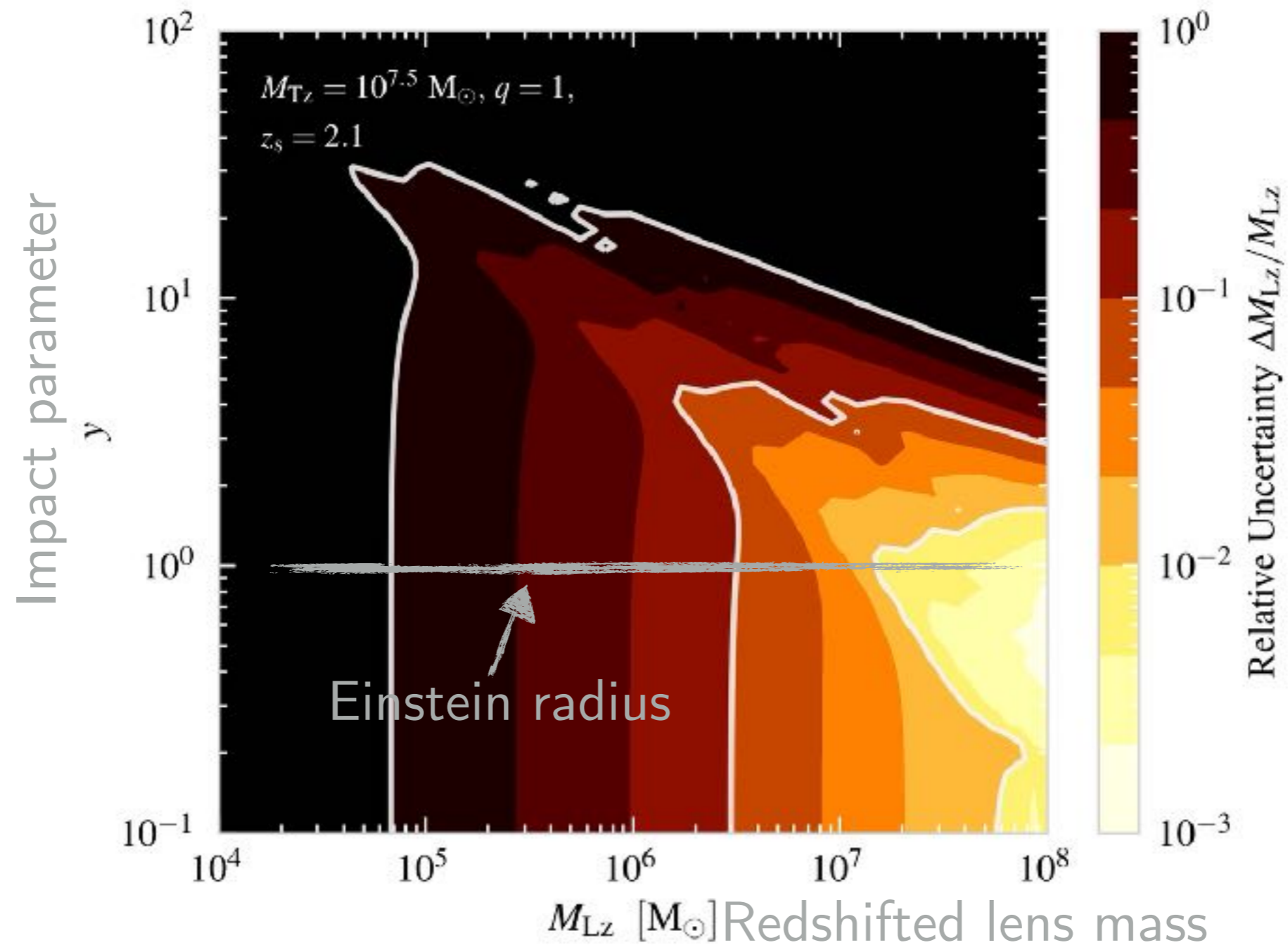


[see Goyal's talk]

Substructures - subhalos

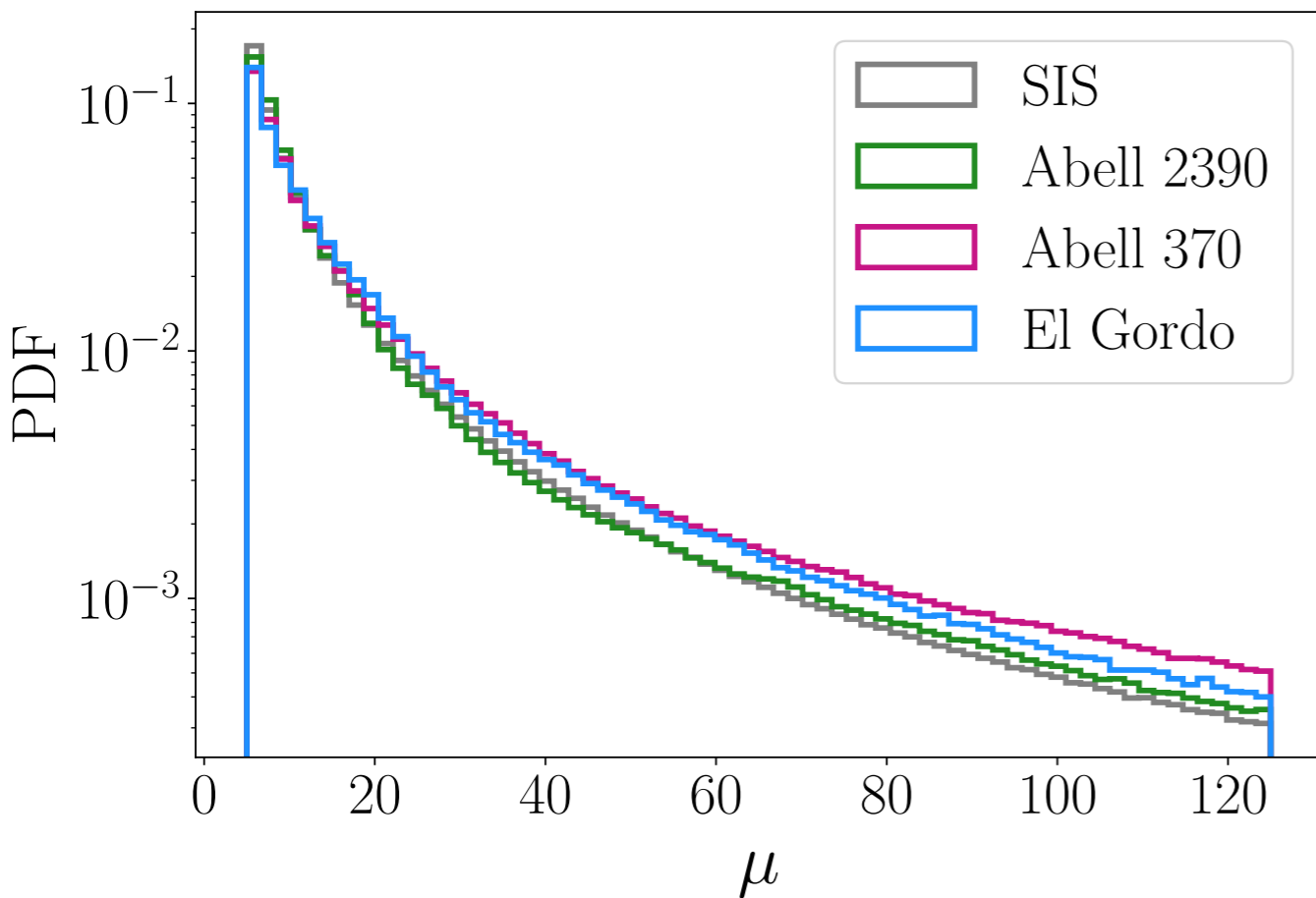


- *Sensitive to lensing beyond the Einstein radius!*

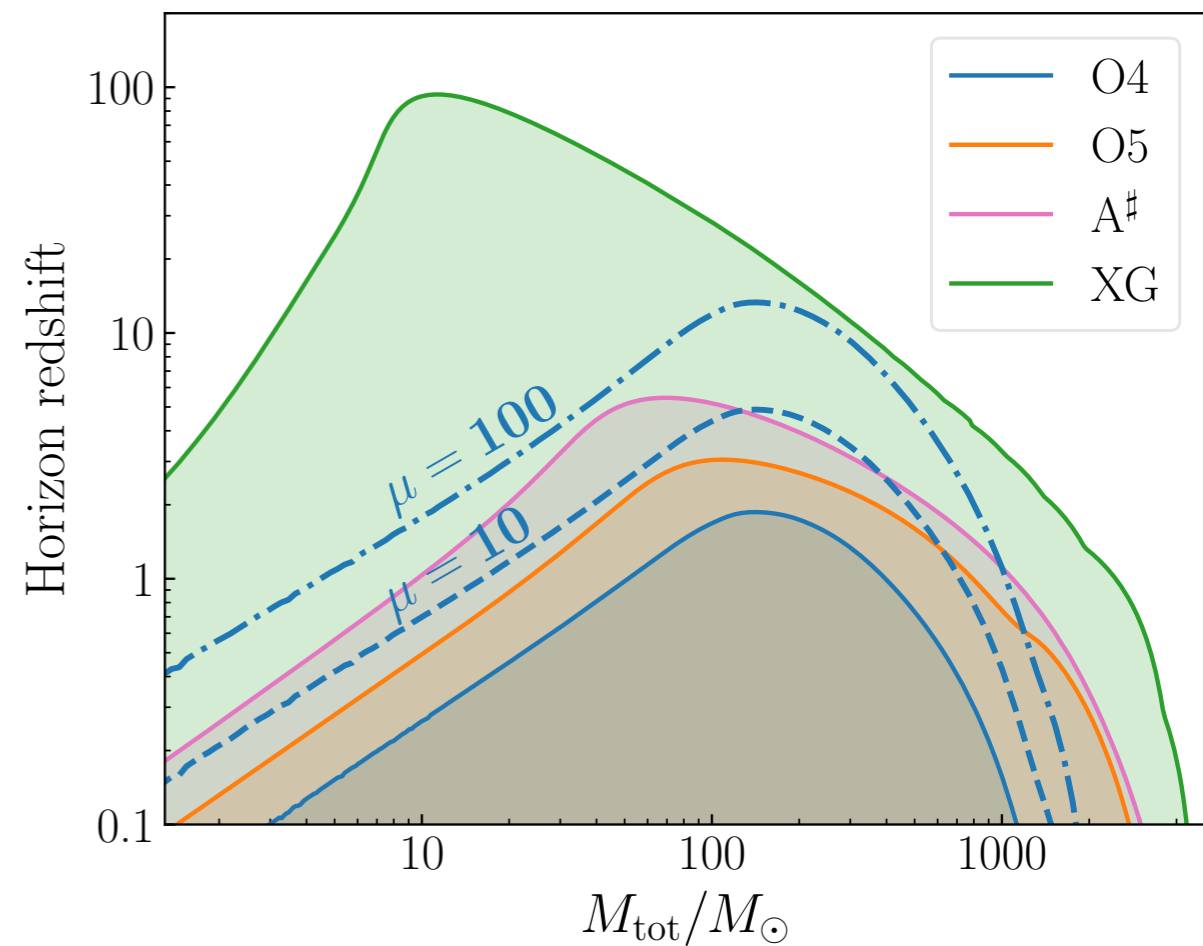


Highly magnified gravitational waves

- Substructures can enhance high magnification tail
- Even more sensitive to small lenses close to the caustics!



[[Vujeva](#), [Ezquiaga](#), [Lo](#), [Chan](#); *to appear*]



[[Lo](#), [Vujeva](#), [Ezquiaga](#), [Chan](#); 2024]

Multi-messenger lensing

- Observe multi-messenger lensed events, e.g. *binary neutron stars* with ground-based or *super-massive black hole binaries* with LISA

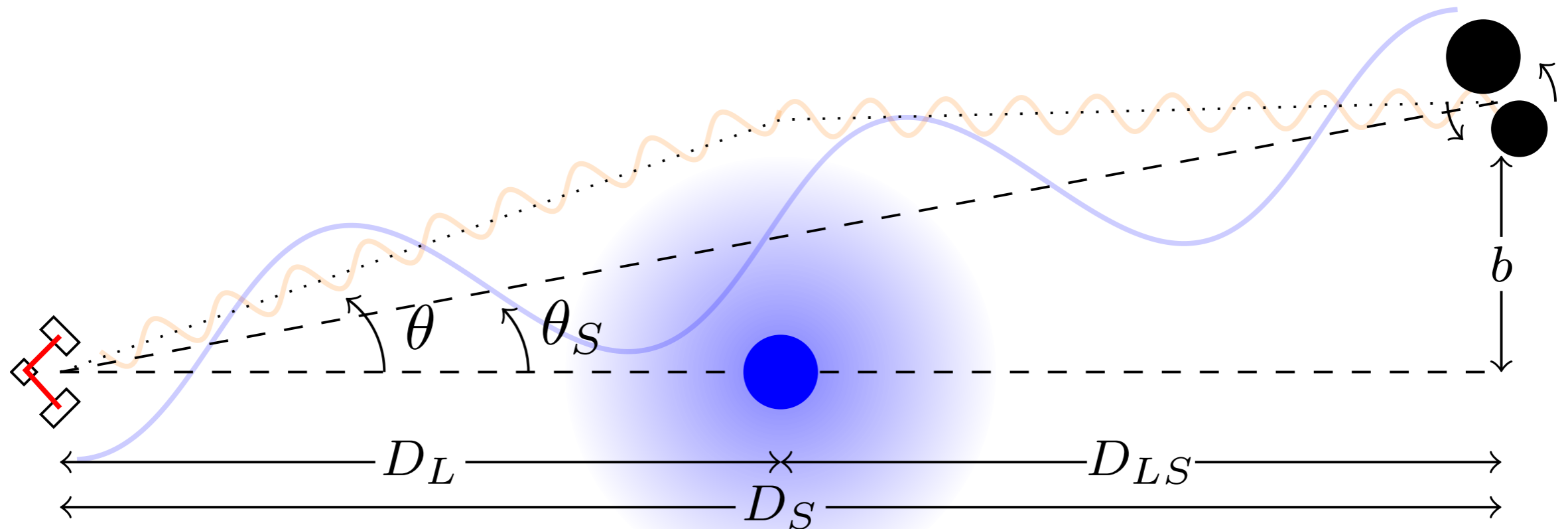


[see Smith's talk]

- Great target for future detectors!
- Will open many science cases

Multi-messenger lensing & wave optics

- Gravitational waves and photons could suffer lensing in different regimes

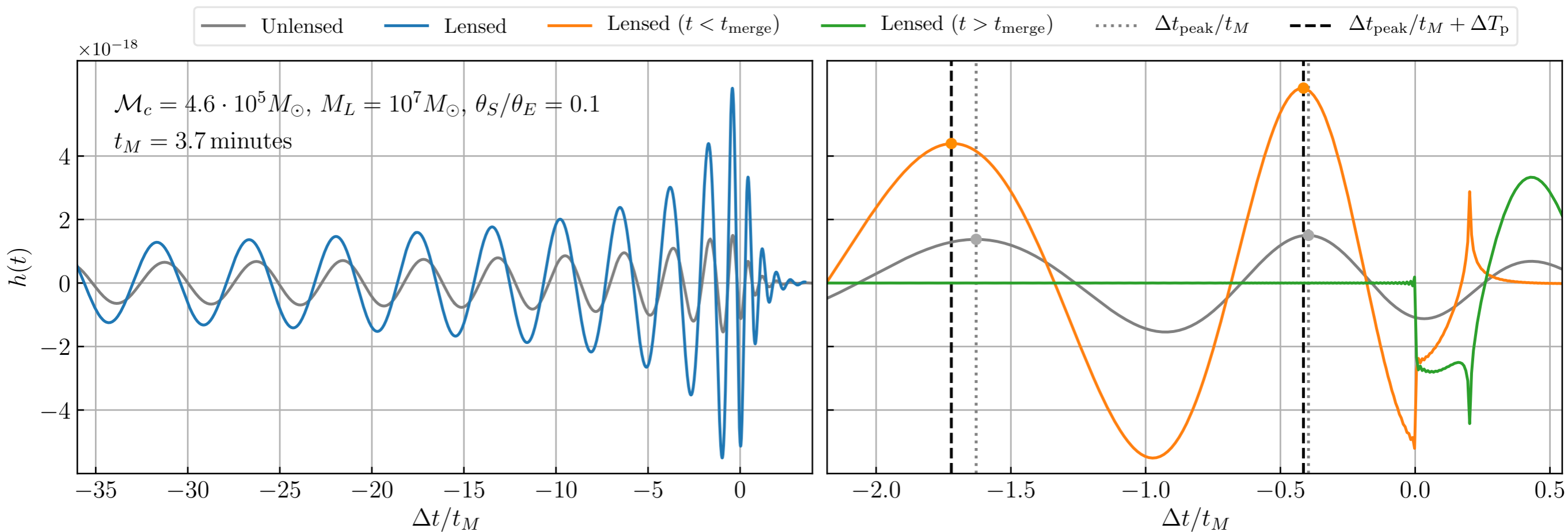


- Phase* and *group velocity* may change in wave optics

$$t_p(\omega, \vec{\theta}_S) = -\frac{i}{\omega} \ln \left(\frac{F(\omega, \vec{\theta}_S)}{|F(\omega, \vec{\theta}_S)|} \right). \quad t_g(\omega, \vec{\theta}_S) = t_p(\omega, \vec{\theta}_S) + \omega \frac{\partial t_p(\omega, \vec{\theta}_S)}{\partial \omega}$$

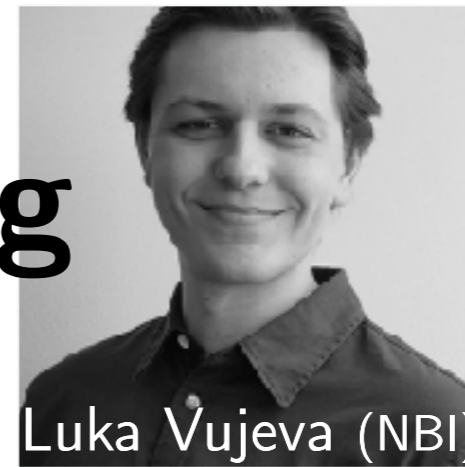
Multi-messenger lensing & wave optics

- There is an *apparent superluminality* due to the waveform distortions

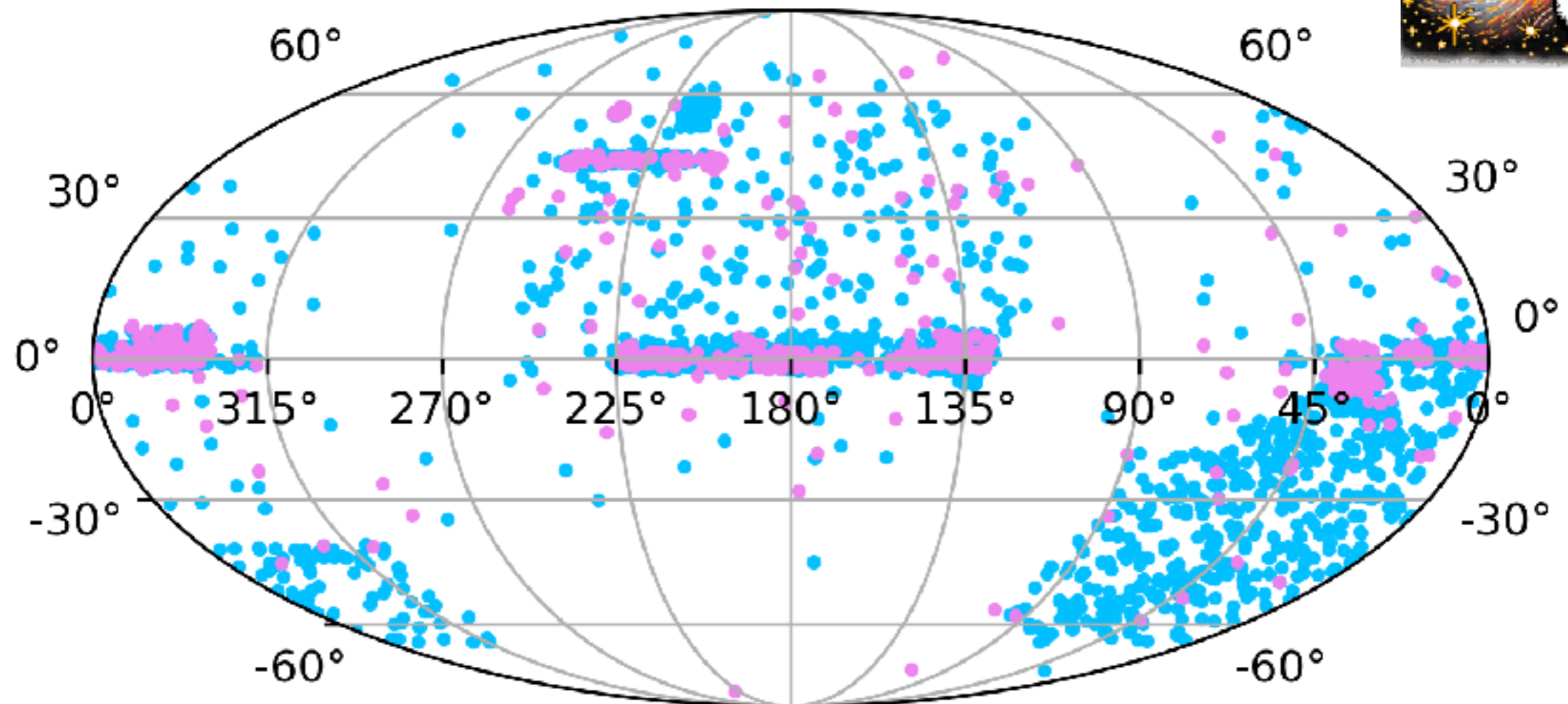


[example for super-massive black hole binary]

Multi-messenger lensing

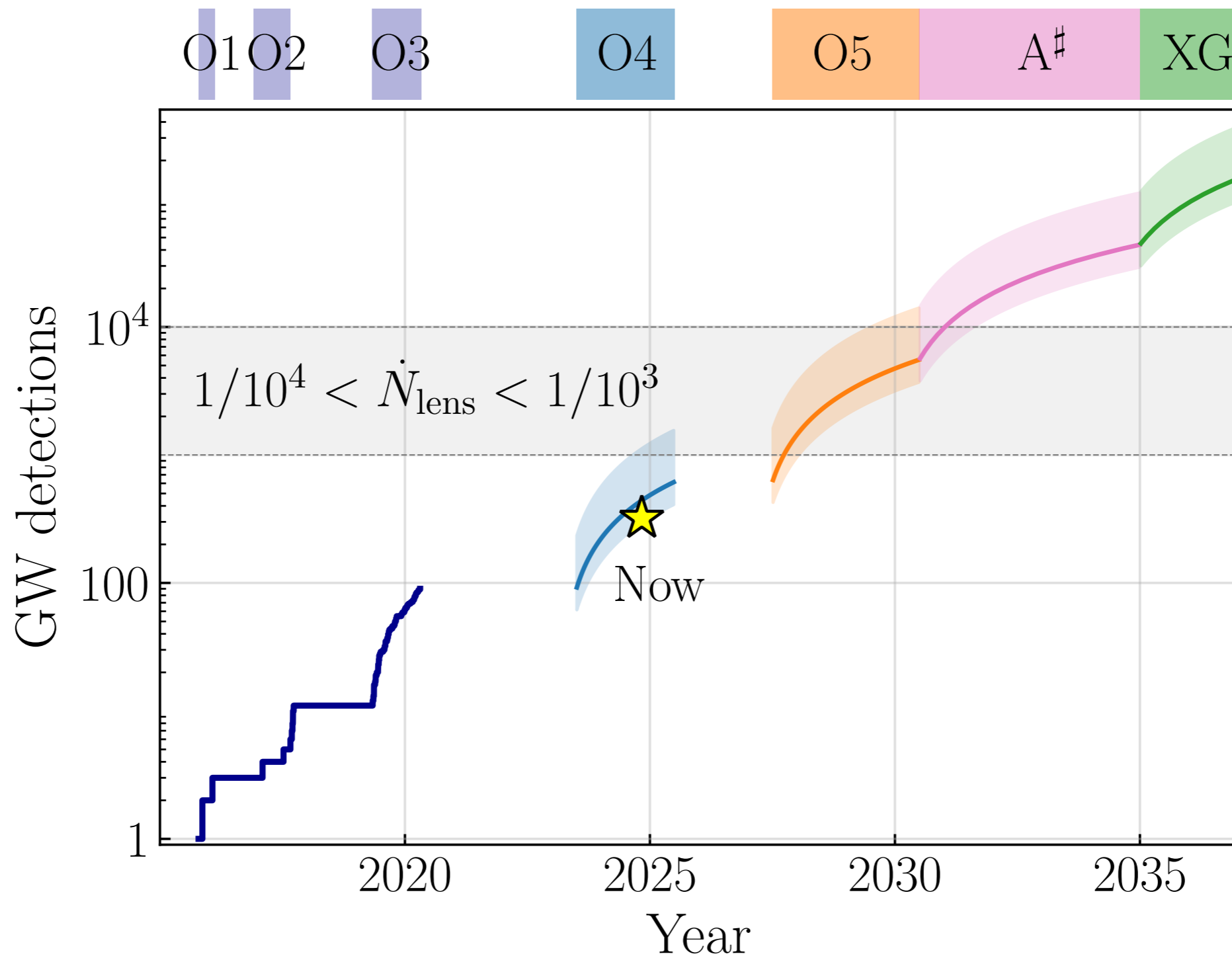


- Cross match GWs with *lens catalogs*
- Identify lensed host galaxy (*difficult!*)
- Watchlist for efficient lenses



GW lensing with next-generation detectors

- Large number of detections enable statistical studies

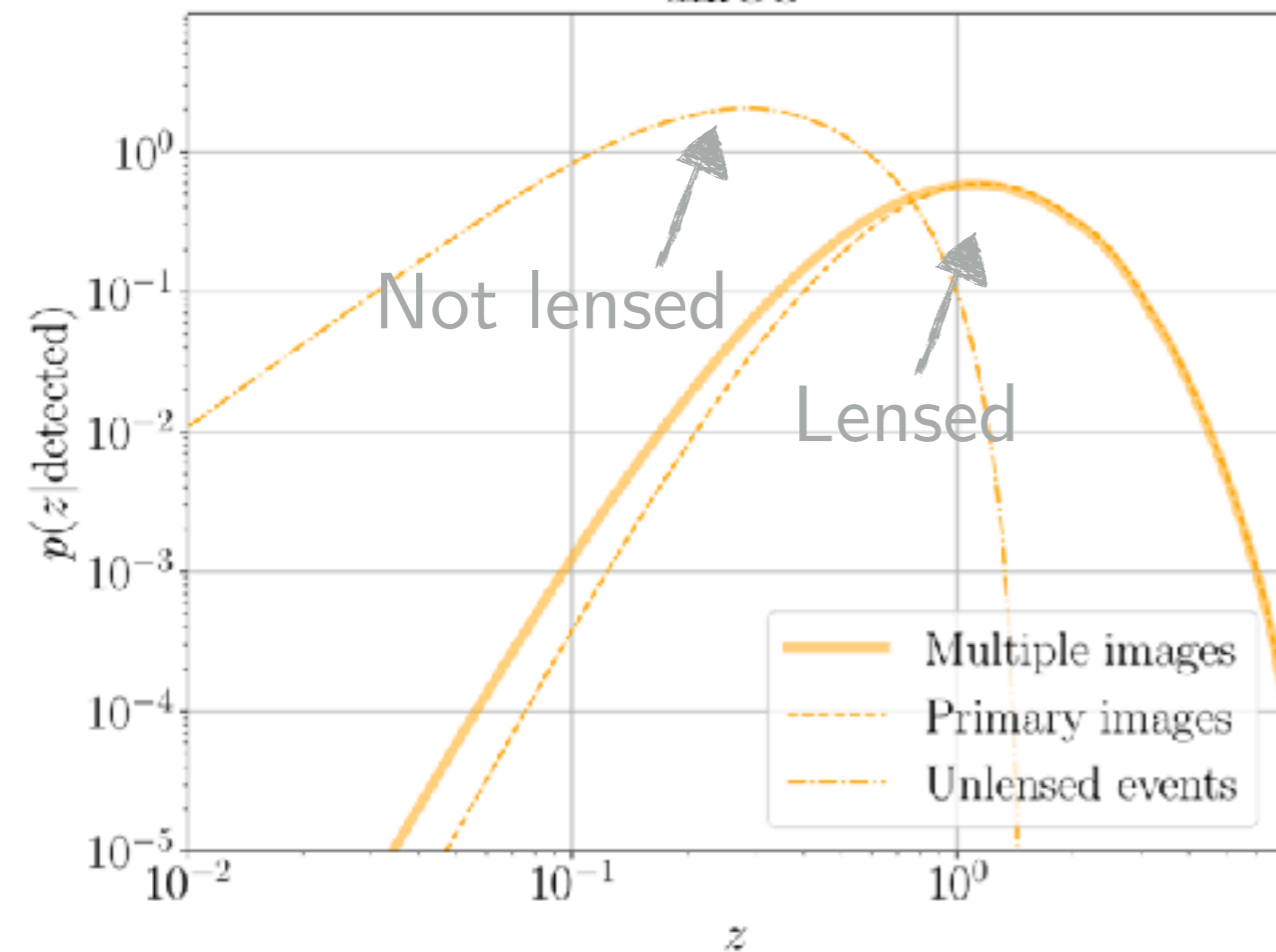


Populations & Cosmology

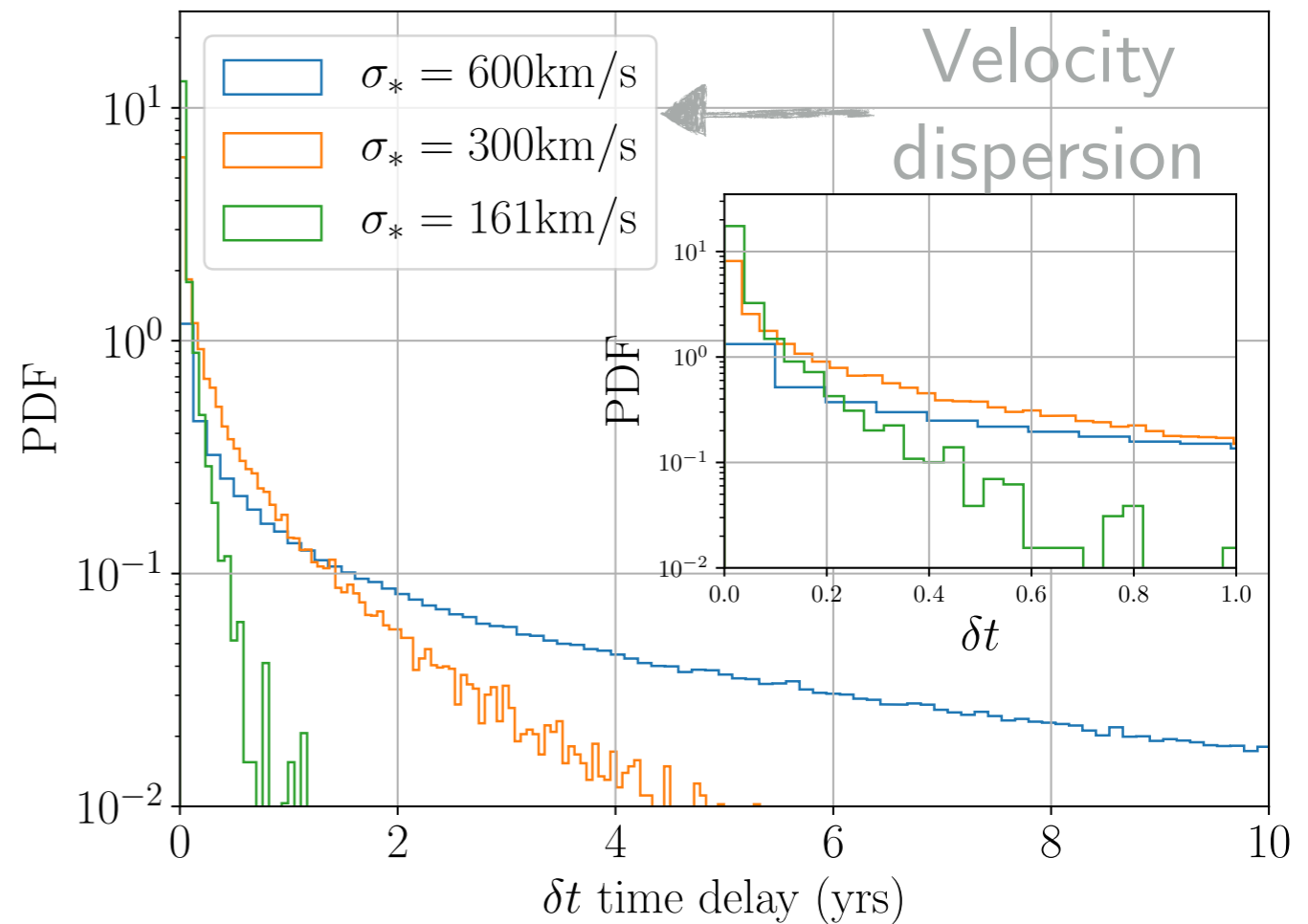
- Rates and time delay distributions inform about populations

Detected populations

aLIGO



Time delay distributions



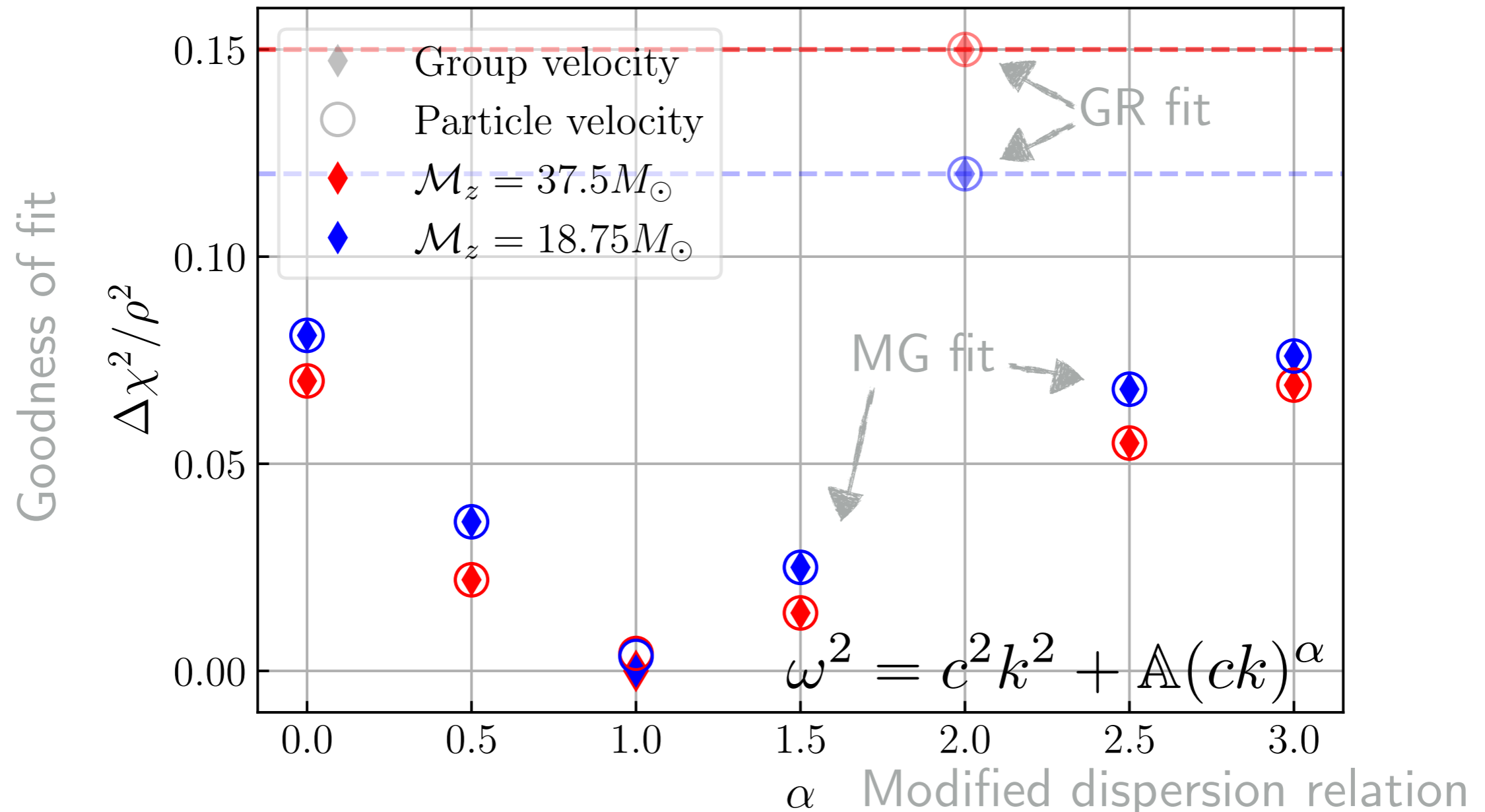
[Xu, Ezquiaga, Holz; ApJ'21]

- If you know the source and lens populations, rates and time delay distributions inform about cosmology [see Ajith's talk]

False violations of general relativity

- Lensed waveforms can be different from (unlensed) general relativity waveforms
- **E.g.** type II images

[Ezquiaga, *et al.*; JCAP'22]



Conclusions

Gravitational waves are precious cosmological probes:

- Well understood signals from general relativity
- Coherent detection of waveform
- Only distorted by gravitational lensing
- Current searches focus on repeated chirps and distorted waveforms
- No evidence so far, but first detections is *approaching!*
- Probing origin of the observed black holes and dark matter substructures with gravitational wave lensing



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